

**CRAFT:
the Cosmic Ray
Analytical Fast Tool**

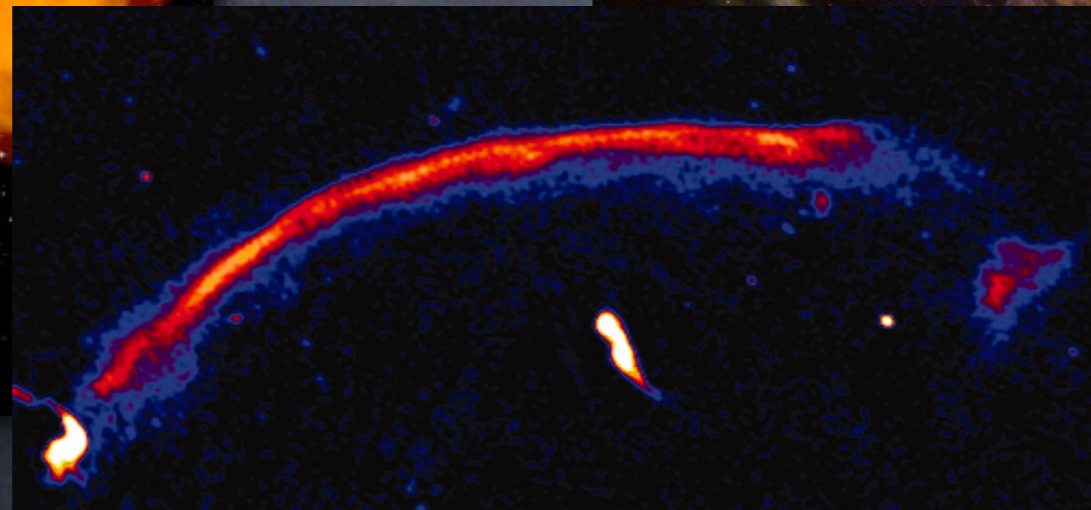
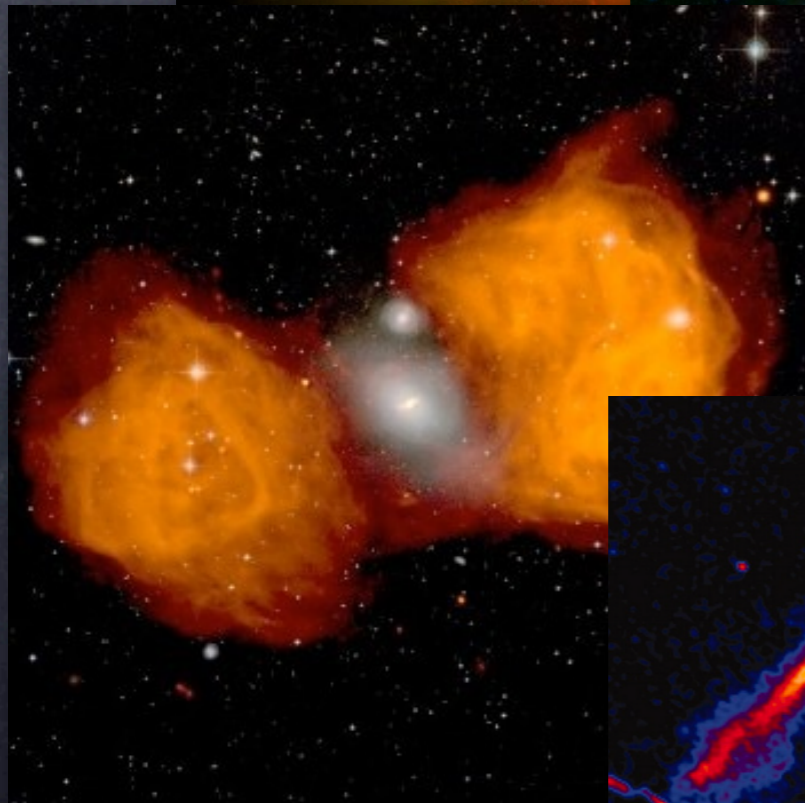
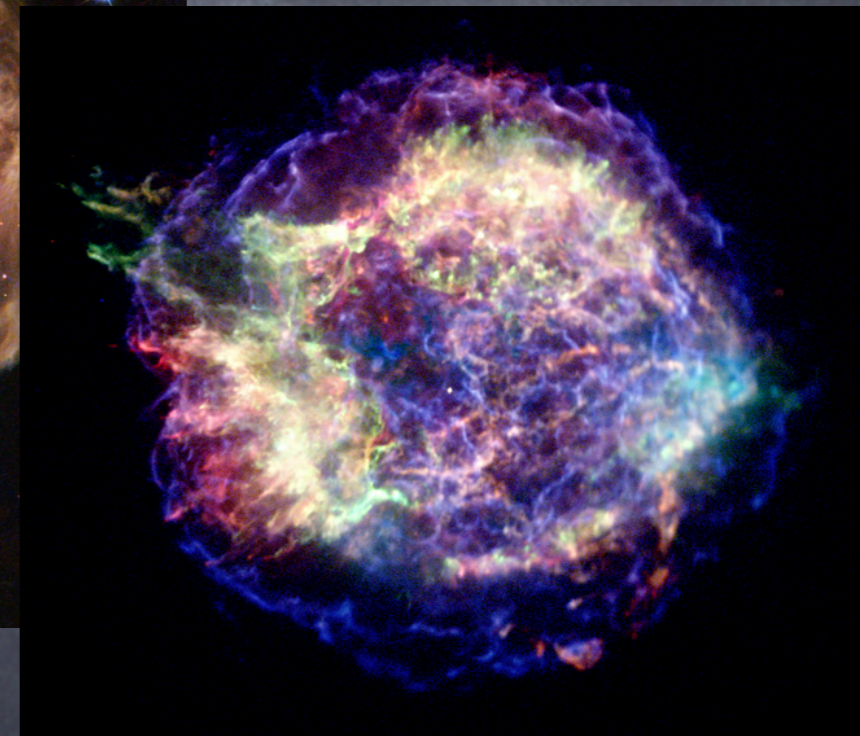
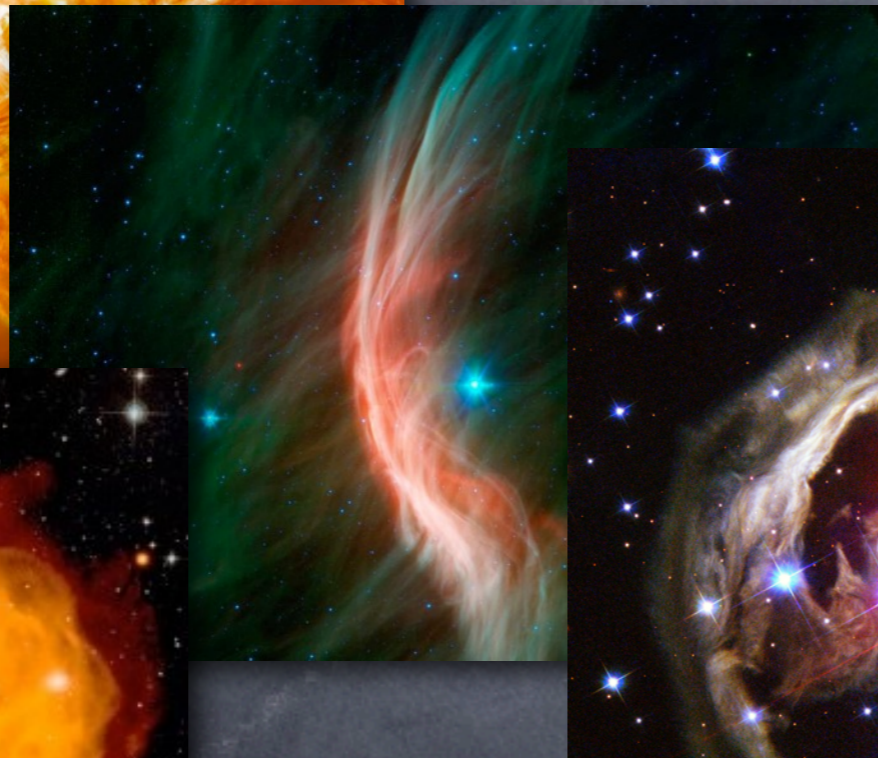
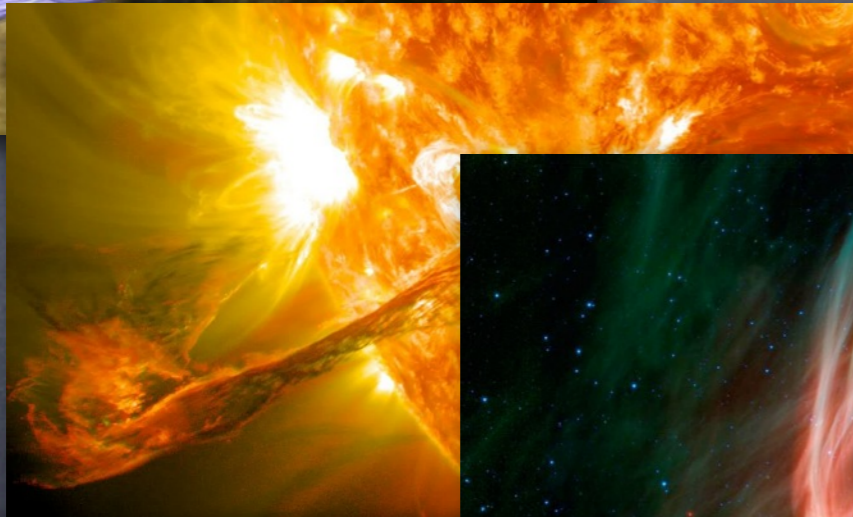
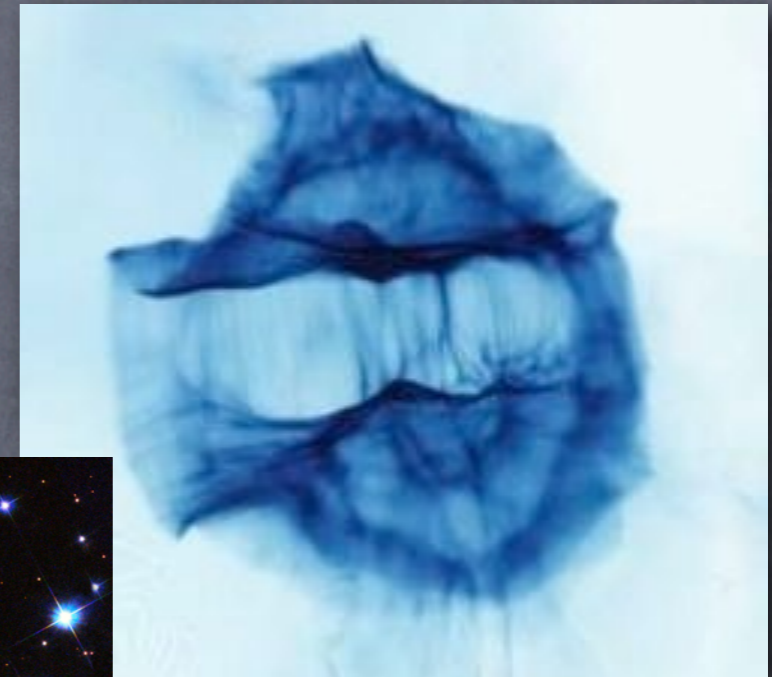
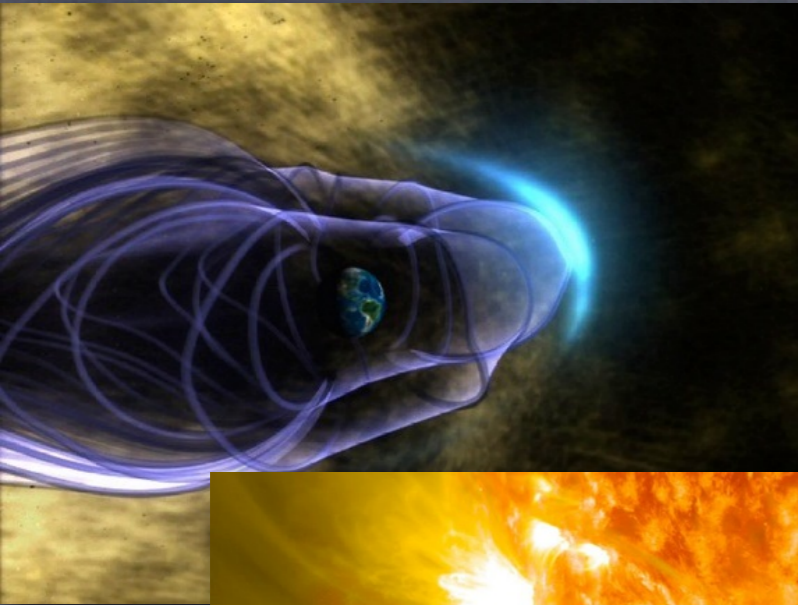
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A. Spitkovsky (Princeton)**

Collisionless shocks



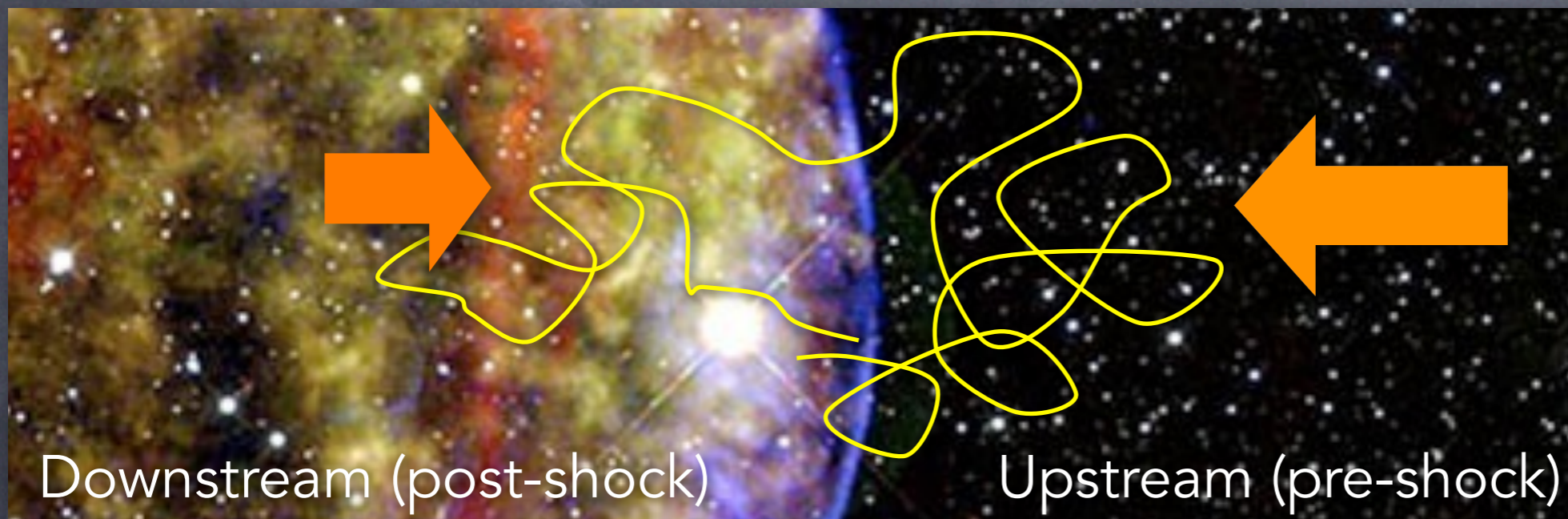
- Mediated by **collective** electromagnetic interactions
- Sources of **non-thermal** particles and emission
- Reproducible in laboratory



Diffusive Shock Acceleration



- **Fermi mechanism** (Fermi, 1949): random scattering leads to energy gain
- In **shocks** particles gain energy at any interaction (Blandford & Ostriker; Bell; Axford et al.; 1978): **Diffusive Shock Acceleration (DSA)**



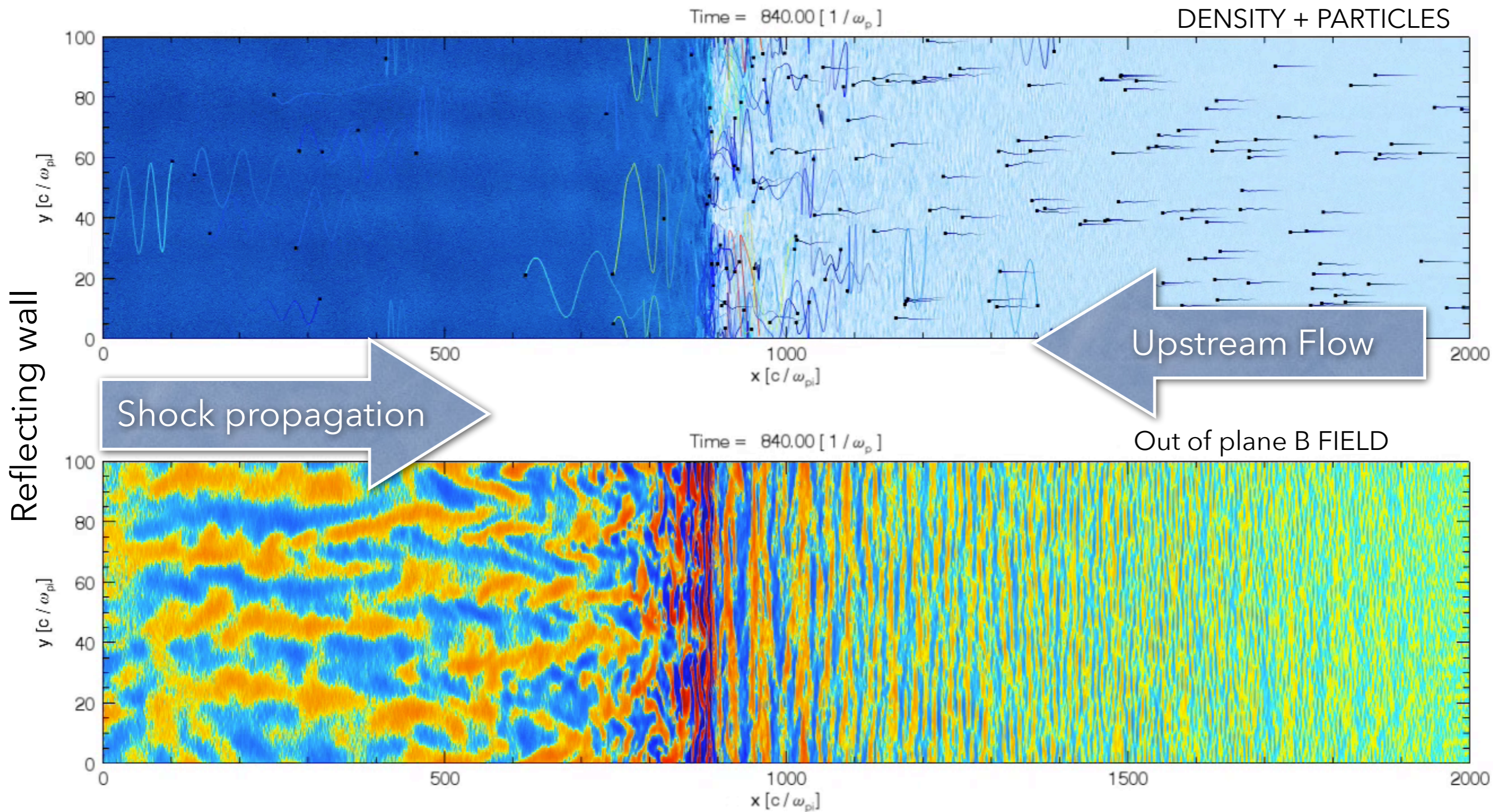
Test-particle
squeezed
between
converging
flows

- DSA produces **power-laws** $N(p) \propto 4\pi p^2 p^{-\alpha}$ in momentum, depending on the **compression ratio** $R = \rho_d / \rho_u$ only

$$R = \frac{4M_s^2}{M_s^2 + 3} \quad \alpha = \frac{3R}{R - 1}$$

- For strong shocks (Mach number $M_s \gg 1$): $R=4$ and $\alpha=4$

Kinetic simulations of collisionless shocks



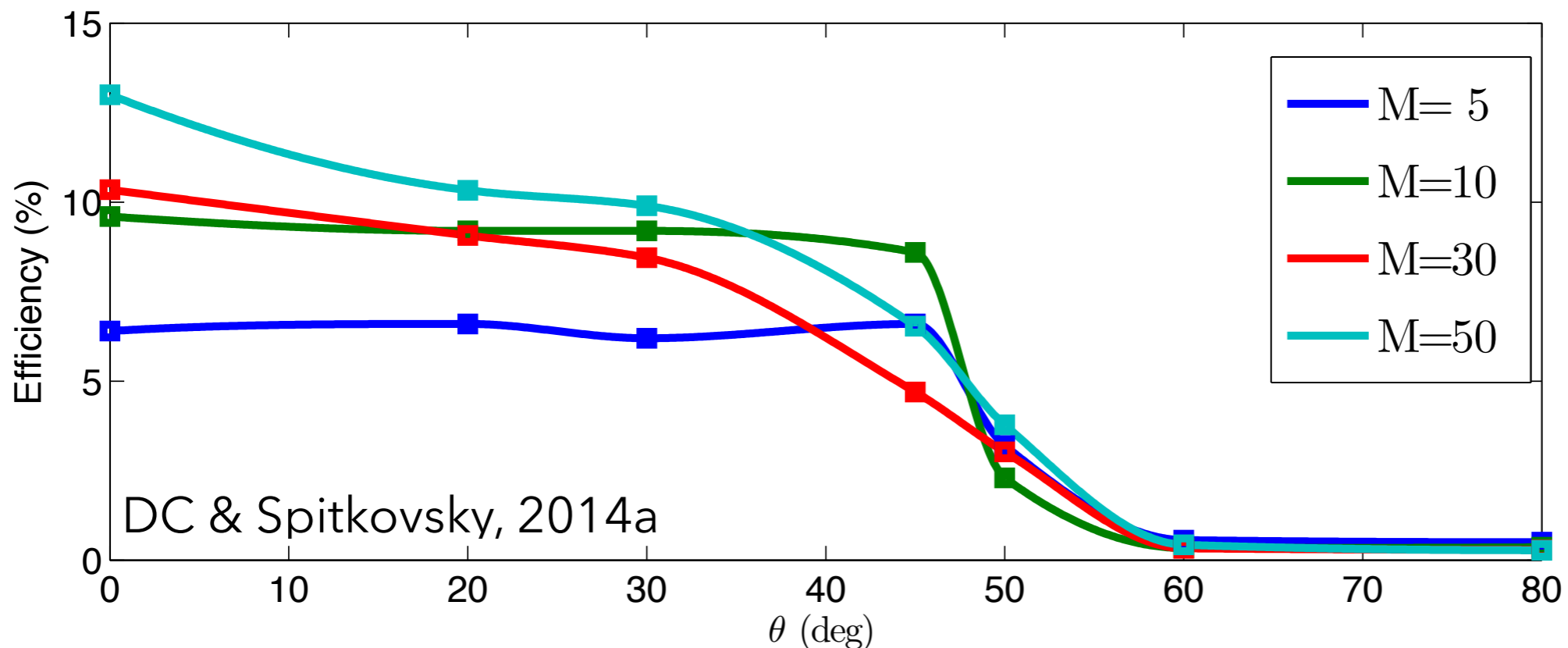
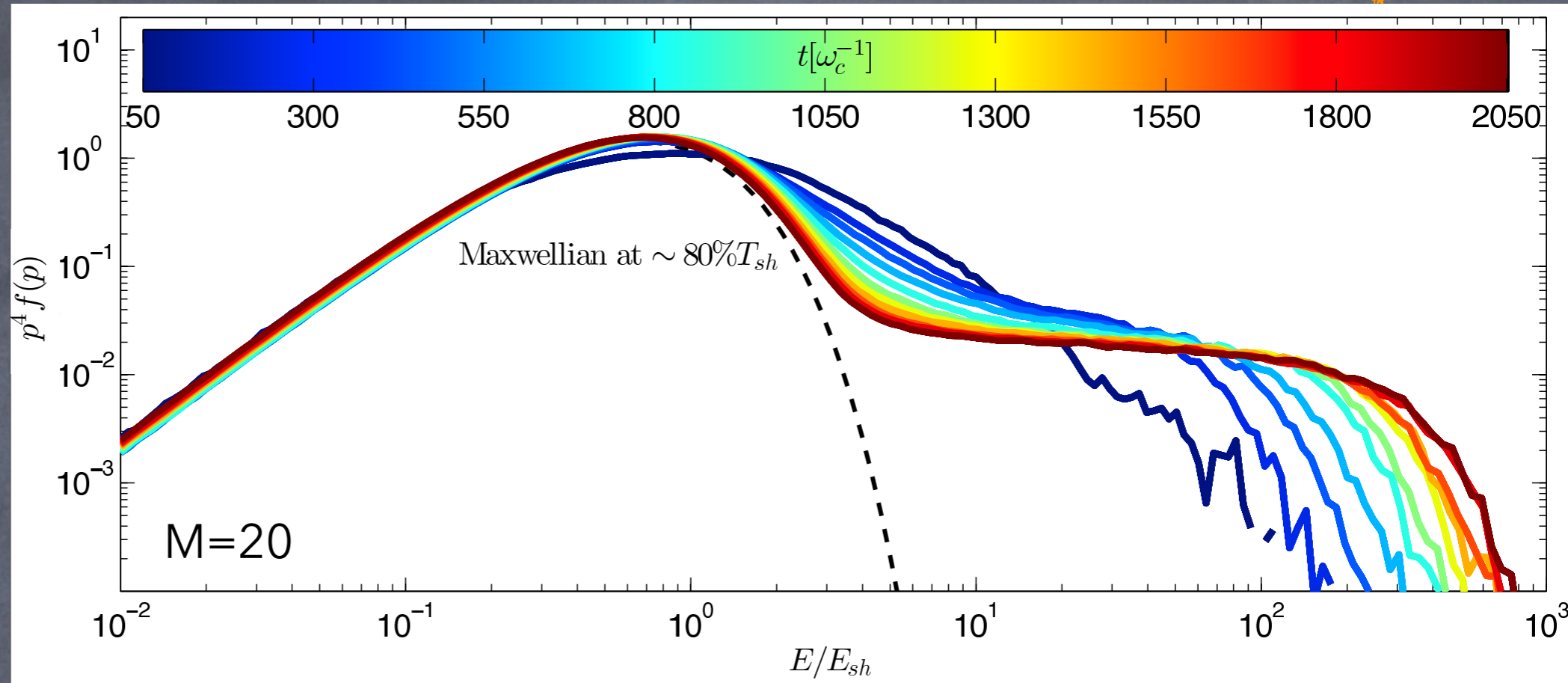
With **dHybrid** (DC & Spitkovsky 2014a,b,c)

Initial B field

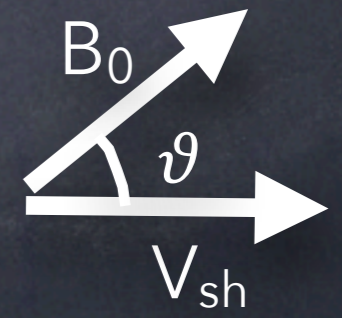
DSA efficiency



DSA from first-principles: ~15% of the shock kinetic energy converted in CRs with universal power-law spectrum



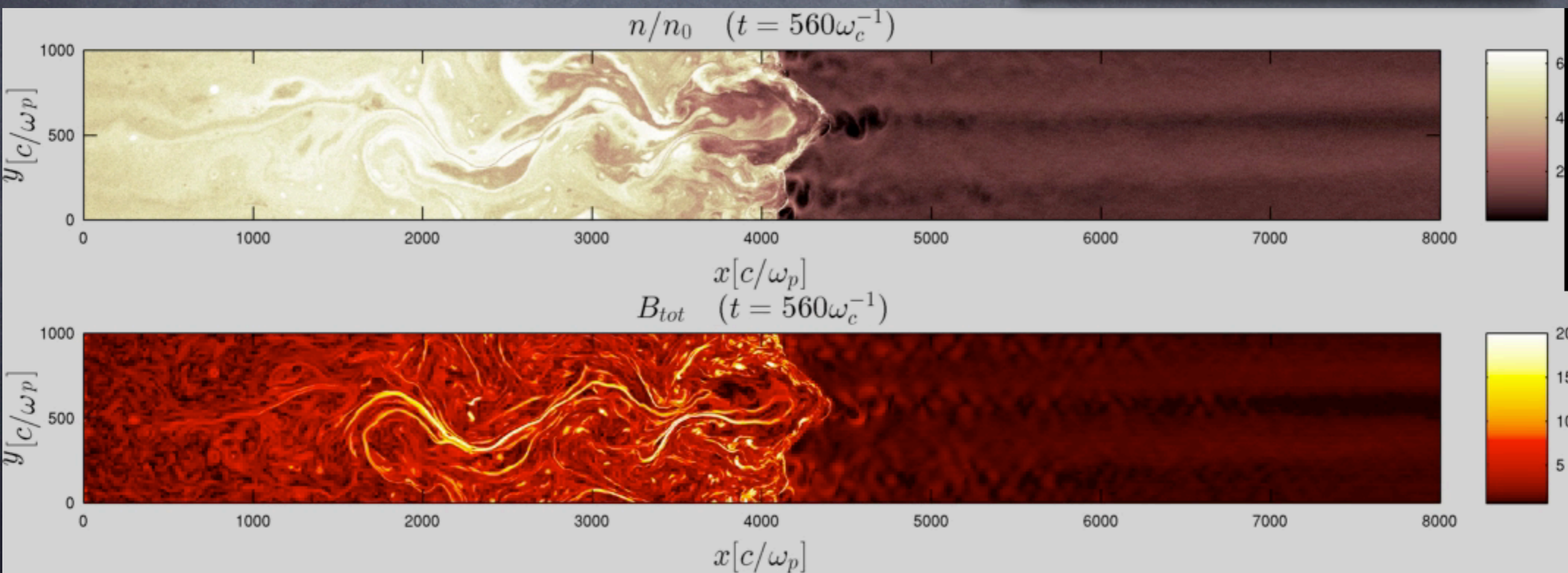
CR acceleration depends on shock strength and inclination



Non-linear DSA



- Efficient CR acceleration **modifies** the shock structure (Jones & Ellison 1991, Malkov & Drury 2001, DC 2012,...)
- Hydrodynamics** \longleftrightarrow CR spectra
- CR-driven **plasma instabilities**



Hybrid simulation of a strong parallel shock (DC & Spitkovsky 2014a)

Kinetic approaches to non-linear DSA



- Solve **CR transport** and shock hydrodynamics self-consistently

1D Diffusion-convection eq.

$$\frac{\partial f(t, x, p)}{\partial t} + \tilde{u}(x) \frac{\partial f(t, x, p)}{\partial x} = \frac{\partial}{\partial x} \left[D(x, p) \frac{\partial f(t, x, p)}{\partial x} \right] + \frac{p}{3} \frac{\partial f(t, x, p)}{\partial p} \frac{d\tilde{u}(x)}{dx}$$

- **FULLY NUMERICAL**: time-dependent
 - Kang & Jones 1997-2008; Berezhko & Völk 1997-2007; Zirakashvili & Aharonian 2009; ...
- **MONTE CARLO**: account for anisotropic distributions $f(p_x, p_y, p_z)$
 - Jones & Ellison 1991; Ellison et al. 1990;1995; Vladimirov, Ellison, & Bykov 2006; ...
- **SEMI-ANALYTICAL**: versatile, computationally extremely fast
 - Malkov 1997; Blasi 2002; Amato & Blasi 2006, DC et al. 2009; 2010; DC 2012, ...
- Require an **a priori** description of
 - **Magnetic field** generation
 - Particle **scattering** (diffusion)
 - **CR injection**

This information can be provided only by **kinetic simulations**

CRAFT: Cosmic-Ray Fast Analytic Tool



(DC et al. 2009-2015, to be publicly released soon)

- Iterative solution of the CR transport equation:

$$\tilde{u}(x) \frac{\partial f(x, p)}{\partial x} = \frac{\partial}{\partial x} \left[D(x, p) \frac{\partial f(x, p)}{\partial x} \right] + \frac{p}{3} \frac{d\tilde{u}(x)}{dx} \frac{\partial f(x, p)}{\partial p} + Q(x, p)$$

$$Q(x, p) = \eta \frac{\rho_1 u_1}{4\pi m_p p_{inj}^2} \delta(p - p_{inj}) \delta(x)$$

Injection

Mass+momentum conservation eqs.

$$\frac{p(x)}{\rho(x)^\gamma} = \frac{p_0}{\rho_0^\gamma};$$

$$f(x, p) = f_2(p) \exp \left[- \int_x^0 dx' \frac{\tilde{u}(x')}{D(x', p)} \right] \left[1 - \frac{W(x, p)}{W_0(p)} \right]$$

$$\Phi_{esc}(p) = -D(x_0, p) \left. \frac{\partial f}{\partial x} \right|_{x_0} = -\frac{u_0 f_2(p)}{W_0(p)};$$

$$W(x, p) = \int_x^0 dx' \frac{u_0}{D(x', p)} \exp \left[\int_{x'}^0 dx'' \frac{\tilde{u}(x'')}{D(x'', p)} \right].$$

$$f_2(p) = \frac{\eta n_0 q_p(p)}{4\pi p_{inj}^3} \exp \left\{ - \int_{p_{inj}}^p \frac{dp'}{p'} q_p(p') \left[U_p(p') + \frac{1}{W_0(p')} \right] \right\}$$

$$U_p(p) = \frac{\tilde{u}_1}{u_0} - \int_{x_0}^0 \frac{dx}{u_0} \left\{ \frac{\partial \tilde{u}(x)}{\partial x} \exp \left[- \int_x^0 dx' \frac{\tilde{u}(x')}{D(x', p)} \right] \left[1 - \frac{W(x, p)}{W_0(p)} \right] \right\}$$

u

$$\rho(x)u(x) = \rho_0 u_0$$

$$\rho(x)u(x)^2 + p(x) + p_{cr}(x) + p_B(x) = \rho_0 u_0^2 + p_{g,0} + p_{B,0}$$

$P_B + P_{cr}$

P_{cr}

$$2\tilde{u}(x) \frac{dp_B(x)}{dx} = v_A(x) \frac{dp_{cr}(x)}{dx} - 3p_B(x) \frac{d\tilde{u}(x)}{dx}$$

Magnetic turbulence transport eq.

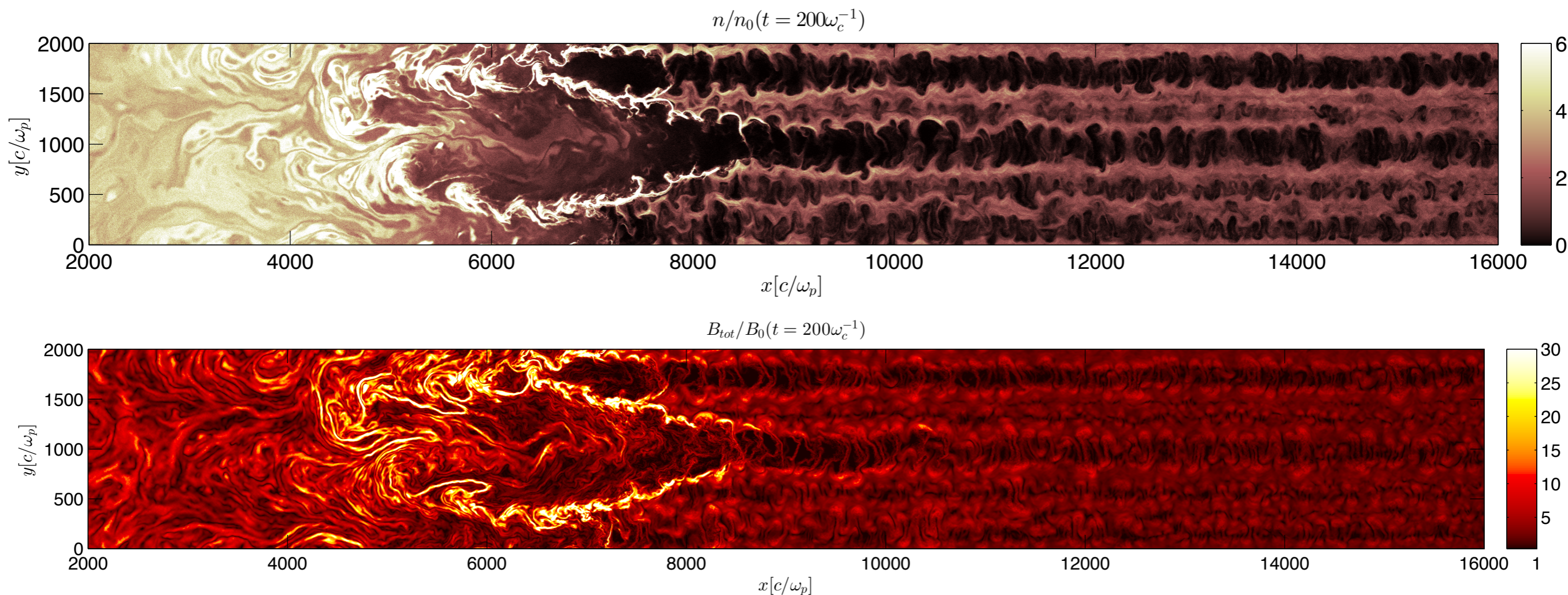
CR distribution function

- CR spectrum at the shock + spectrum of escaping CRs
- Very fast: a few seconds on a laptop

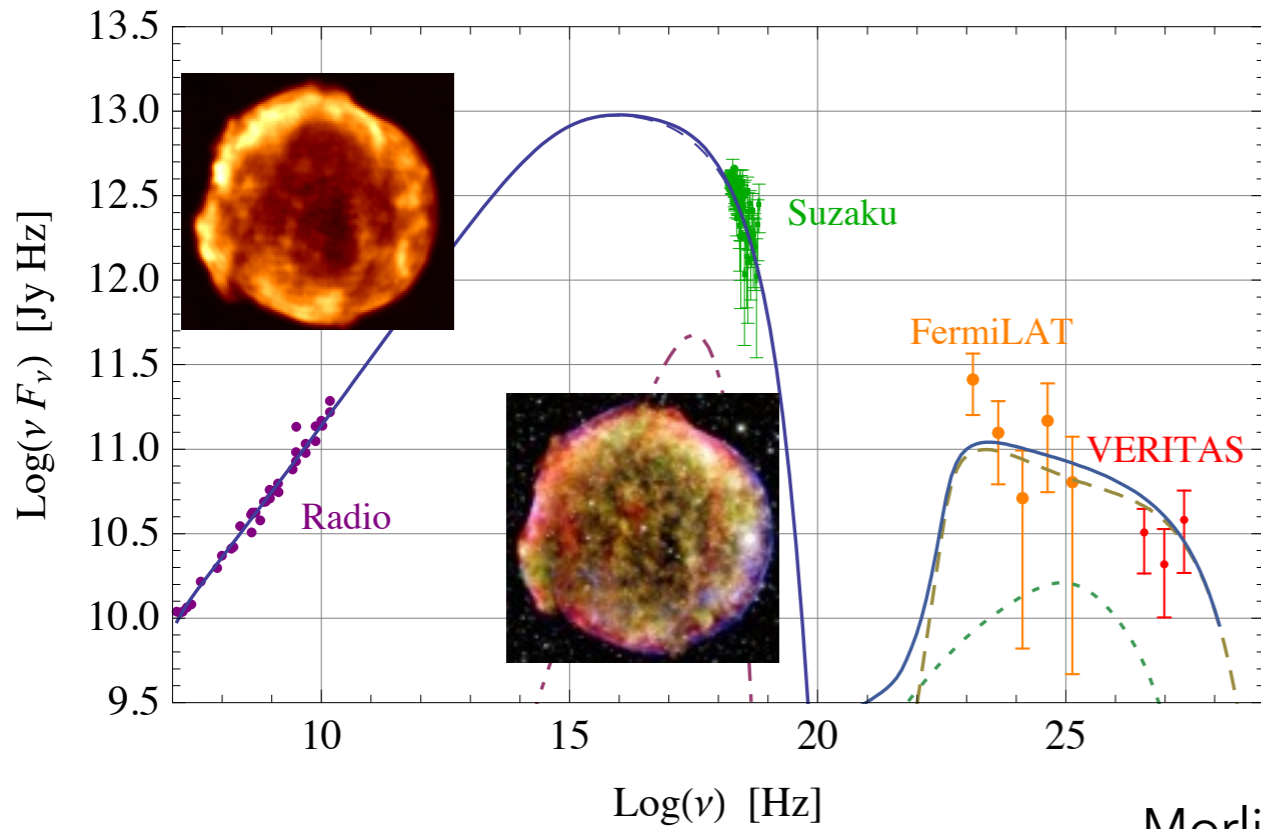
Microphysics under the hood



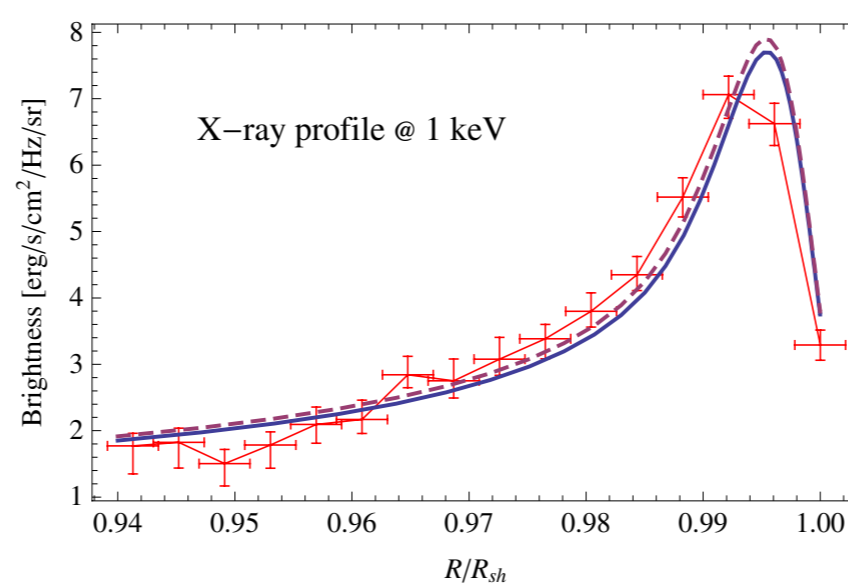
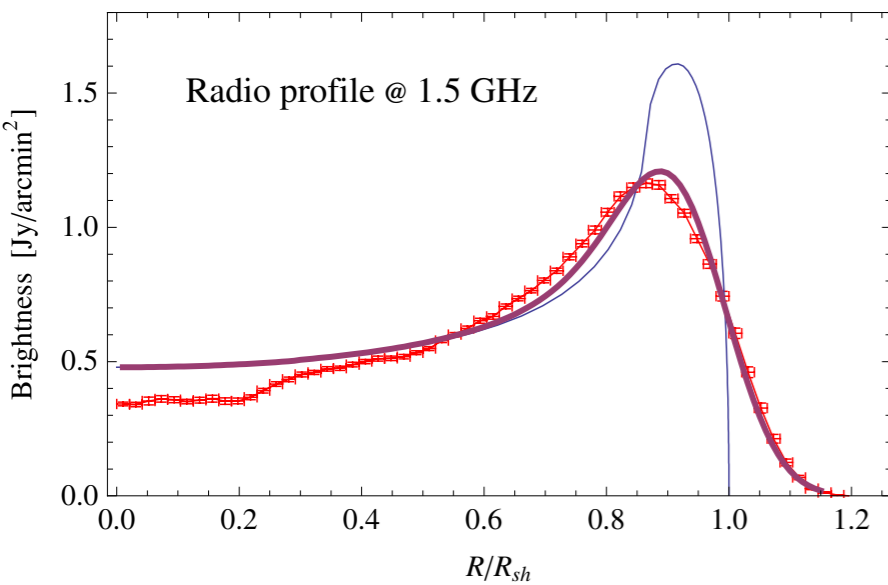
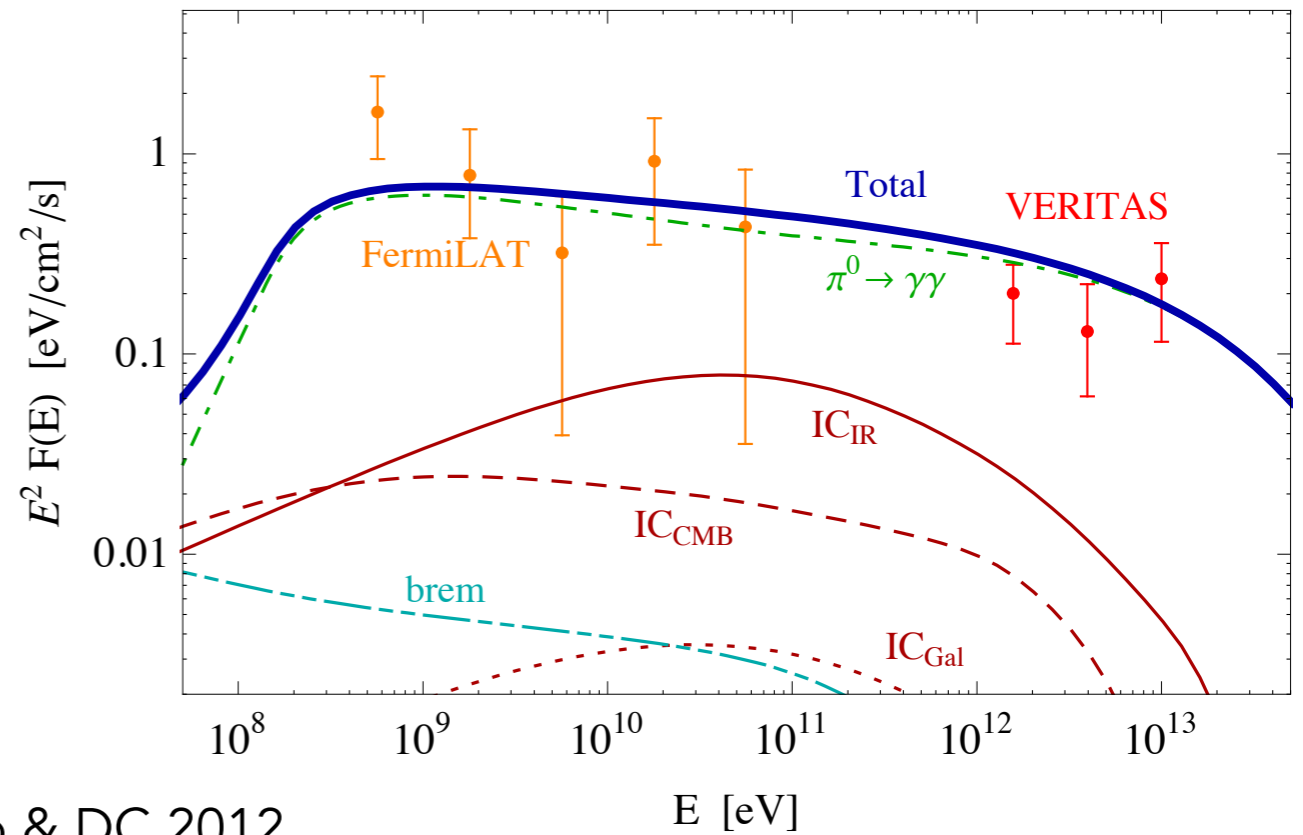
- **CR injection:** minimal model by DC, Pop & Spitkovsky 2015
- **B-field amplification:** resonant and Bell's streaming instabilities (DC & Spitkovsky 2014b)
- **Diffusion:** Bohm, in the amplified B (DC & Spitkovsky 2014c)
- **Electron/proton ratio:** from PIC simulations (Park, DC & Spitkovsky 2015, DC et al, in prep.)



A example of CRAFTwork: Tycho



Morlino & DC 2012



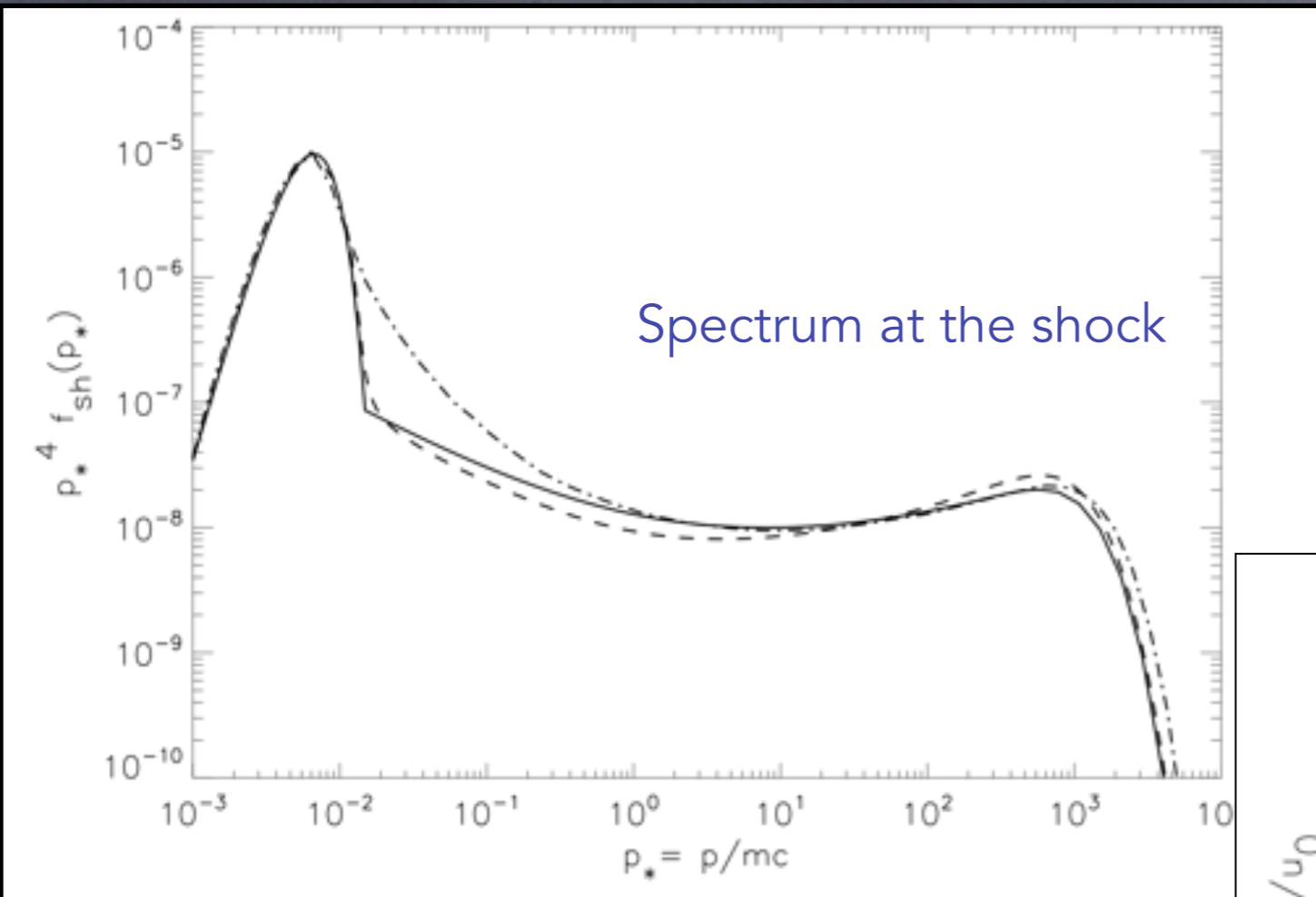
- Account for **spectra**, SNR hydrodynamics, and **morphology**
- Hadron acc. eff. **~10%**
- Protons up to **~0.5 PeV**

Only two free parameters: **injection efficiency** and **electron/proton ratio**

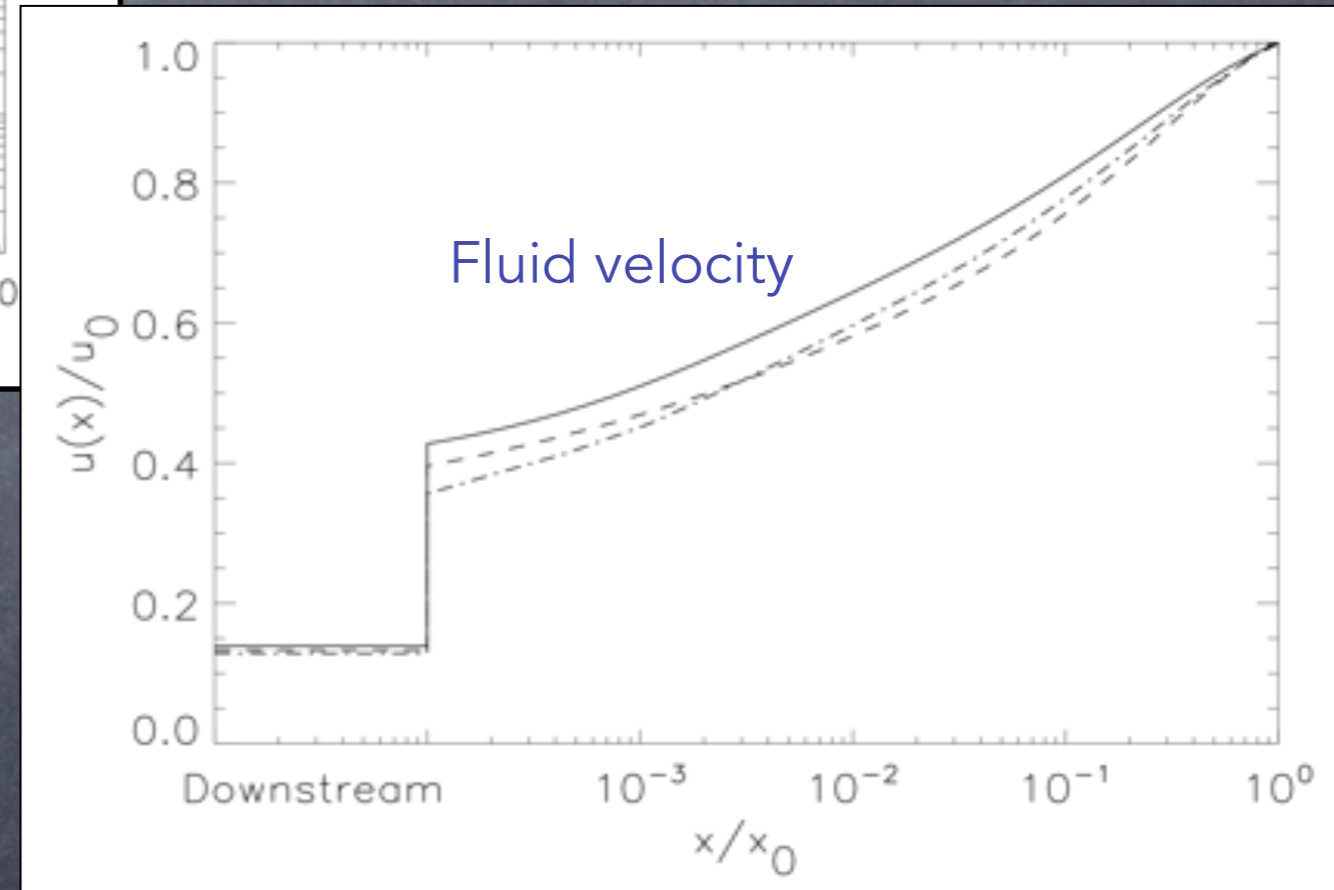
CRAFT

- Currently in **beta version** (email me if interested in testing)
 - soon on **GitHub** with extended **manual** and **references**
- Will come as:
 - **Stand-alone code** for SNR evolution
 - **Subroutine** for subgrid CR physics in hydro/MHD simulations
 - **Python interface + visualization suite** for particle and photon spectra
- Future developments:
 - **Nuclei** heavier than hydrogen (DC et al. 2010)
 - Propagation into **partially-ionized media** (Blasi et al. 2012, Morlino et al. 2013-2015)
 - ... and better and better **microphysics** from kinetic simulations!

Vs Numerical and Monte Carlo approaches



Solid: Caprioli et al. (Semi-Analytic)
Dashed: Kang & Jones (Numerical)
Dot-Dashed: Ellison et al. (Monte Carlo)



- Numerical (AMR): ~ 2 weeks
- Monte Carlo: ~ 3 days
- **Semi-analytic**: ~ 7 seconds

DC, Kang, Vladimirov & Jones, 2010