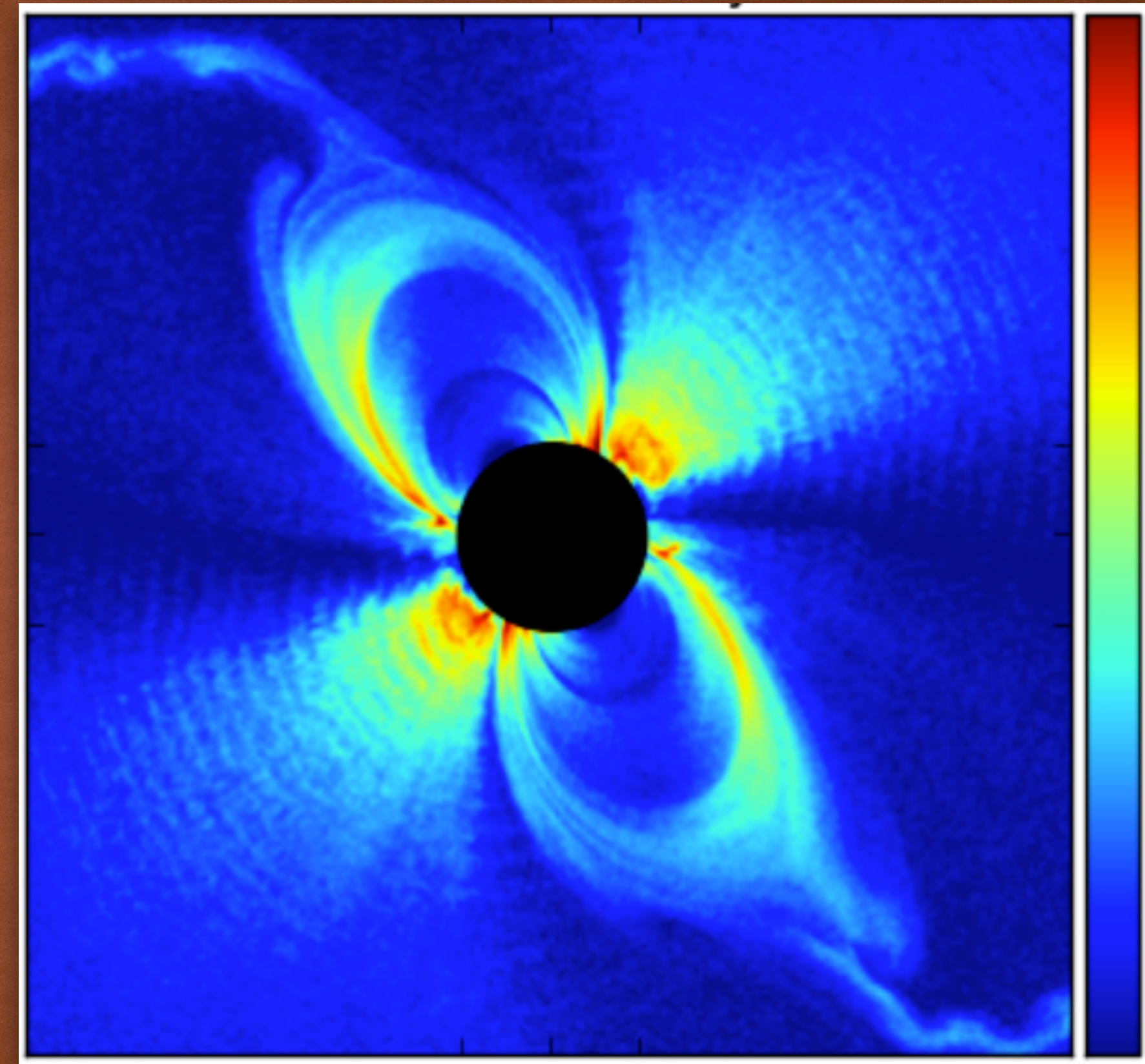
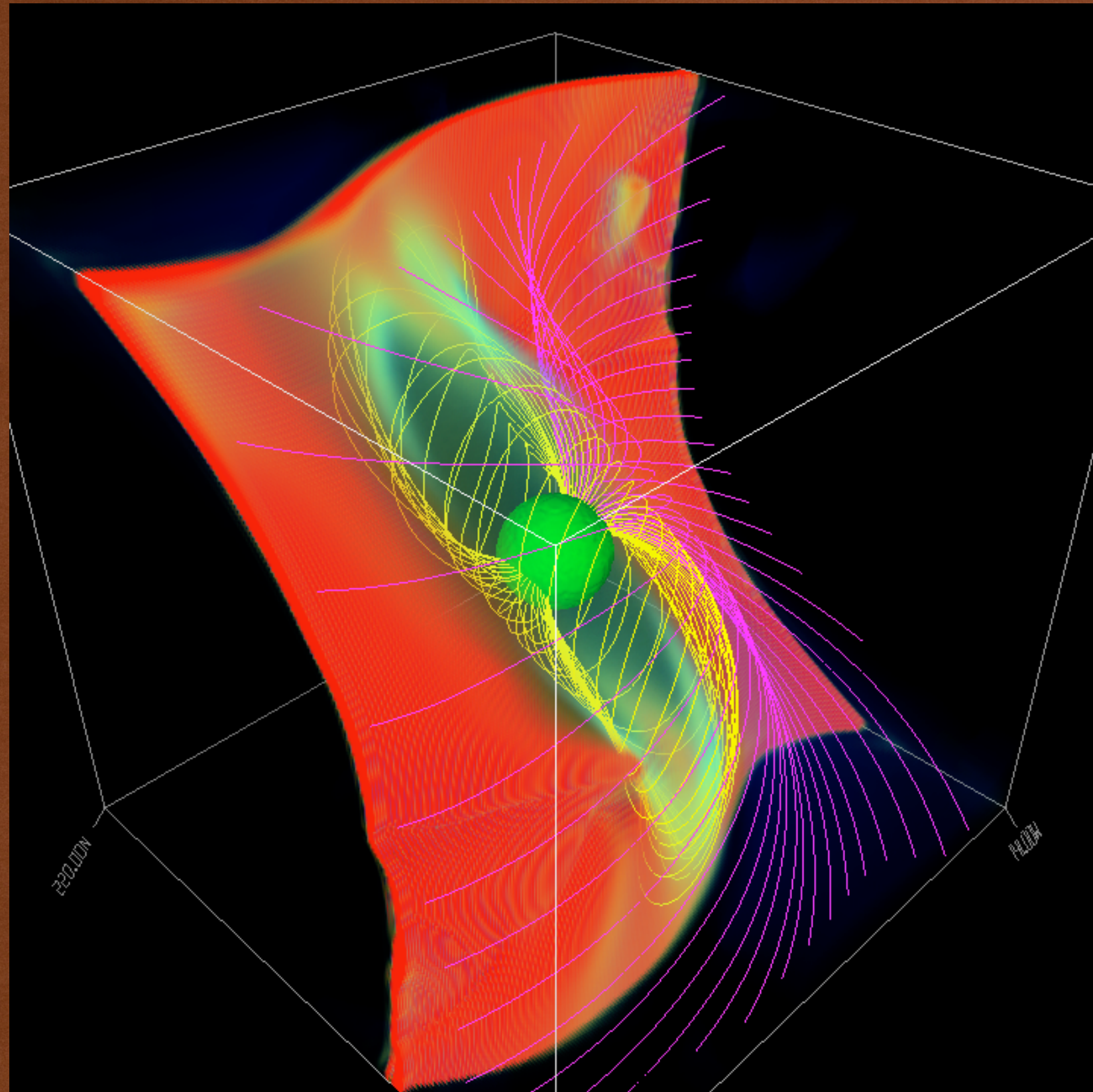


# AB-INITIO SIMULATIONS OF MAGNETOSPHERIC STRUCTURE OF PULSARS AND THEIR HIGH ENERGY EMISSION

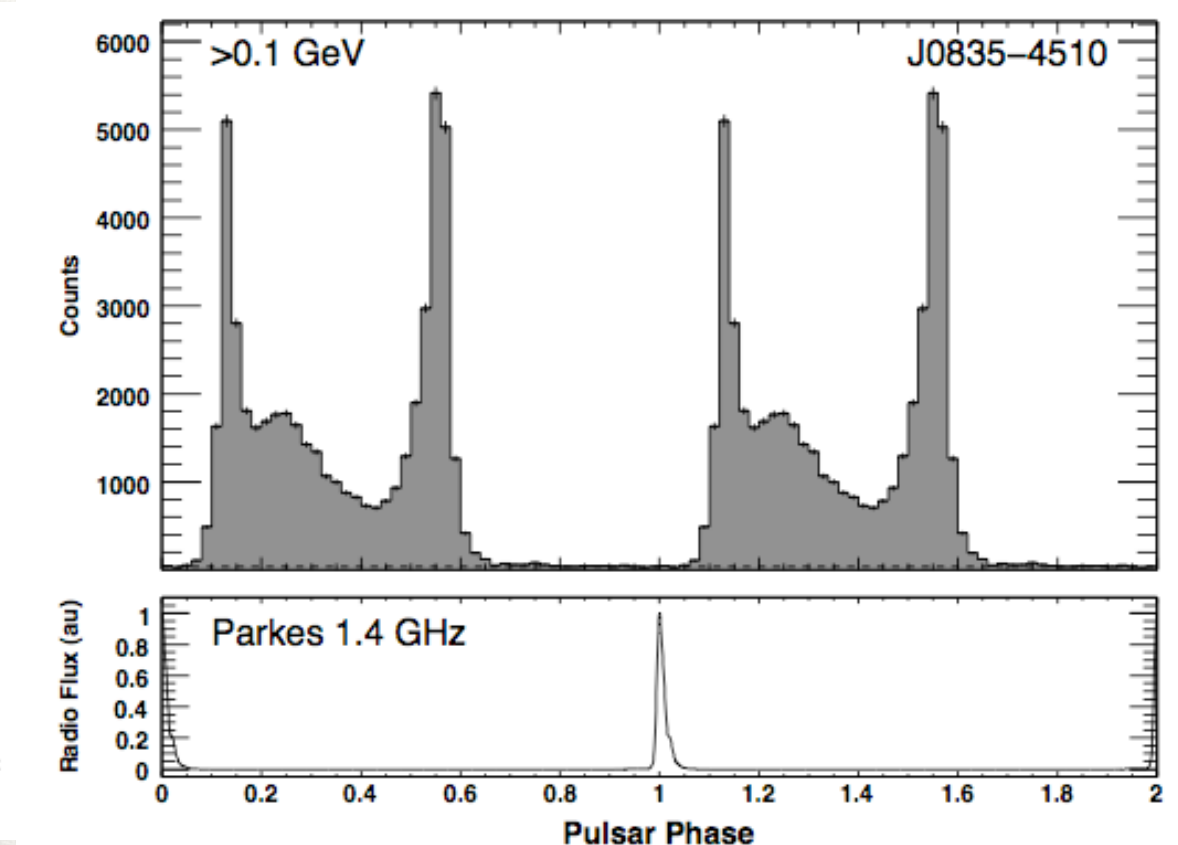
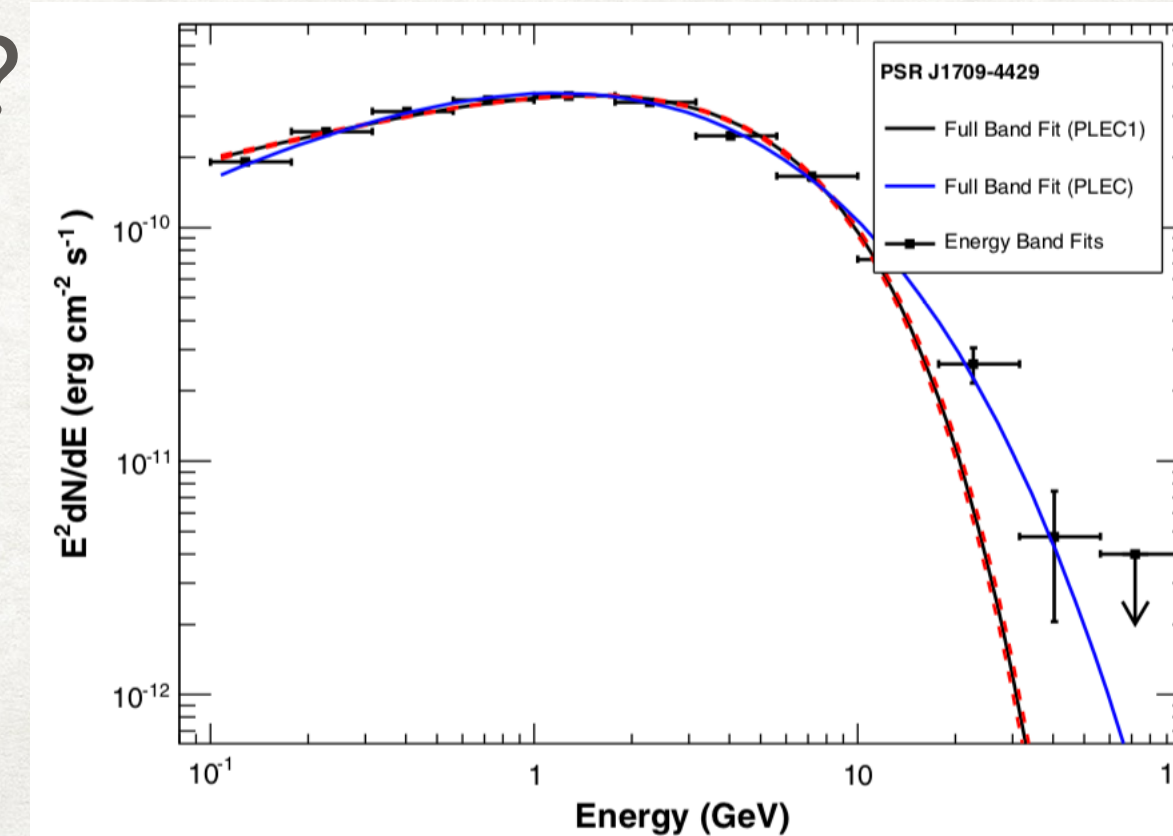
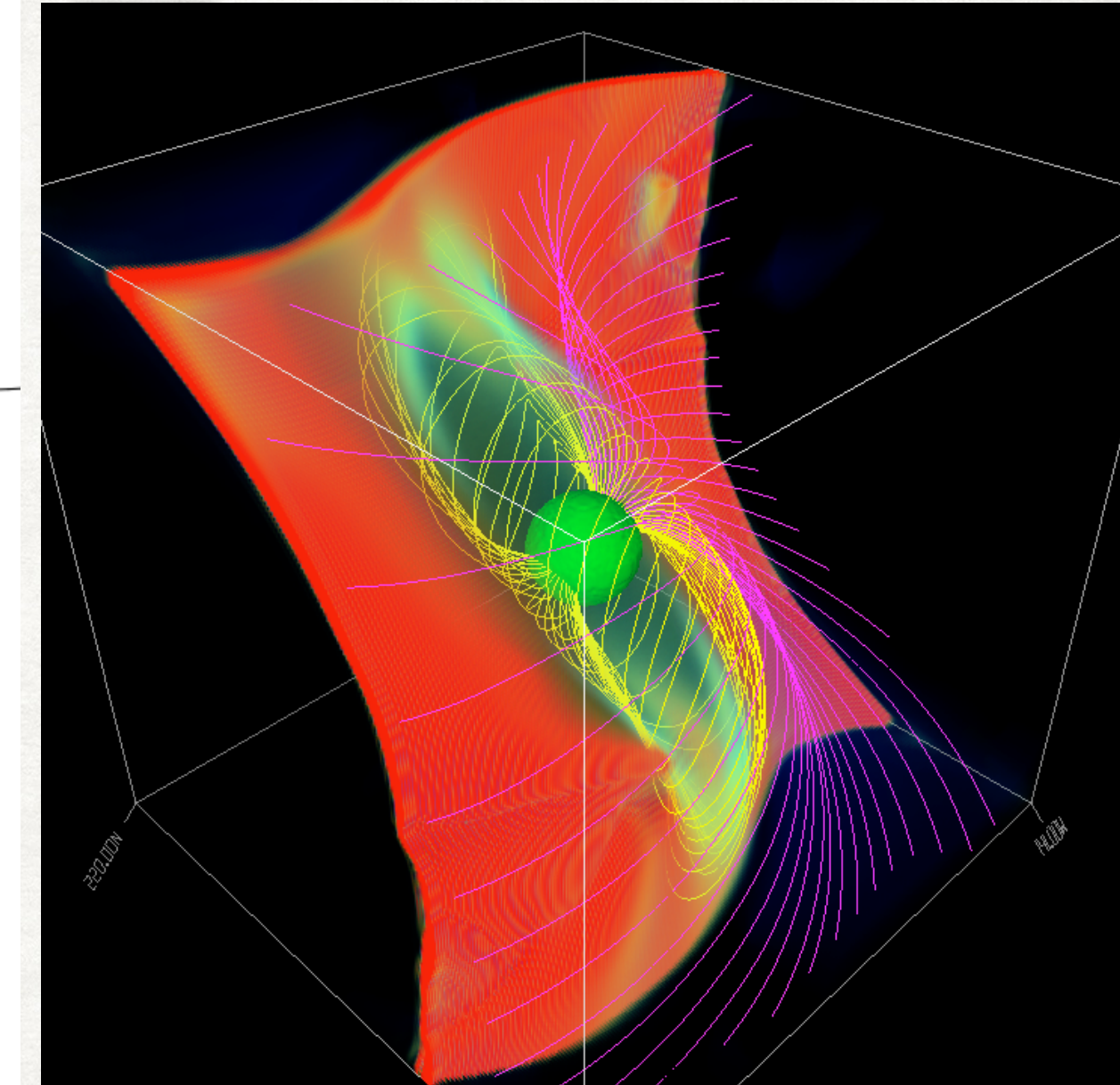
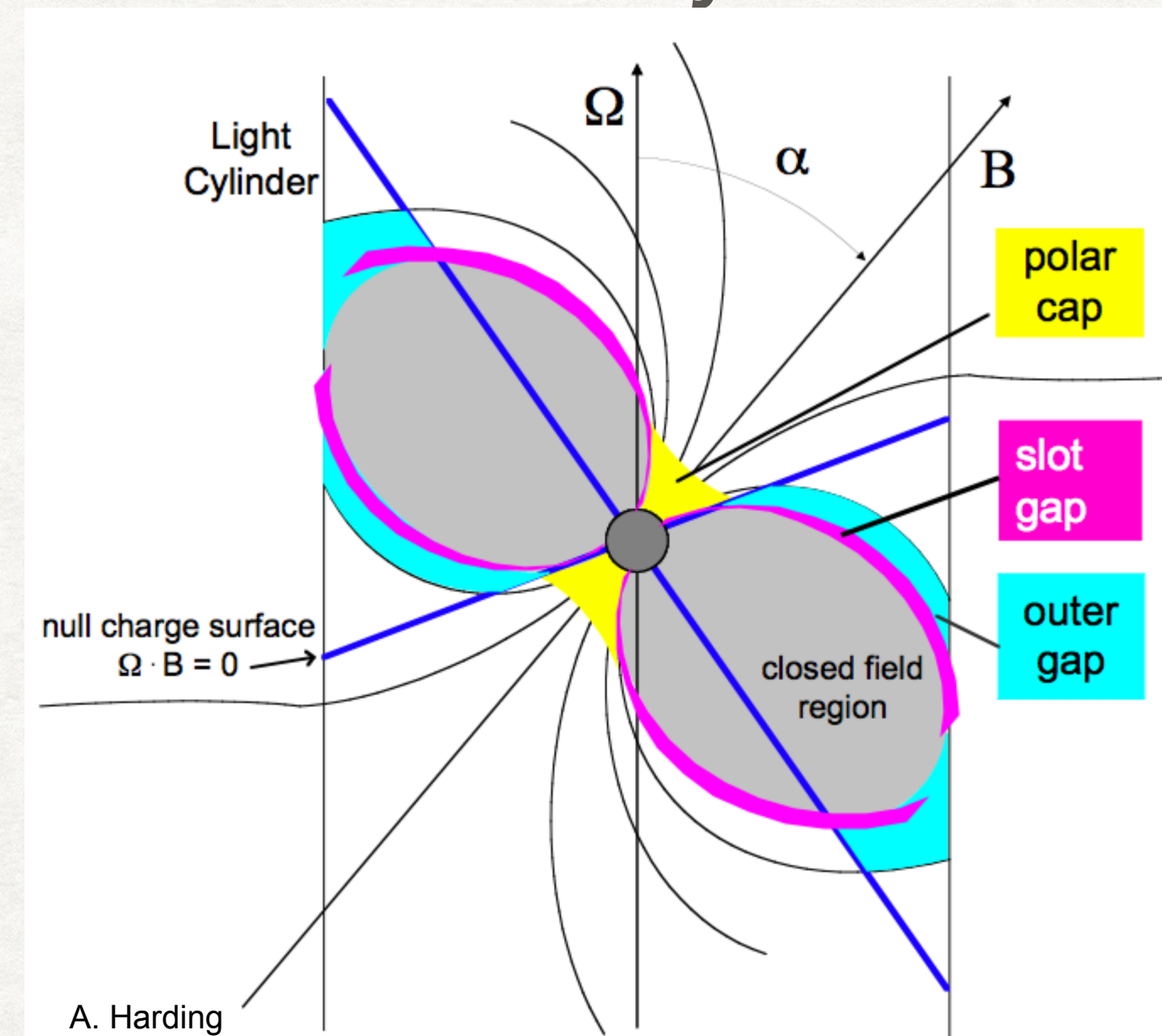


ANATOLY SPITKOVSKY, HAYK HAKOBYAN (PRINCETON)  
ALEXANDER PHILIPPOV (FLATIRON INST)



# FERMI PULSARS: OPEN QUESTIONS

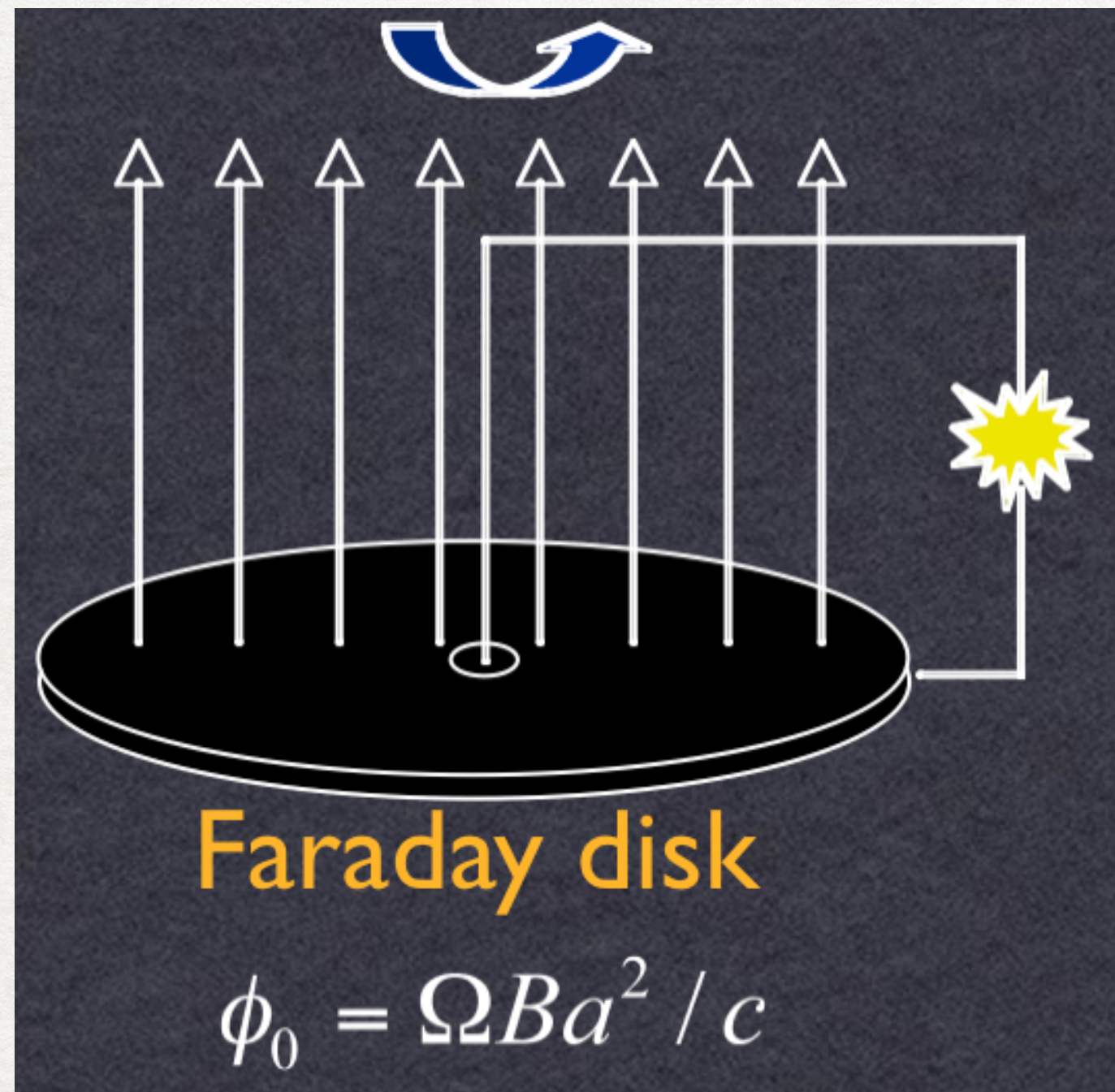
- Origin of gamma-ray emission: gaps, current sheets, synchrotron, curvature, IC?
- Light curves and spectra
- Luminosity-Edot relation
- Multiwavelength correlations
- What is the magnetospheric structure of pulsars, and how is the plasma supplied?



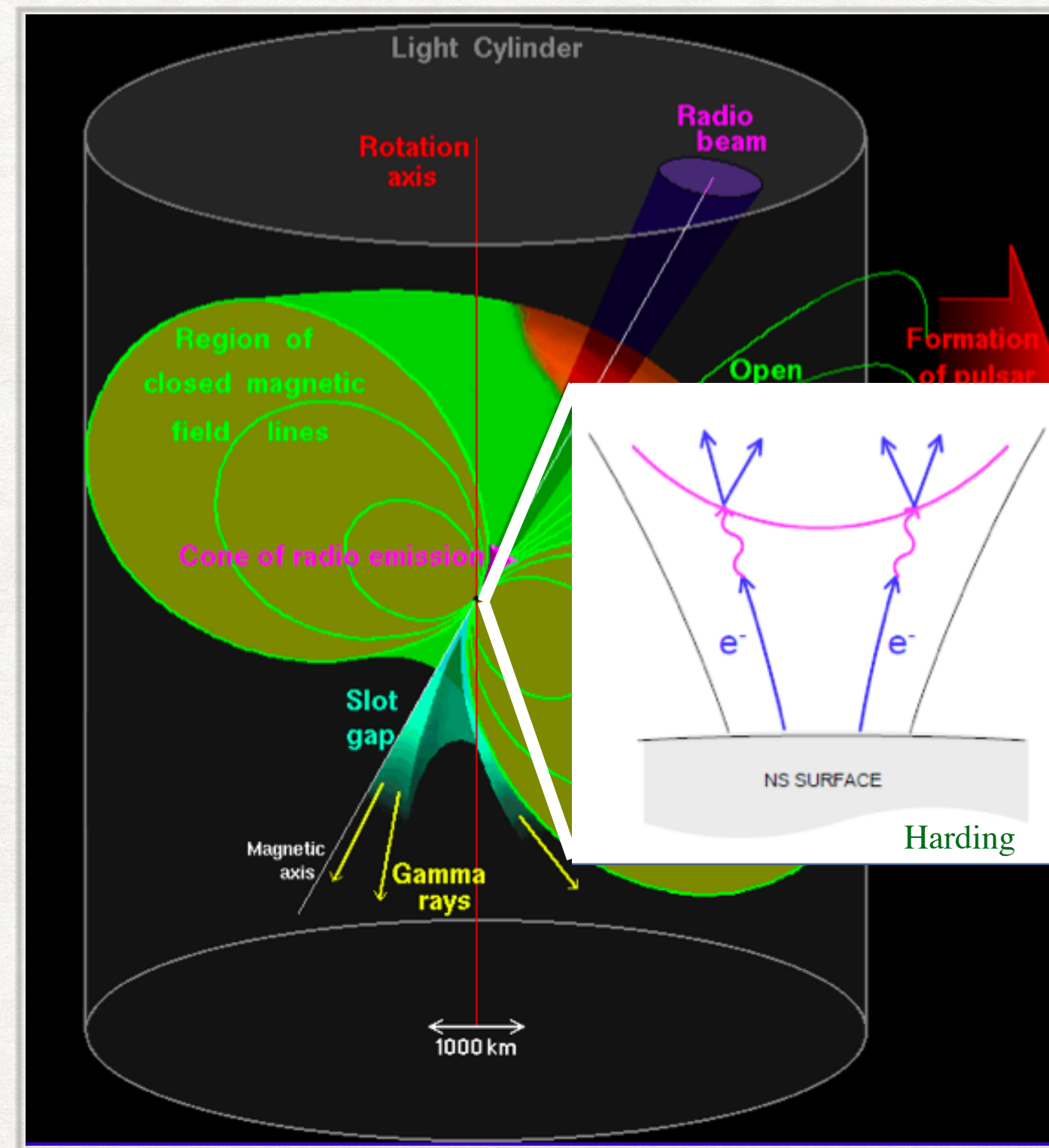


# PULSAR PROBLEM: AN OLD CHESTNUT

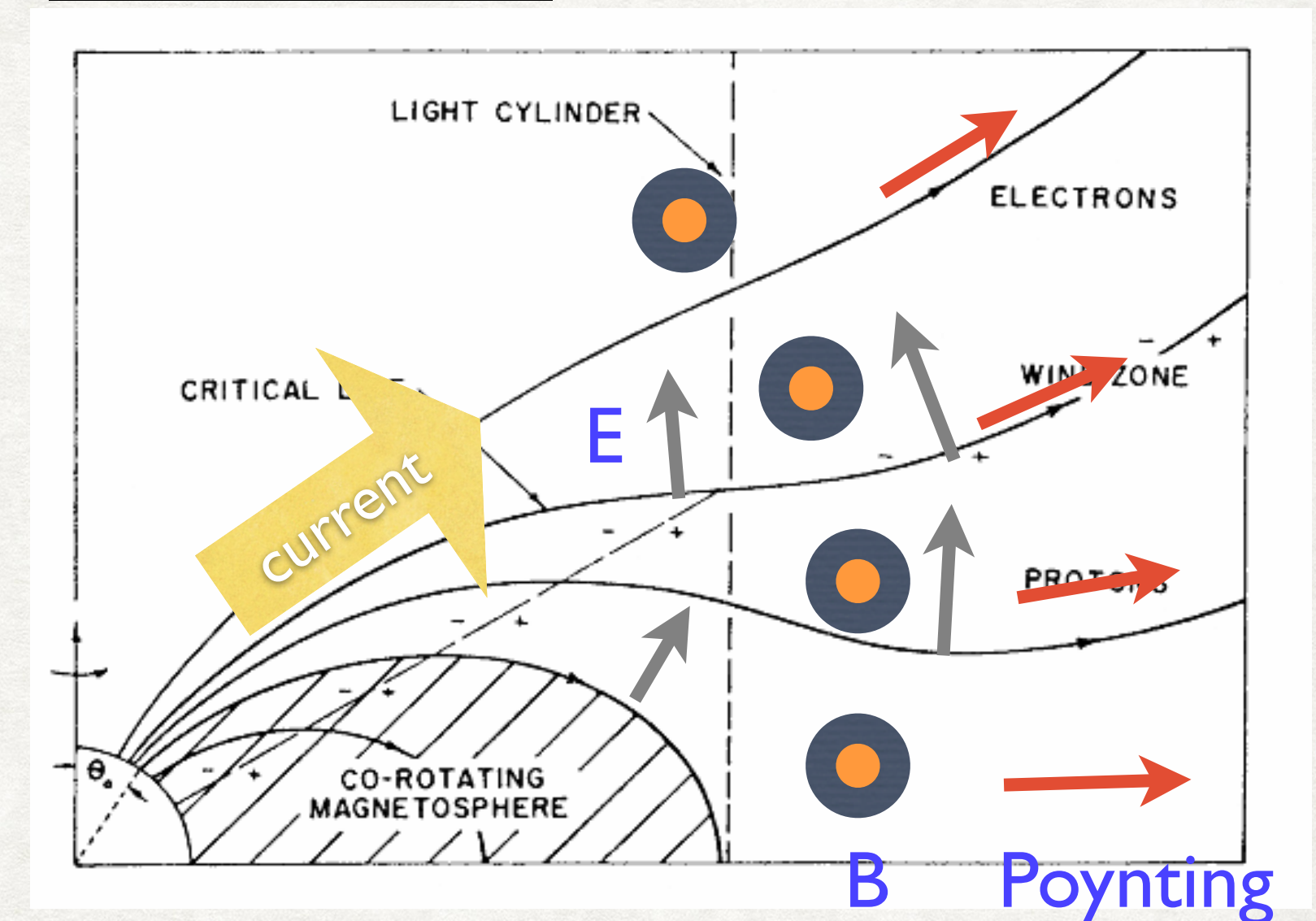
- What is the magnetospheric structure of a rotating magnetized conducting sphere with small surface work function?



Pole-equator potential



$$j_{GJ} = \rho_{GJ} c = -\frac{\vec{\Omega} \cdot \vec{B}}{2\pi}$$



Goldreich & Julian 1969

Current closure is essential



# PULSAR PROBLEM: METHODS AND APPROXIMATIONS I

- Force-free paradigm. Assume plasma is abundant and light.

$$\frac{1}{c} \frac{\partial \mathbf{E}}{\partial t} = \nabla \times \mathbf{B} - \frac{4\pi}{c} \mathbf{j}, \quad \frac{1}{c} \frac{\partial \mathbf{B}}{\partial t} = -\nabla \times \mathbf{E}$$

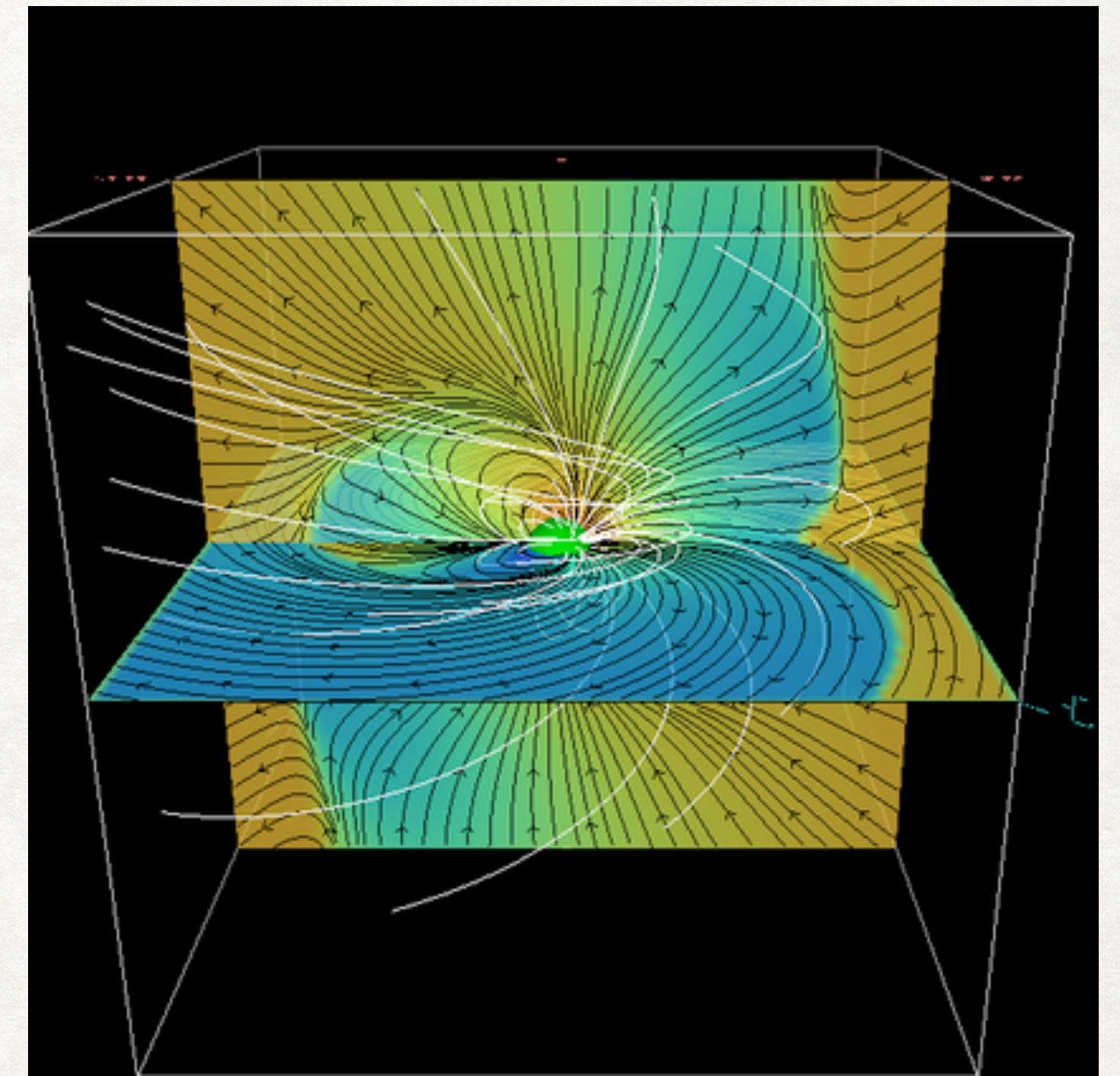
$$\rho_c \mathbf{E} + \mathbf{j} \times \mathbf{B} = \frac{d(\gamma \rho_m \mathbf{v})}{dt} + \text{pressure} \quad \mathbf{E} \cdot \mathbf{B} = 0$$

$$\mathbf{j} = \frac{c}{4\pi} \nabla \cdot \mathbf{E} \frac{\mathbf{E} \times \mathbf{B}}{B^2} + \frac{c}{4\pi} \frac{(\mathbf{B} \cdot \nabla \times \mathbf{B} - \mathbf{E} \cdot \nabla \times \mathbf{E}) \mathbf{B}}{B^2}$$

- Solution properties:

- Y-point
- Closed/open field lines
- Current sheet
- No pathologies at null surface and LC
- Predicts the spindown law
- Field lines are asymptotically radial

- All accelerating fields are shorted out
- Possible to extend to resistive limit (Li et al 2012, Kalapotharakos et al 2012-16)



Oblique: Spitkovsky (2006), Kalapotharakos et al (2009), Petri (2012), Tchekhovskoy et al. (2014) (full MHD)

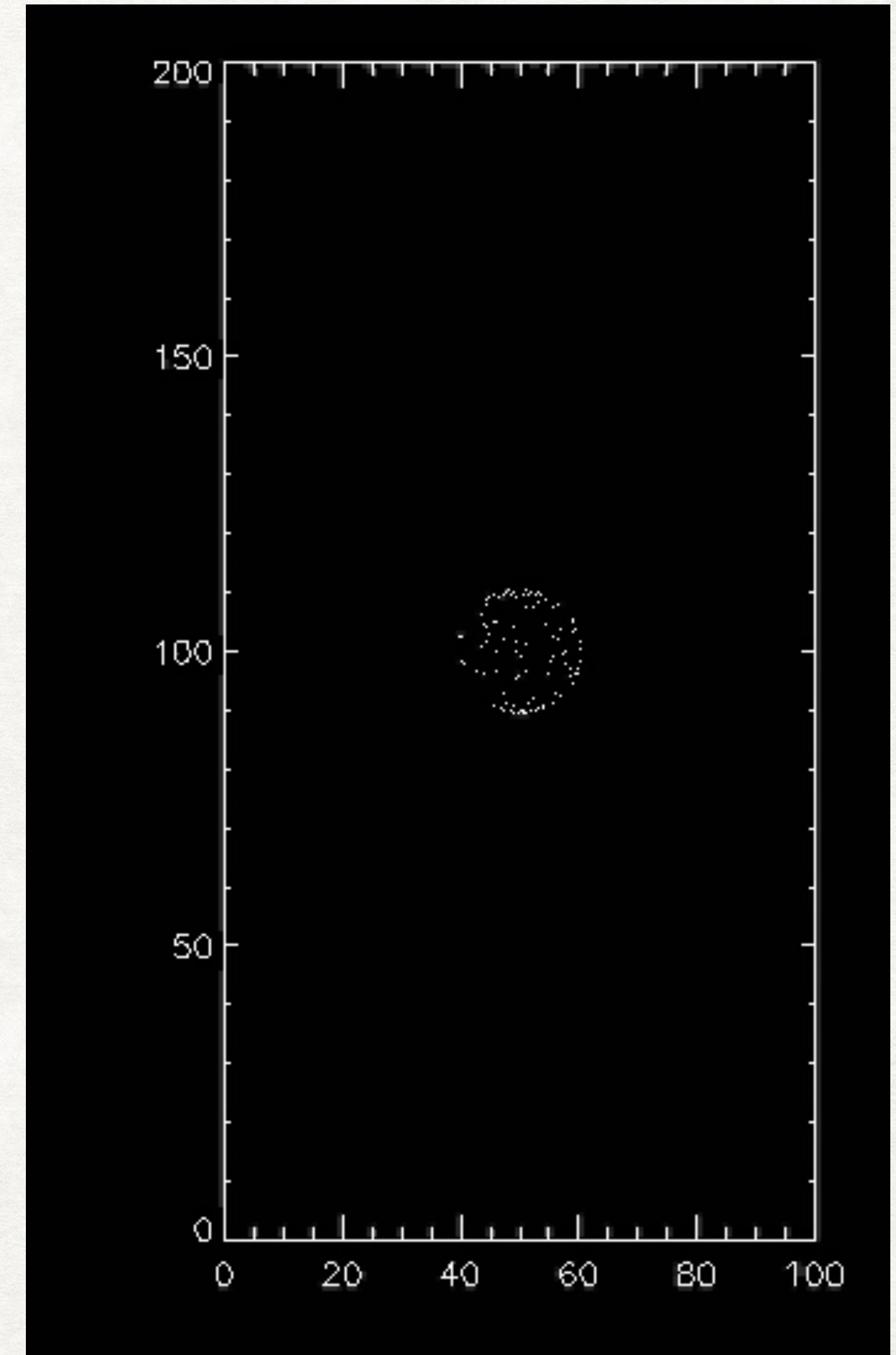
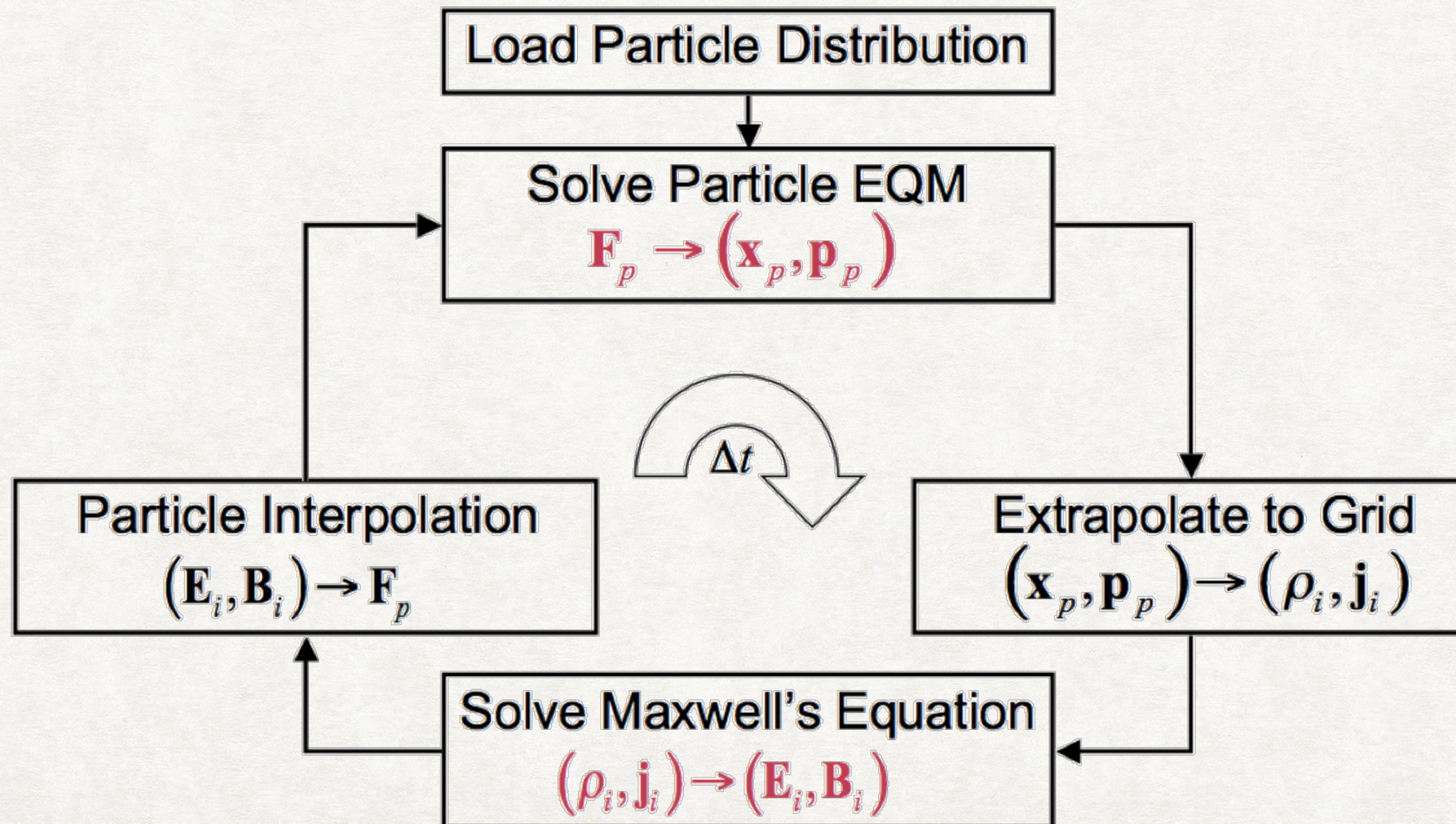
$$\dot{E} = \frac{\mu^2 \Omega^4}{c^3} (1 + \sin^2 \theta)$$



# PULSAR PROBLEM: METHODS AND APPROXIMATIONS II

- Kinetic model. Use particle-in-cell (PIC) method

$$\begin{aligned} \frac{\partial \mathbf{E}}{\partial t} &= c(\nabla \times \mathbf{B}) - 4\pi \mathbf{J}, & \nabla \cdot \mathbf{E} &= 4\pi \rho, & \nabla \cdot \mathbf{B} &= 0 \\ \frac{\partial \mathbf{B}}{\partial t} &= -c(\nabla \times \mathbf{E}), & \frac{d}{dt} \gamma m \mathbf{v} &= q(\mathbf{E} + \frac{\mathbf{v}}{c} \times \mathbf{B}) \end{aligned}$$



c.f. Krauss-Polstorff & Michel 1985, Spitkovsky & Arons 2002

GJ charge-separated solution is dead

Other groups using PIC: Kalapotharakos et al, Chen et al, Belyaev et al



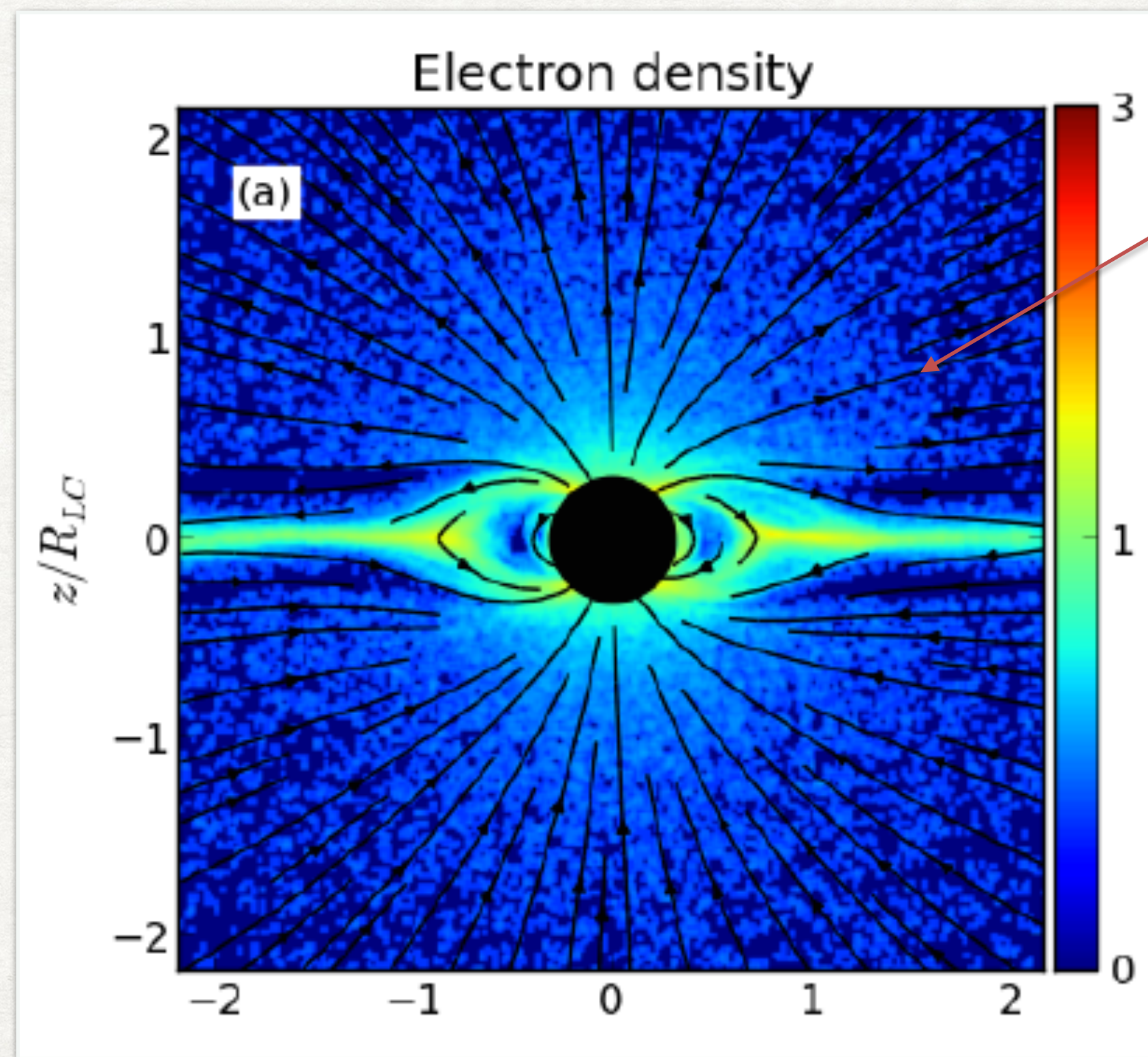
# SOLUTIONS WITH PAIR PRODUCTION

- Add pair production with threshold based on particle energy in the inner magnetosphere. Outer magnetosphere: pair production in photon-photon collisions; keep track of photons

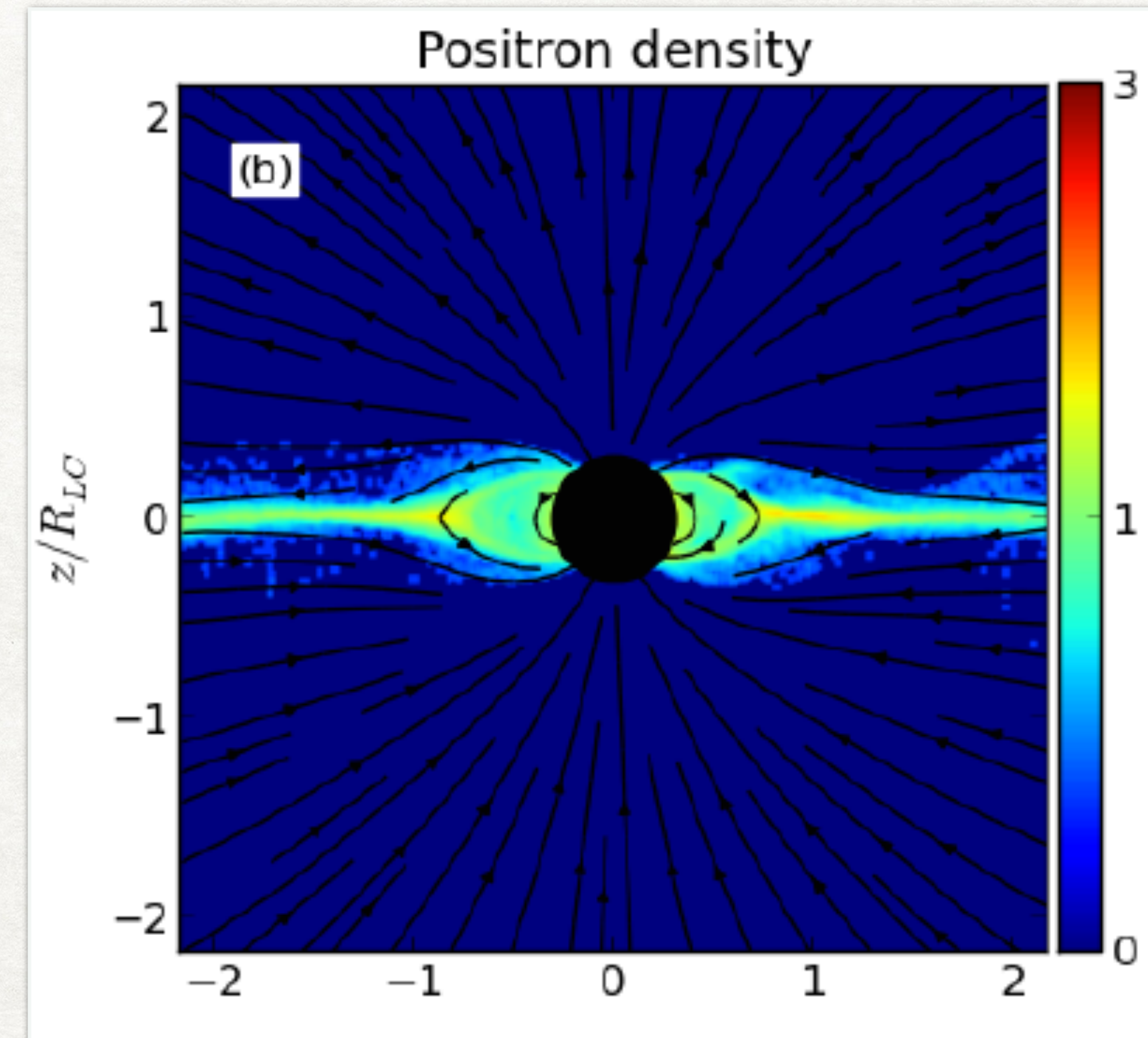
$$R_*/(c/\omega_p) \approx 30 - 40 \gg 1$$

$$R_{LC}/R_* = 3 - 5$$

$$\Phi_{PC} = \mu\Omega^2/c^2 \approx 500 \gg \gamma_{\text{threshold}} = 40$$



$$j < \rho_{GJ} c$$



- $j < \rho_{GJ} c$  is satisfied by non relativistic outflow of electrons!
- Approaches force-free field solution, but no polar pair production!

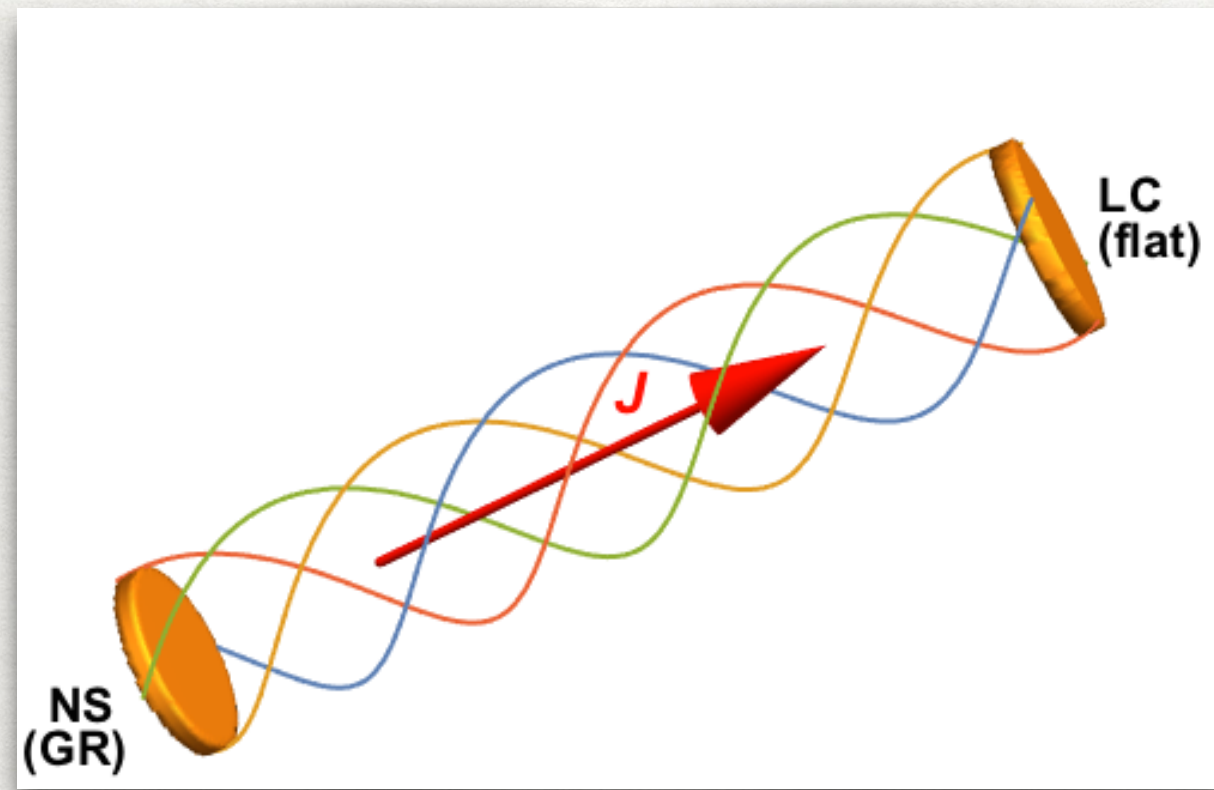
Philippov et al. (2015a)

Chen, Beloborodov (2014)



# SOLUTION: IT'S A *MASSIVE* ROTATING SPHERICAL CONDUCTOR!

- Add GR frame-dragging effect

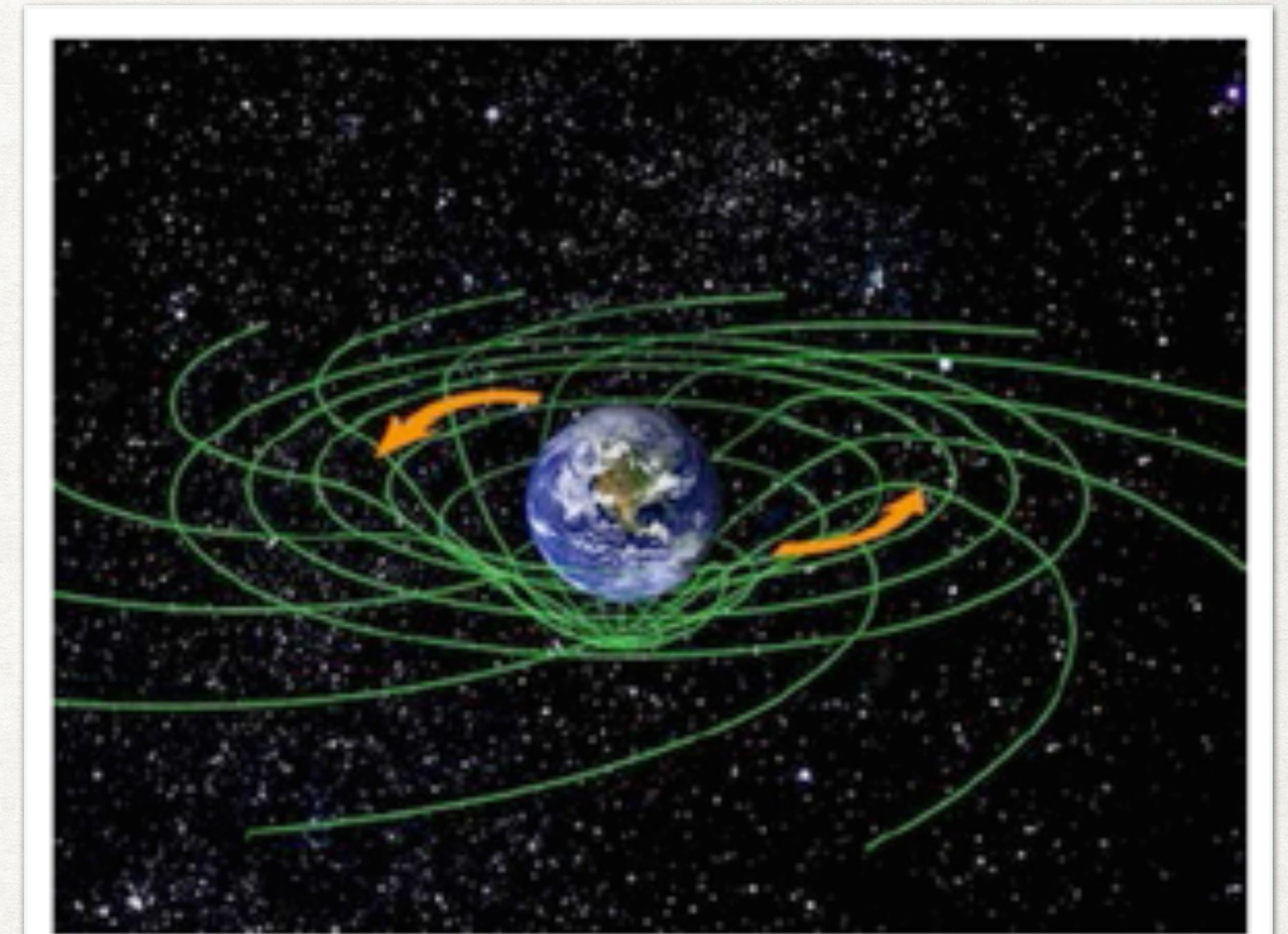


$$\omega_{LT} = \frac{2}{5} \Omega_* \frac{r_s}{R_*} \quad \vec{\beta} = \frac{1}{c} \vec{\omega}_{LT} \times \vec{r}$$

$$\nabla \times \left( \alpha \vec{E} + \frac{\vec{\beta}}{c} \times \vec{B} \right) = -\frac{1}{c} \frac{\partial \vec{B}}{\partial t},$$

$$\nabla \times \left( \alpha \vec{B} - \frac{\vec{\beta}}{c} \times \vec{E} \right) = \frac{1}{c} \frac{\partial \vec{E}}{\partial t} + \alpha \vec{j} - \rho \vec{\beta}.$$

$$\frac{J_{\hat{r}}}{\rho G J C} \approx \left( \frac{J_{\hat{r}}}{\rho G J C} \right)_{\text{flat}} \frac{1}{1 - \omega_{LT} / \Omega_*}$$



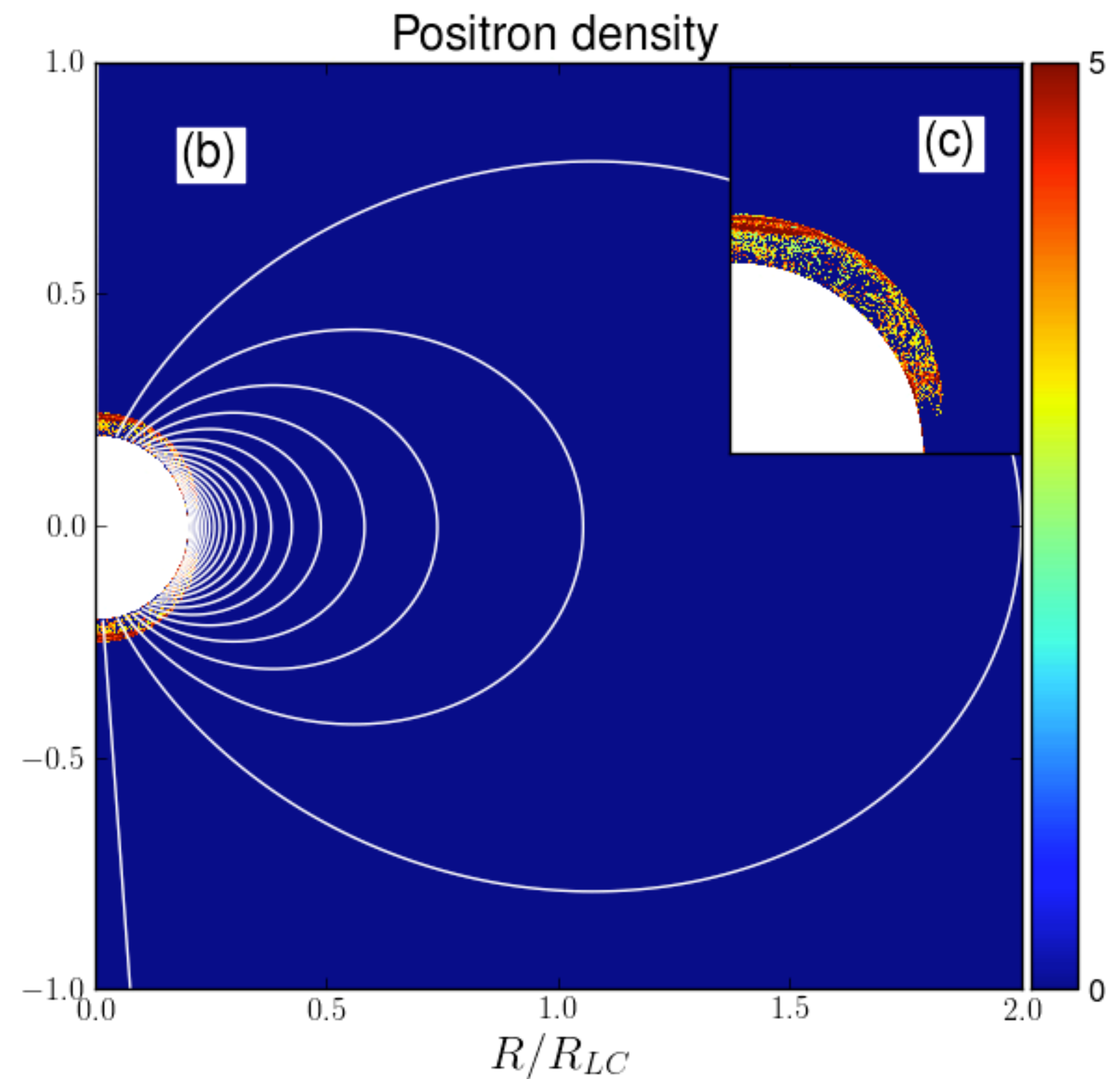
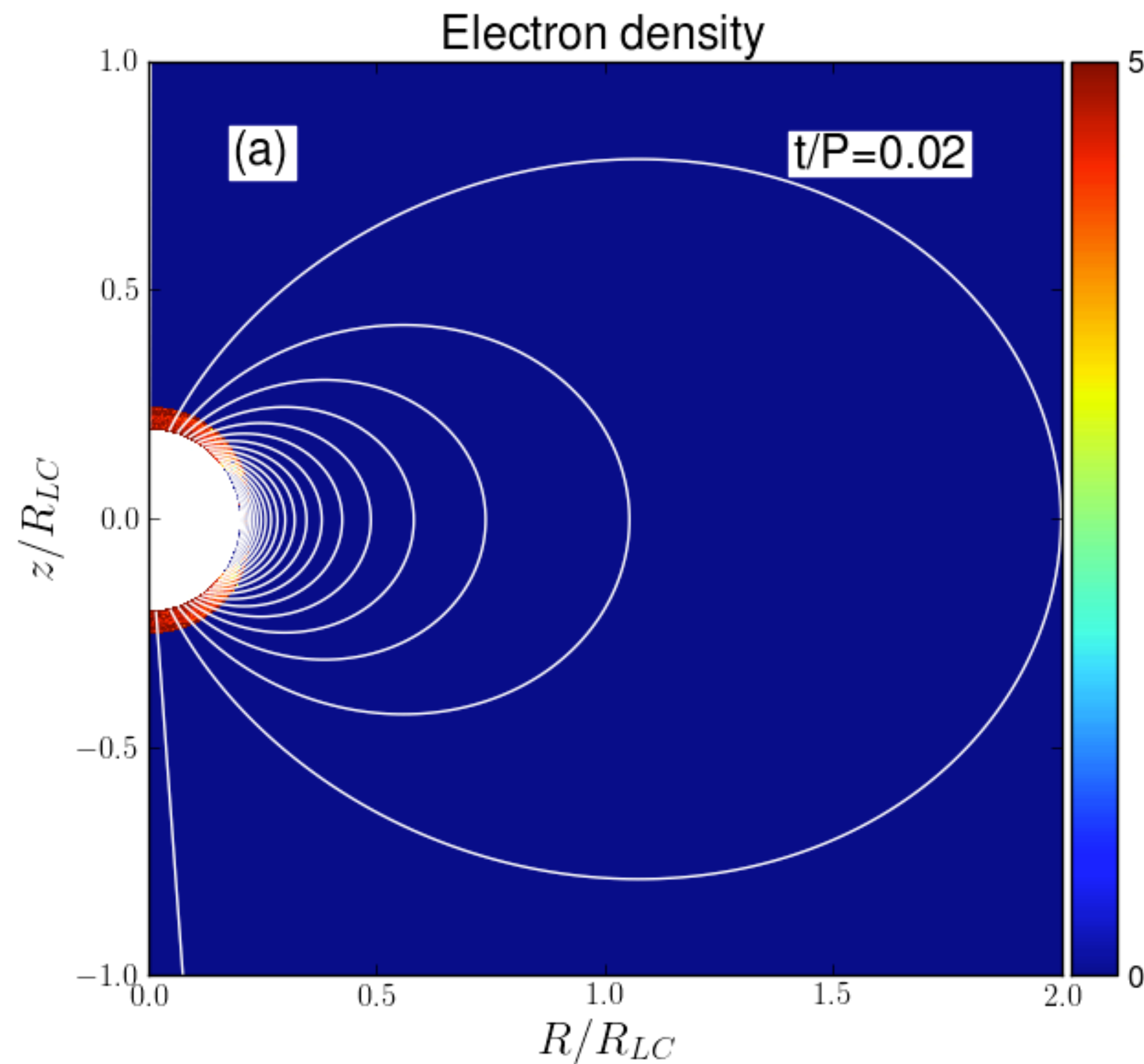
Lense-Thirring frame dragging

Frame-dragging makes effective rotation frequency of the star smaller close to the star (this lowers the necessary corotation charge), but the rotation is still the same far from the star (this keeps the current the same).

Philippov et al. (2015b)  
Beskin 1990  
Muslimov & Tsygan 1992



# ALIGNED ROTATOR WITH GR AND PAIRS

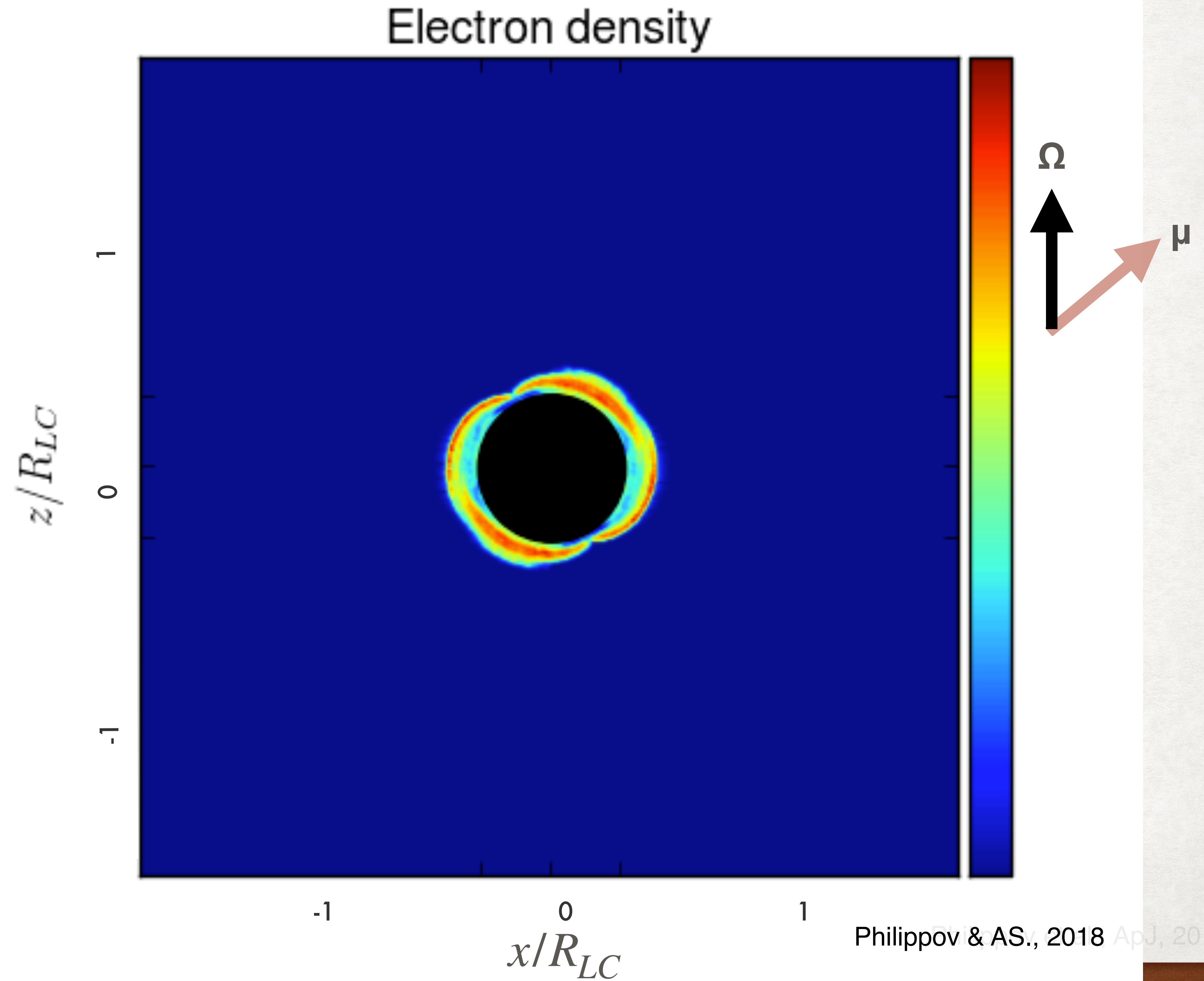


- Polar pair production returns!



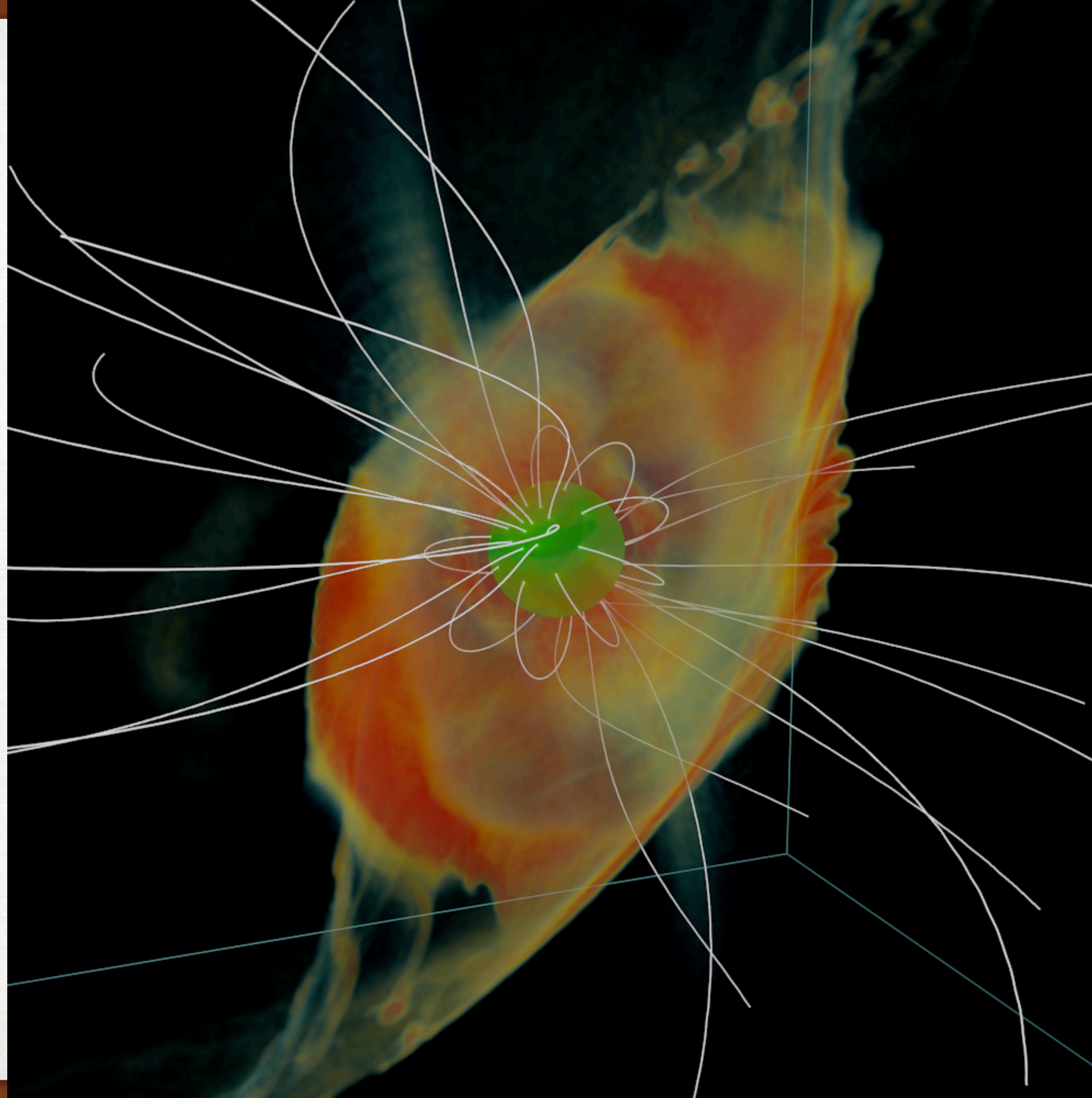
# OBLIQUE ROTATOR WITH GR AND PAIRS

- Pair production happens on the polar cap, in return current layers and in the current sheet beyond LC
- Polar discharge is non-stationary. Electric field screening by advecting plasma clouds generates waves. The plasma motions are collective and coherent — implications for radio emission (see Beloborodov 2008, Timokhin & Arons 2013)





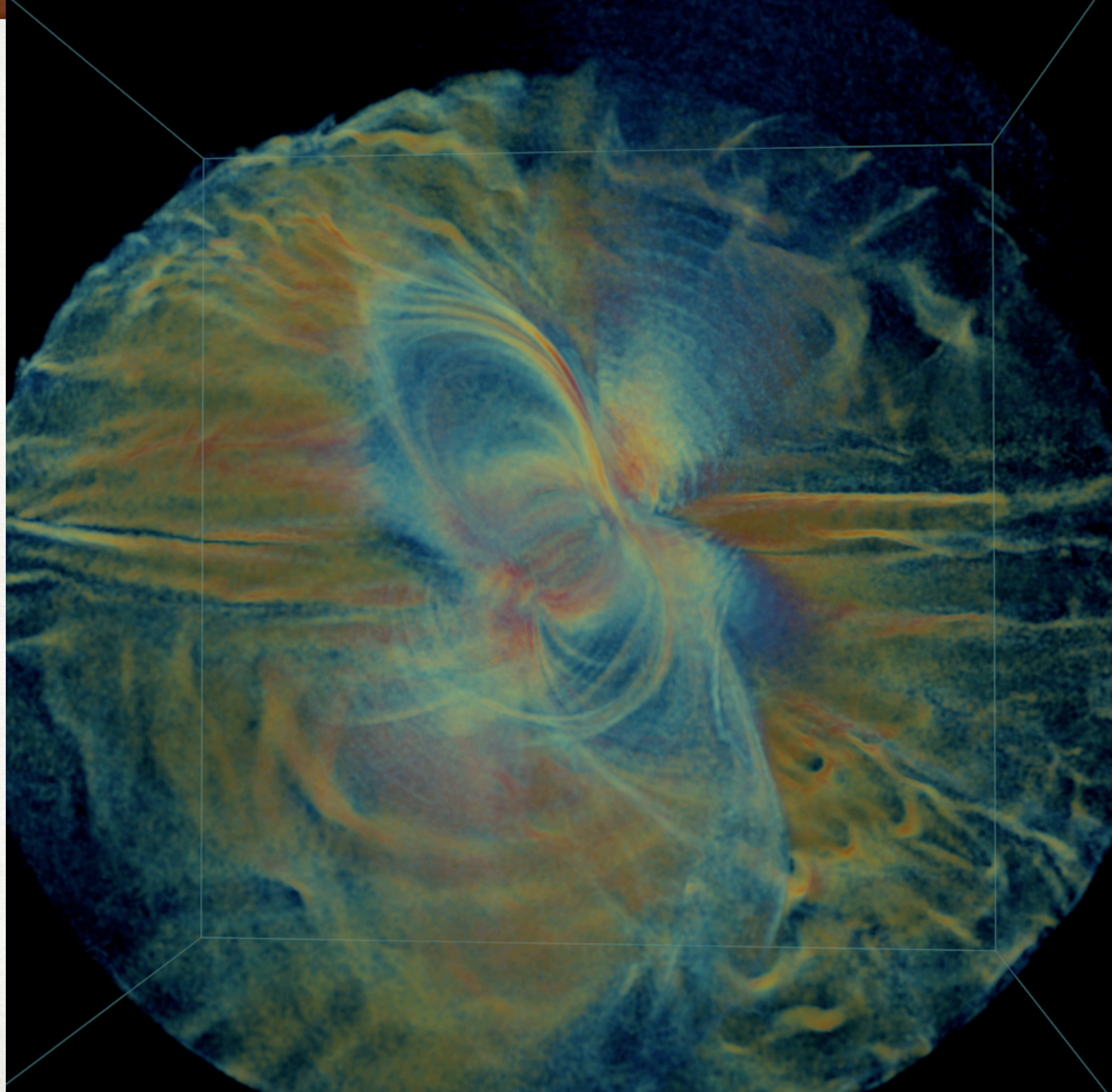
OBLIQUE  
ROTATOR  
WITH GR  
AND PAIRS:  
PLASMA  
DENSITY



ApJ, 2016



OBLIQUE  
ROTATOR  
WITH GR  
AND PAIRS:  
CURRENT  
DENSITY

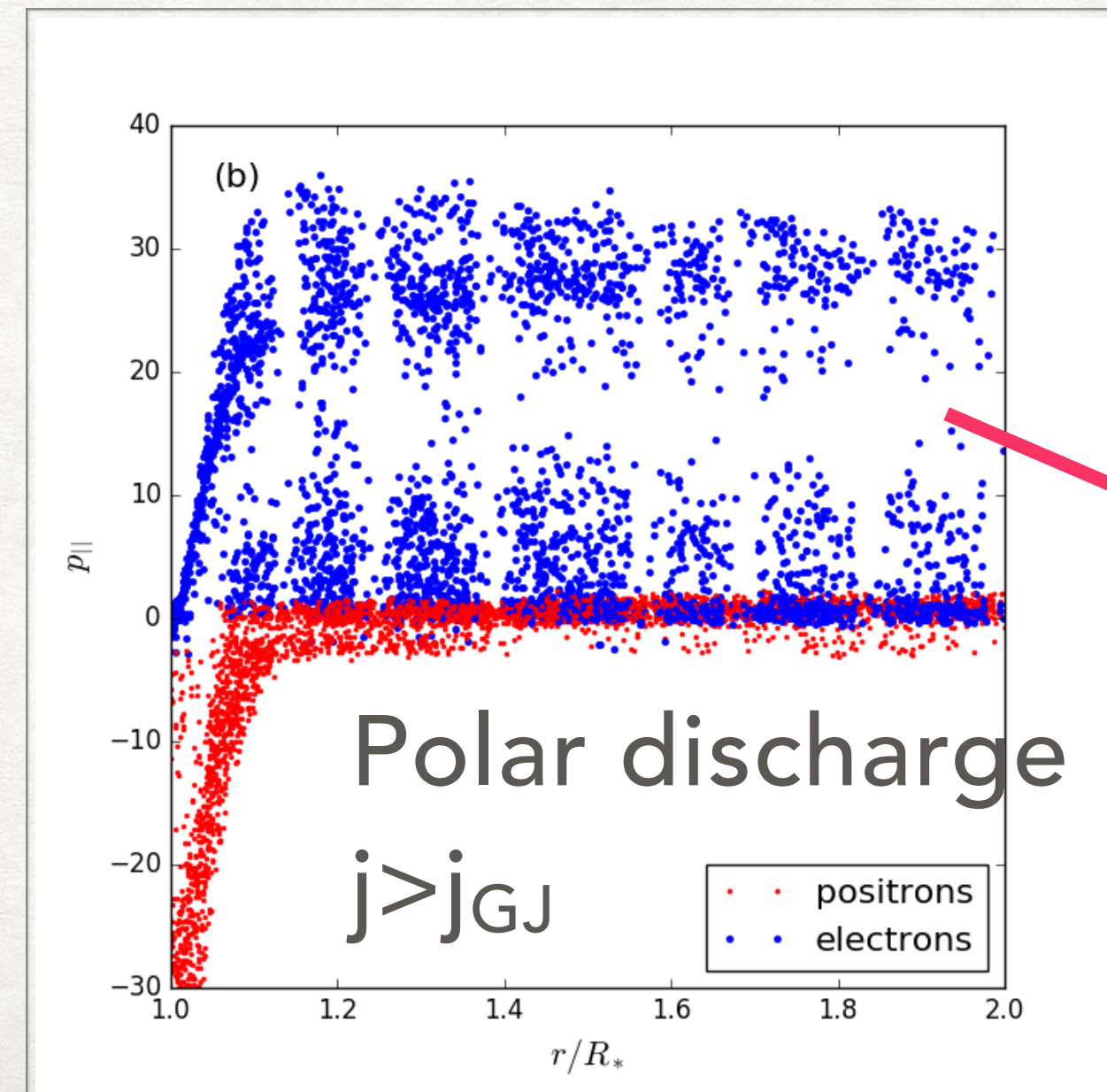




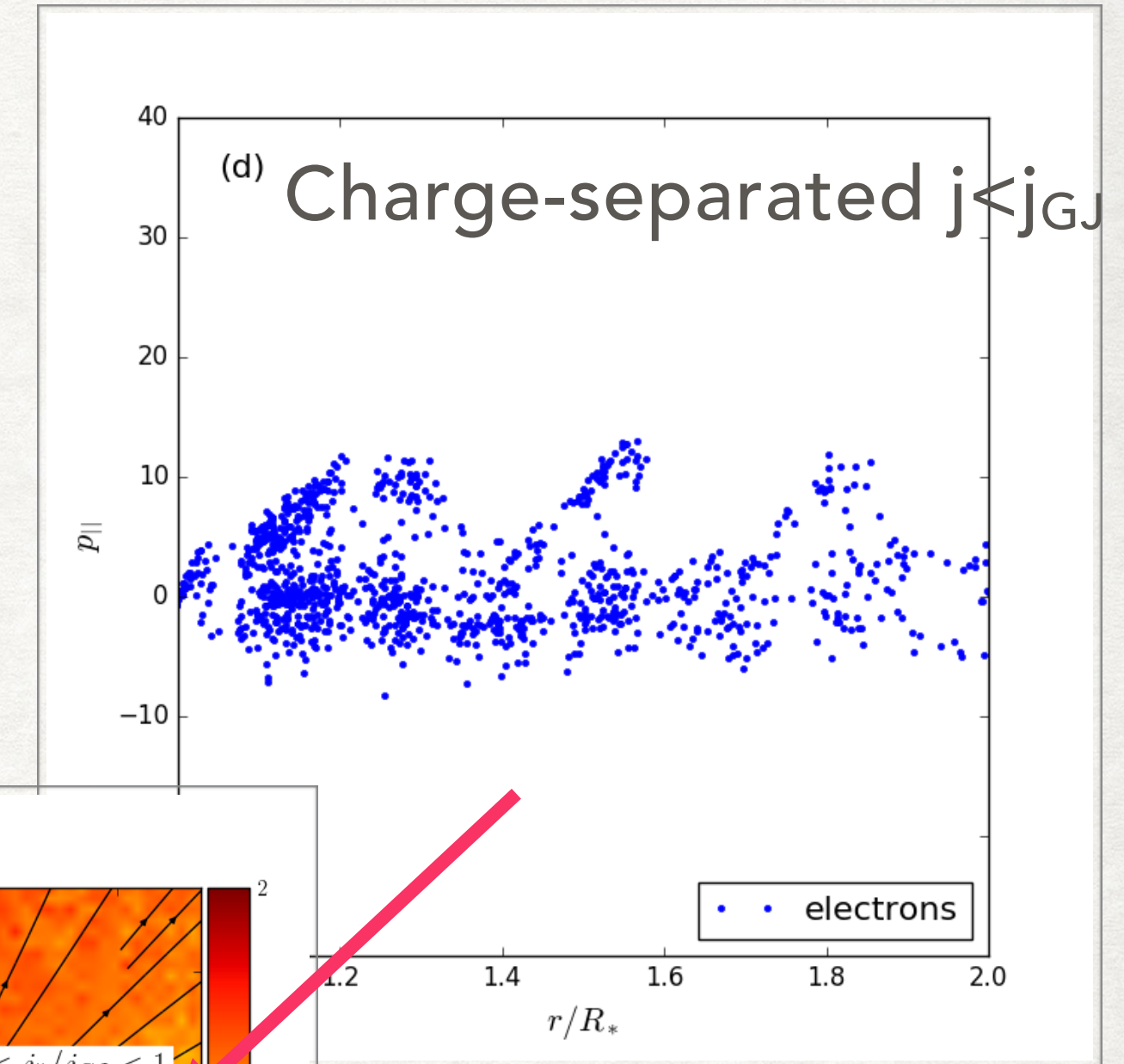
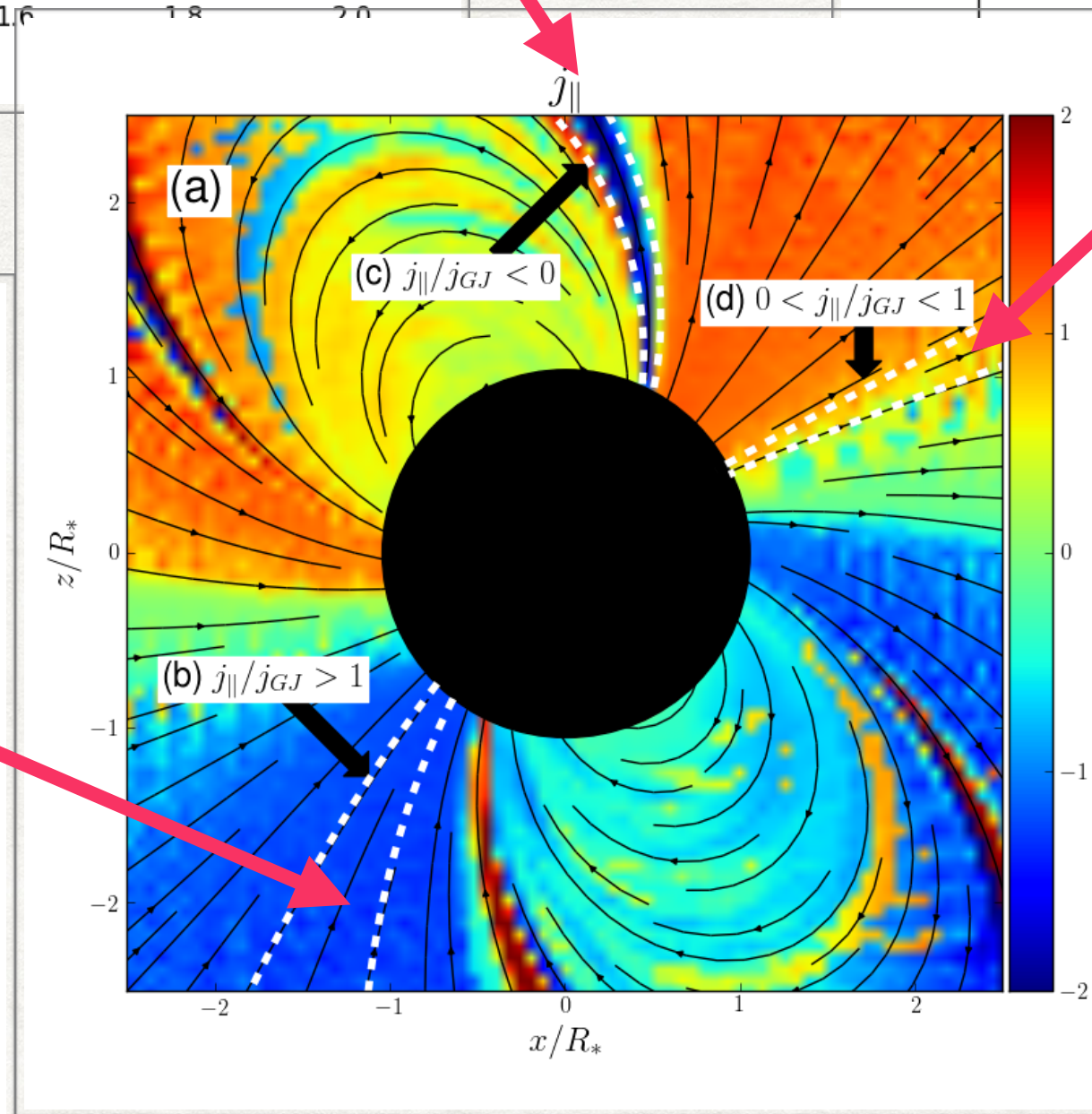
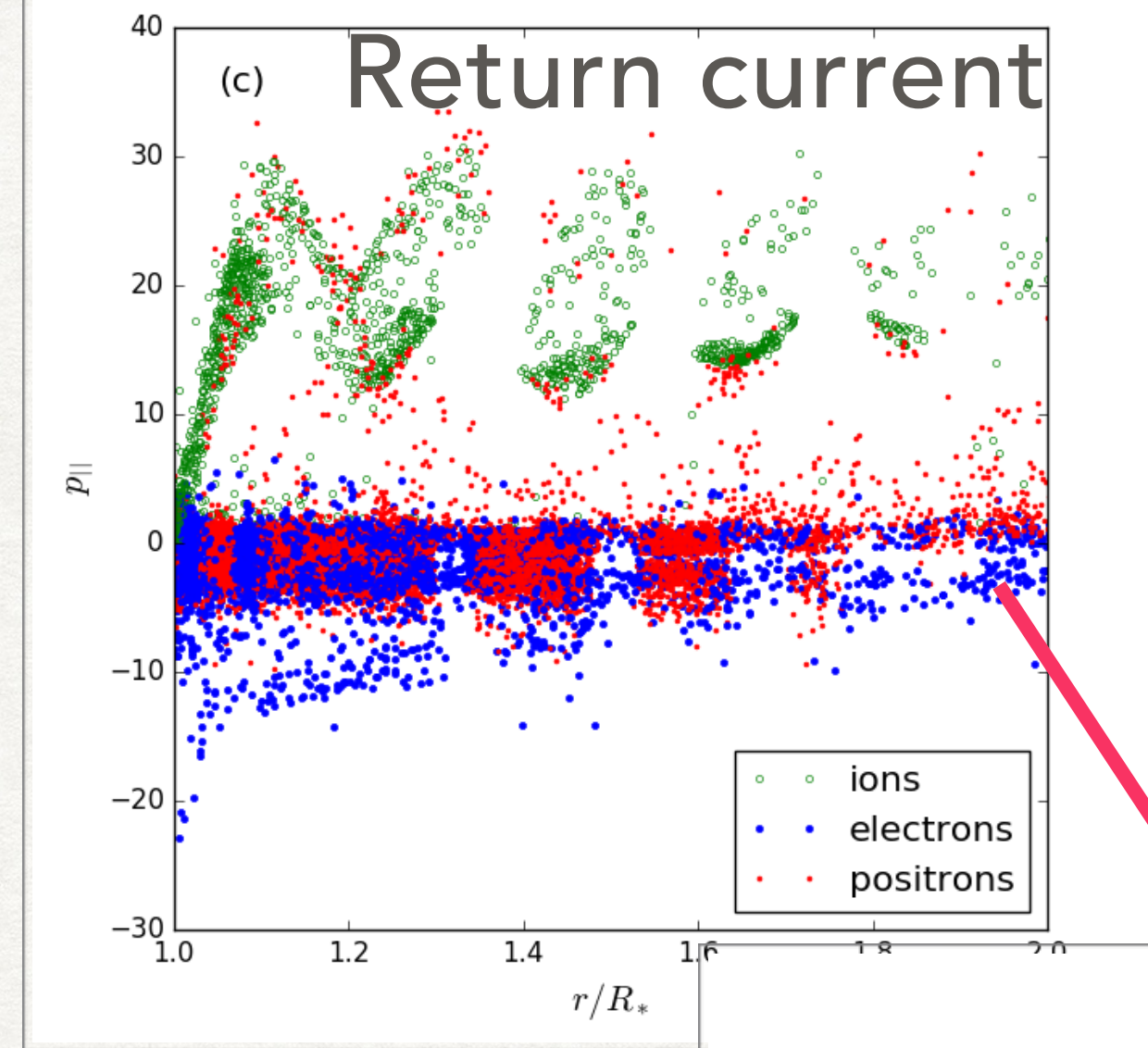
# OBLIQUE ROTATOR WITH GR AND PAIRS

- Counterstreaming is present in polar discharge and in return current
- Opportunities for maser emission from collective instabilities of counterstreaming distributions.

Momentum space



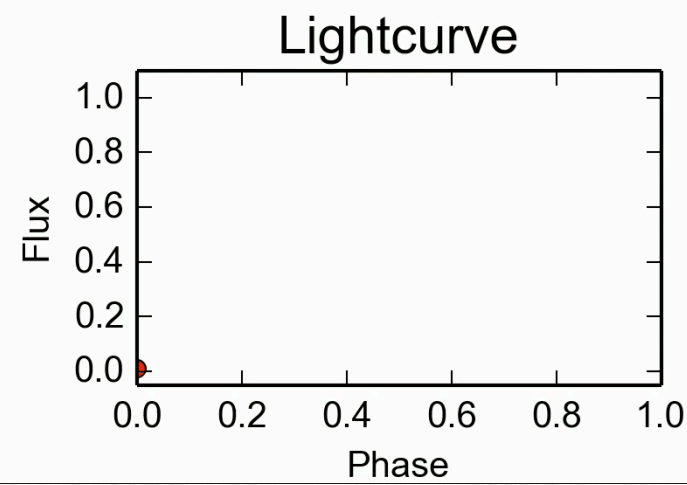
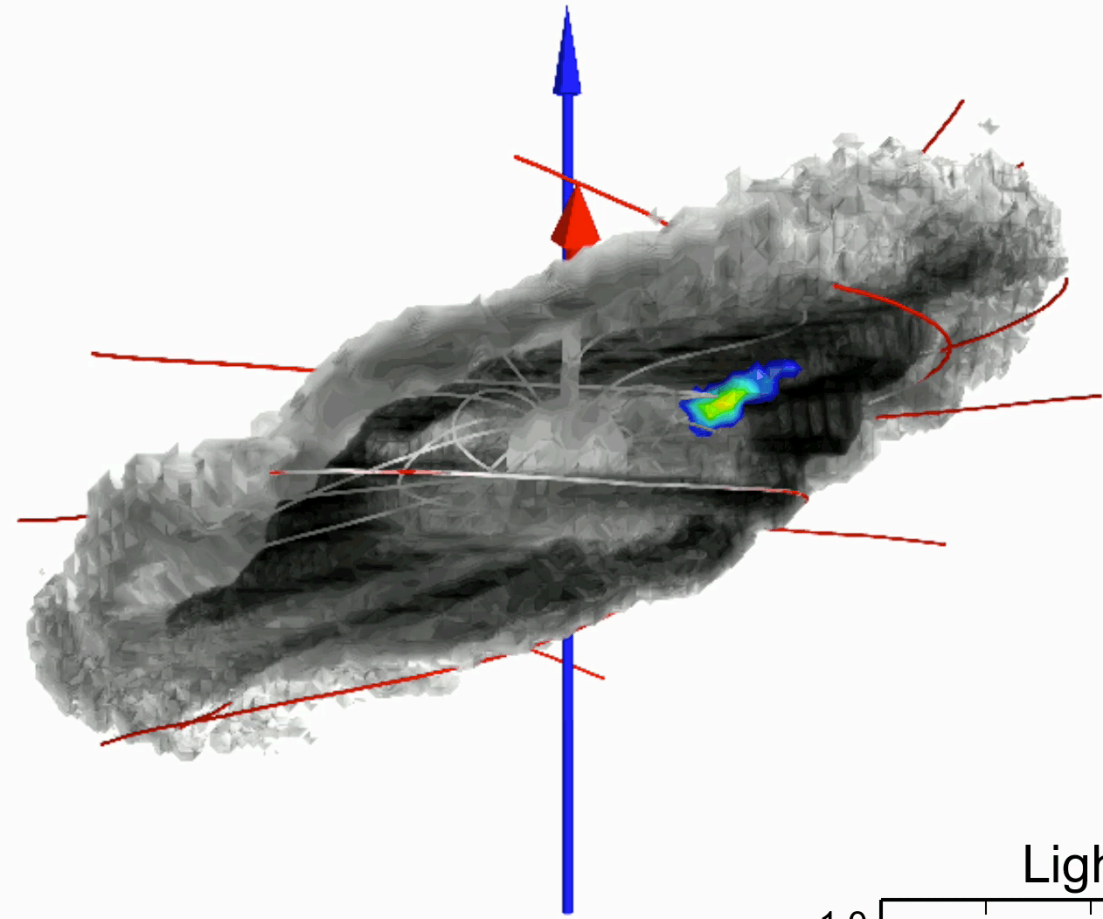
Accelerates ions!



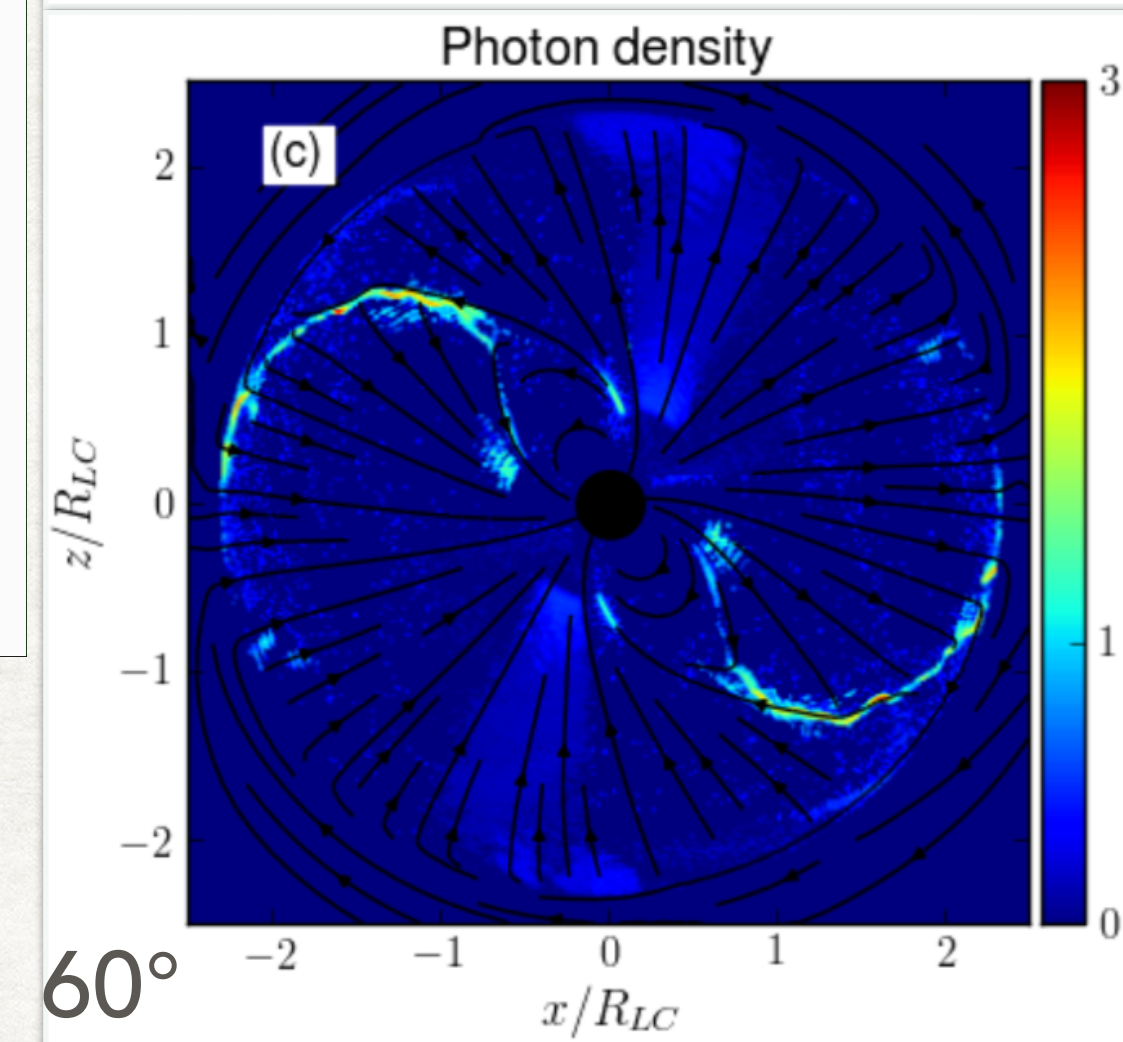
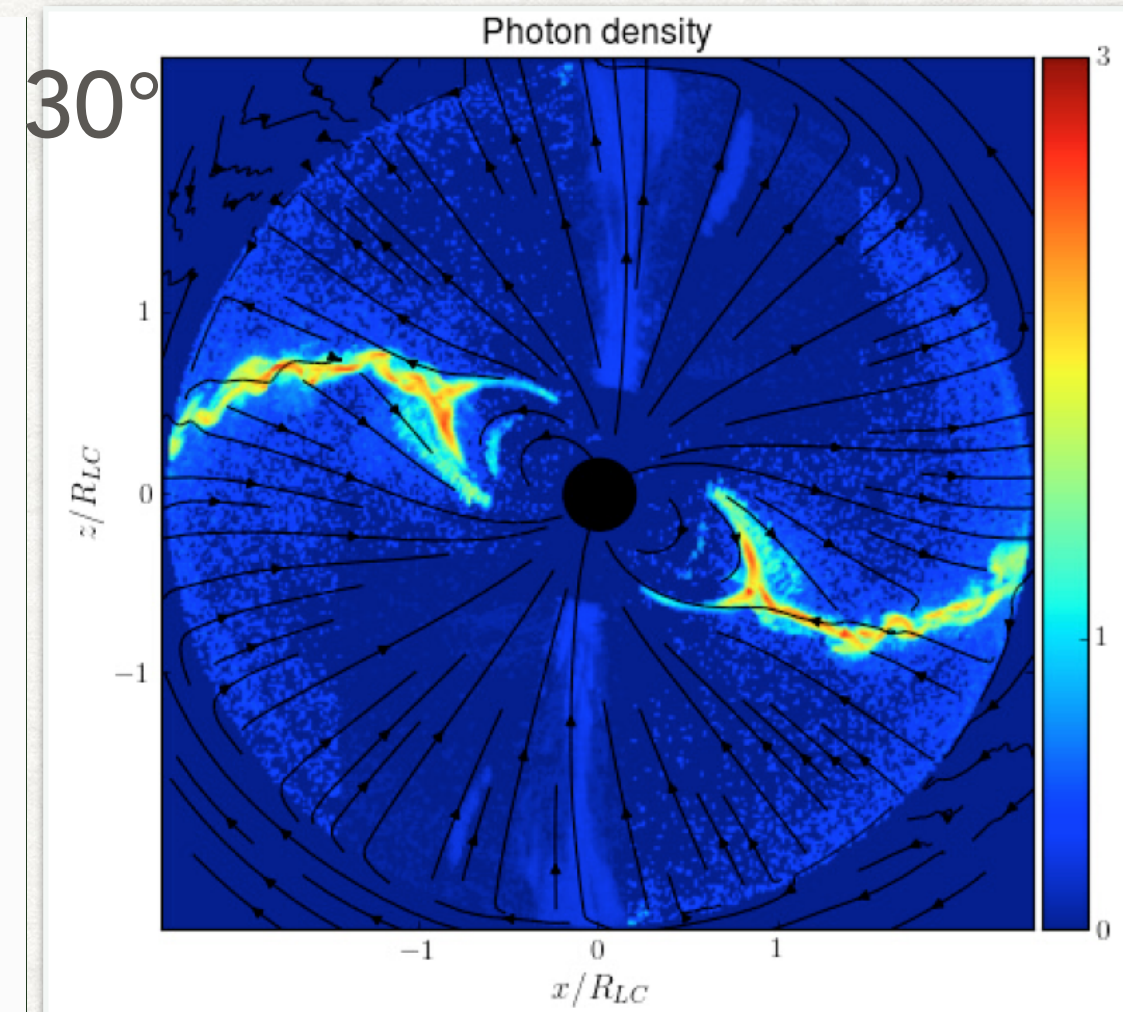


# GAMMA-RAY LIGHTCURVES

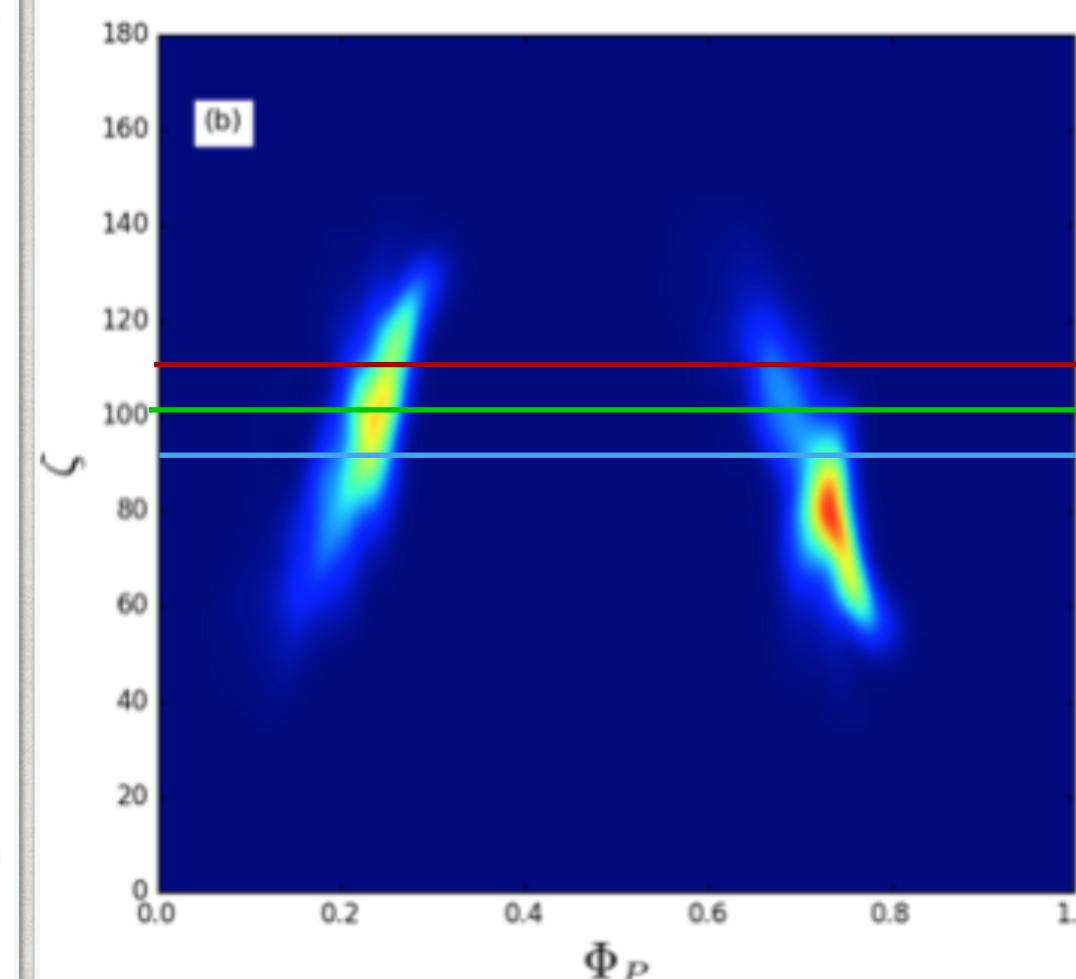
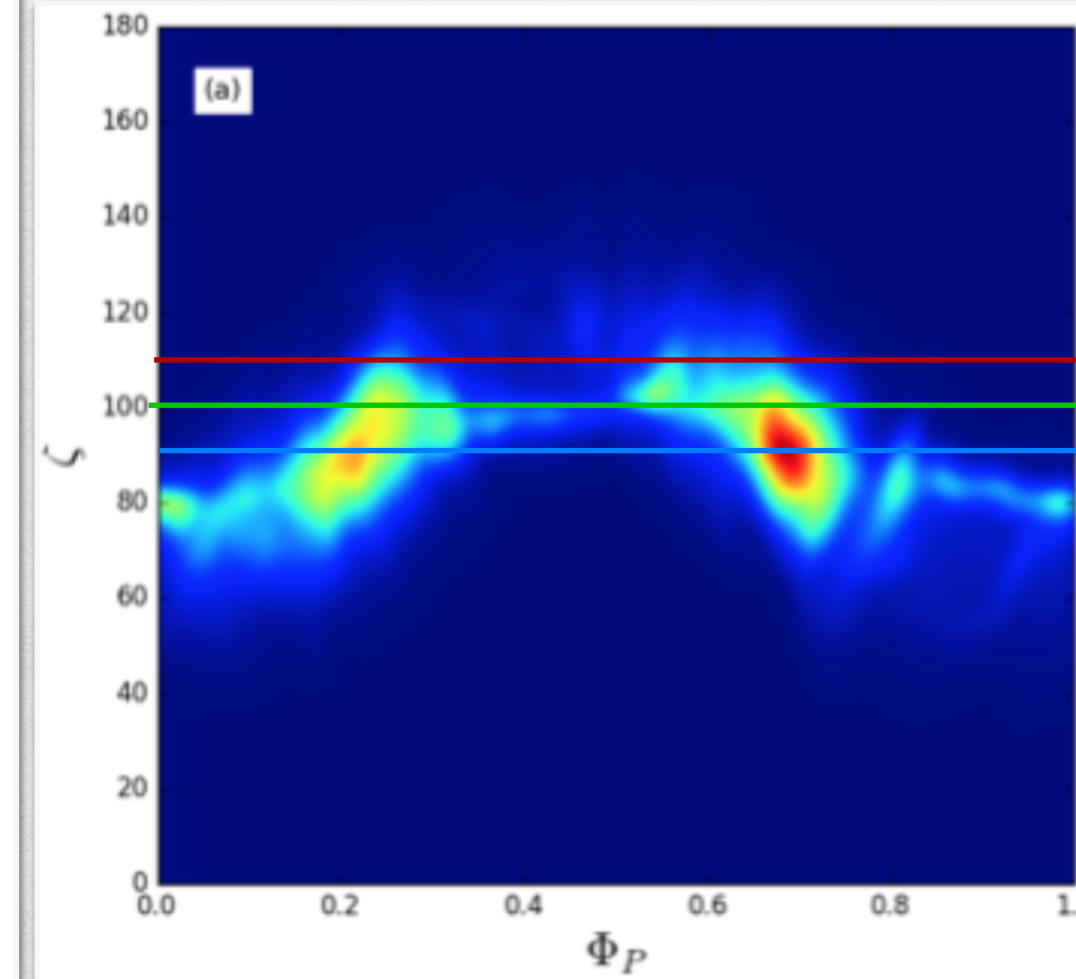
$i=30$  - Phase=0.00 - Positrons -



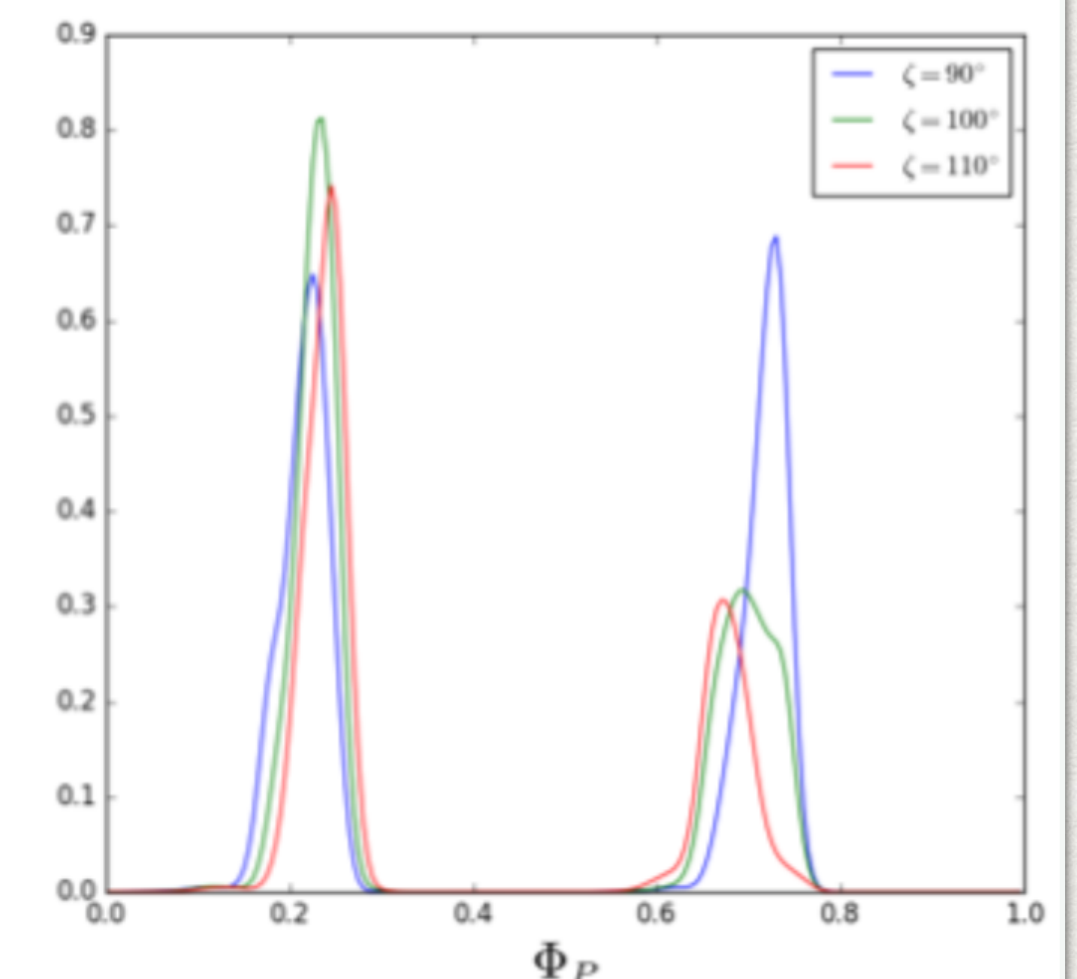
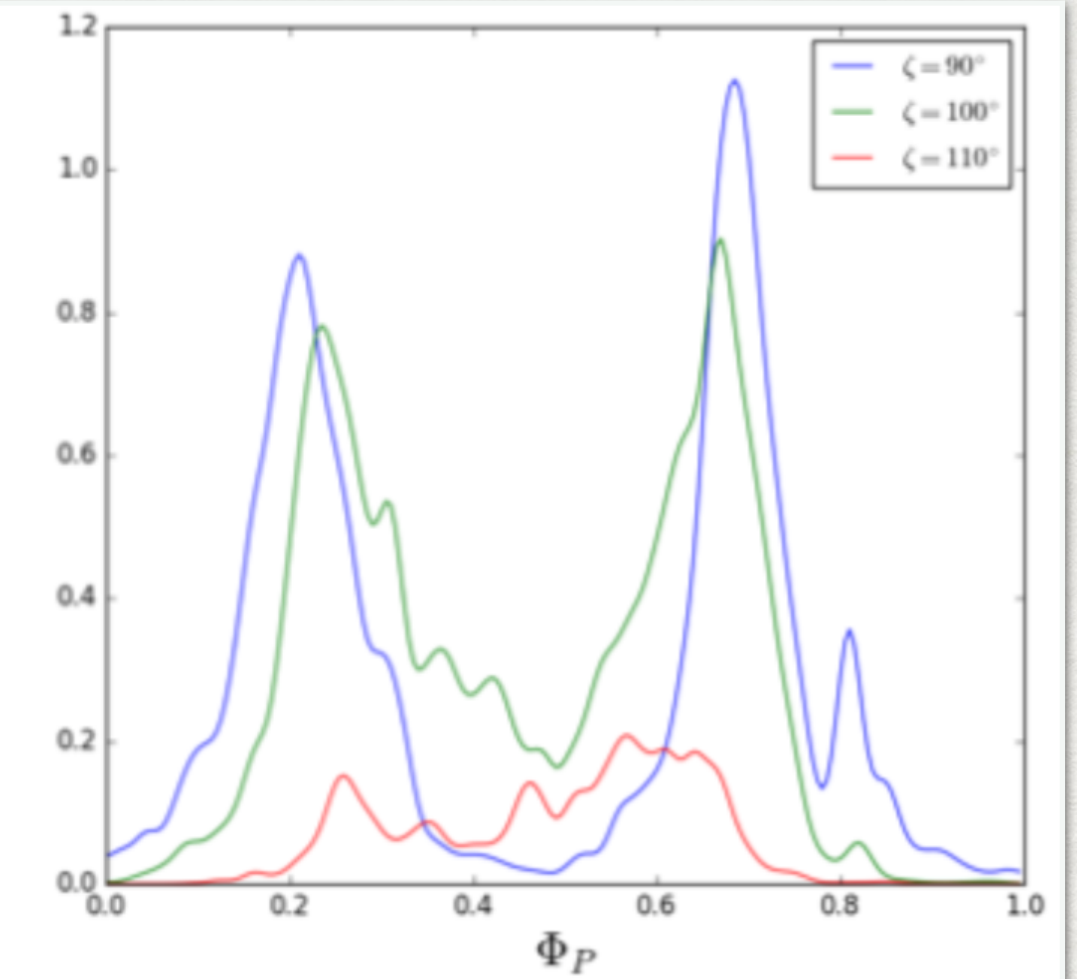
Caustic formation from spiral shape of current sheet



Photon density



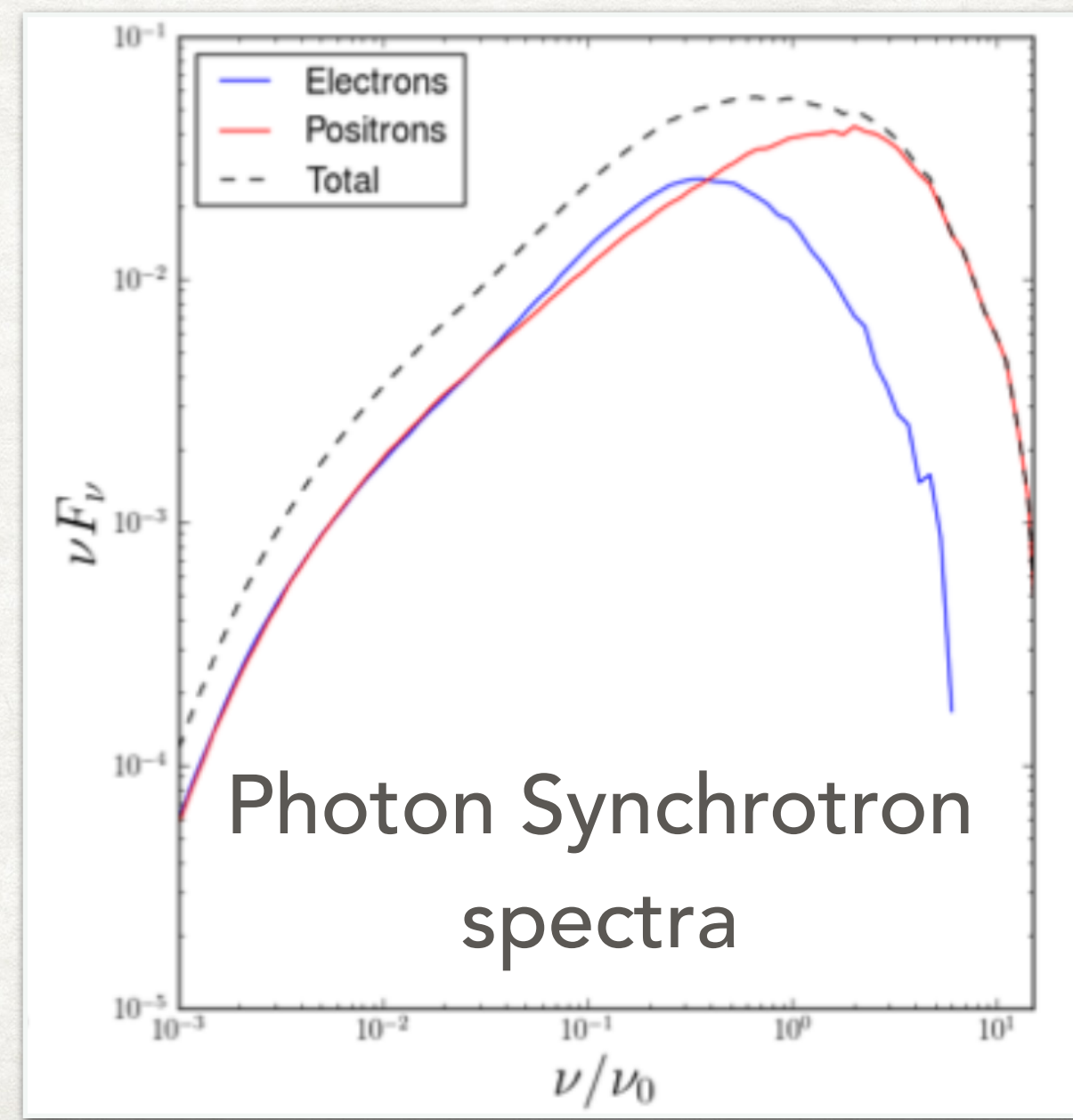
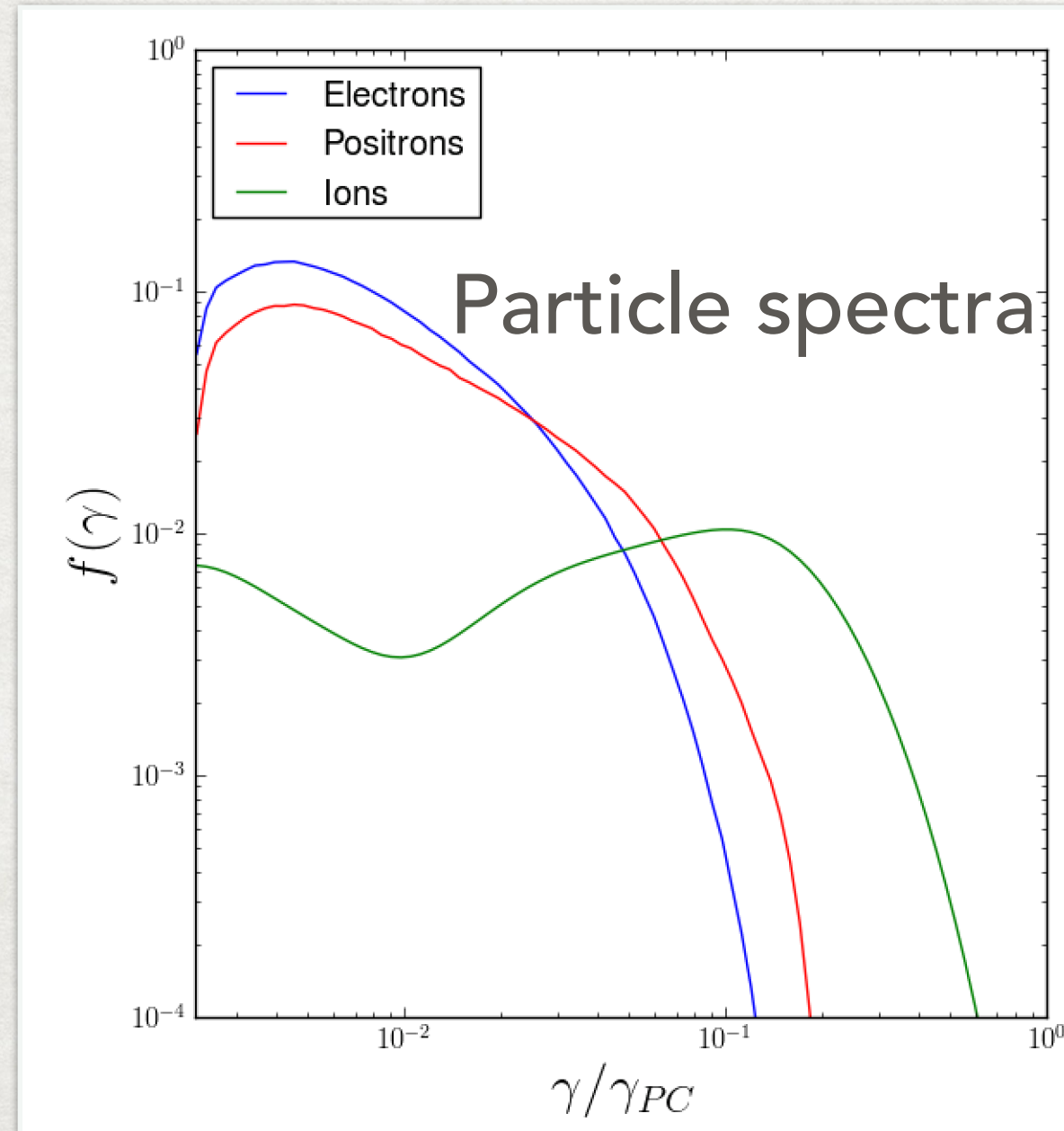
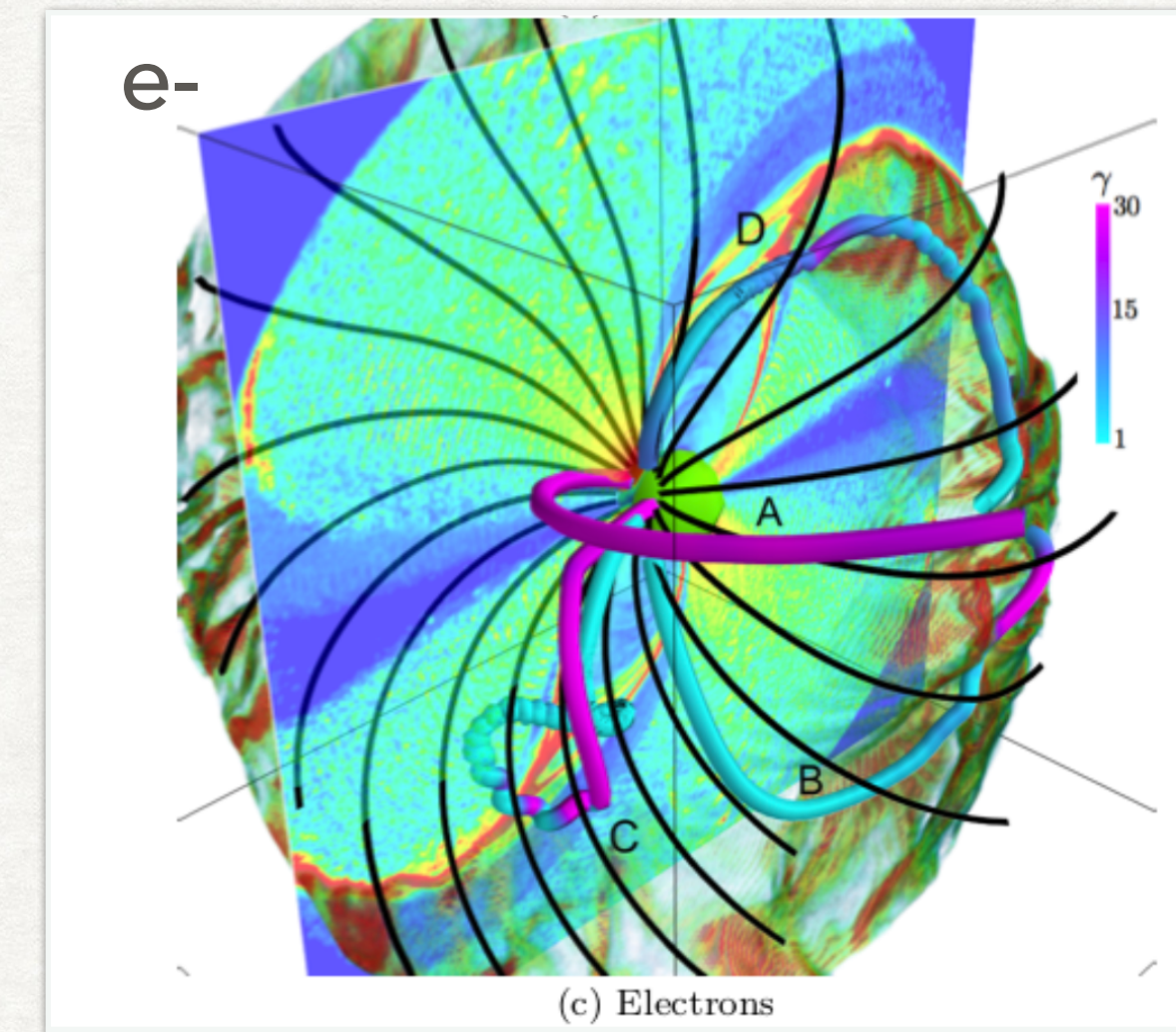
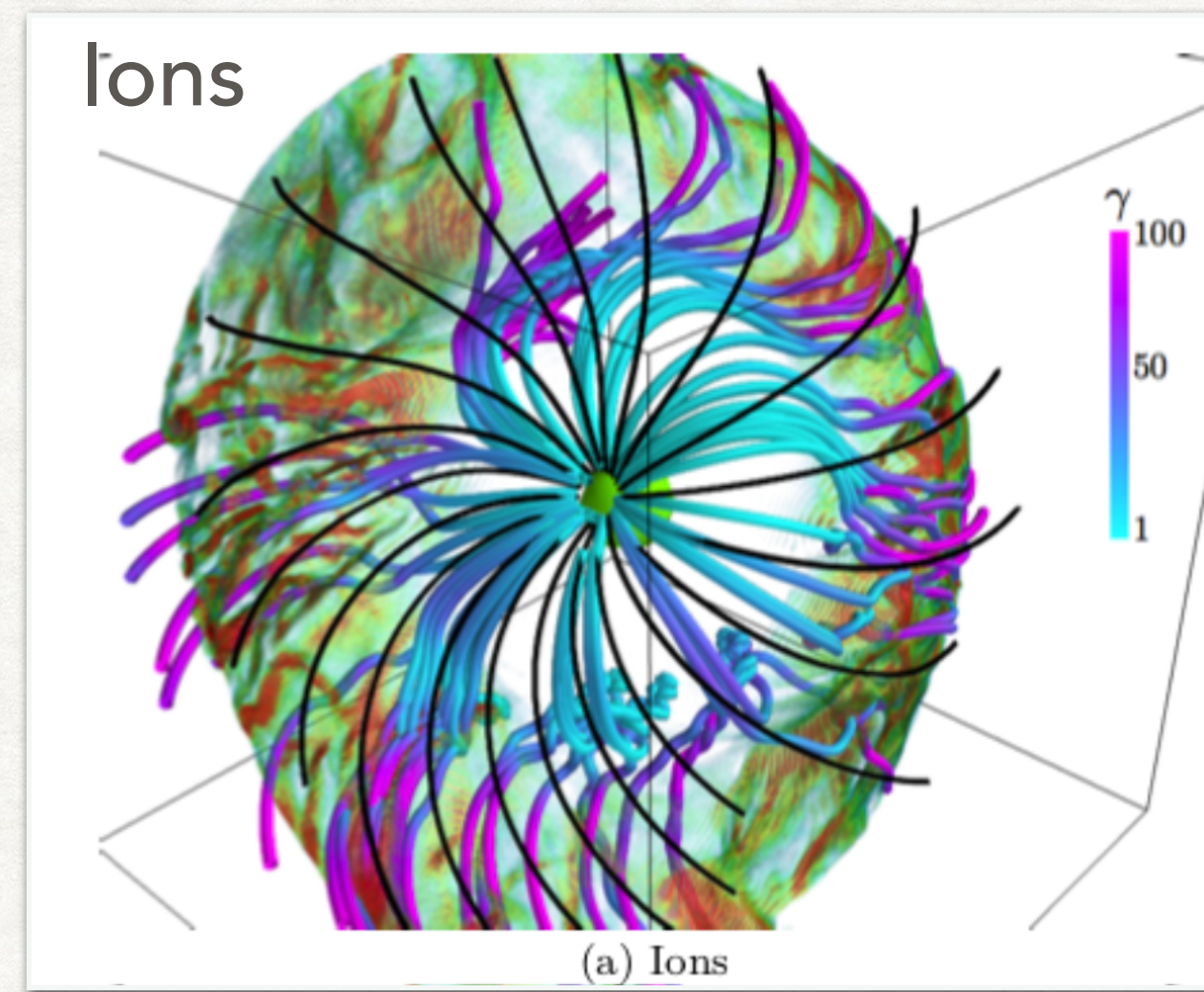
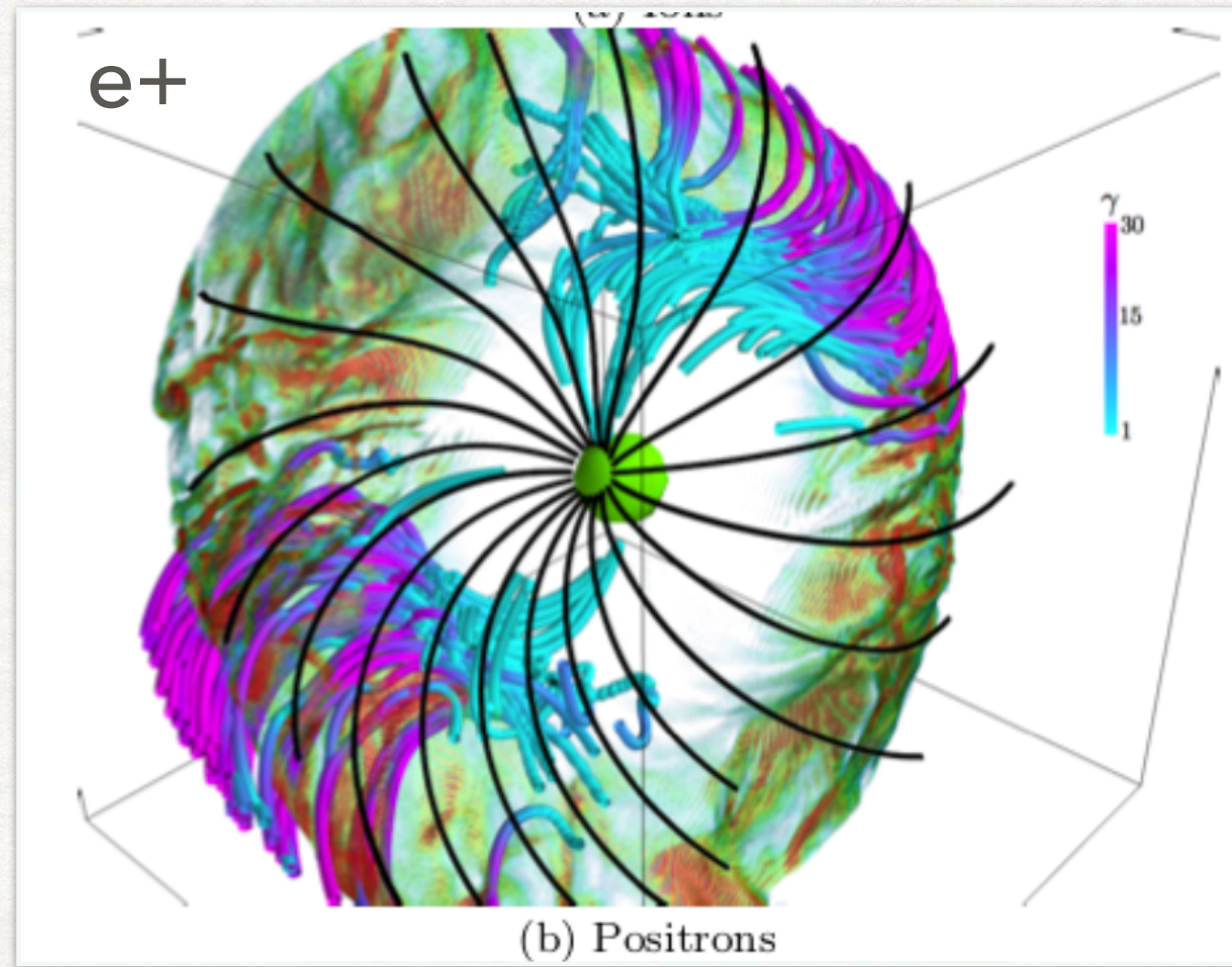
Sky map



Light Curve



# PARTICLE ACCELERATION AND SPECTRA



Particles are accelerated in the current sheet during reconnection. Radiation appears as broad spectral peak. The max frequency is set by magnetization

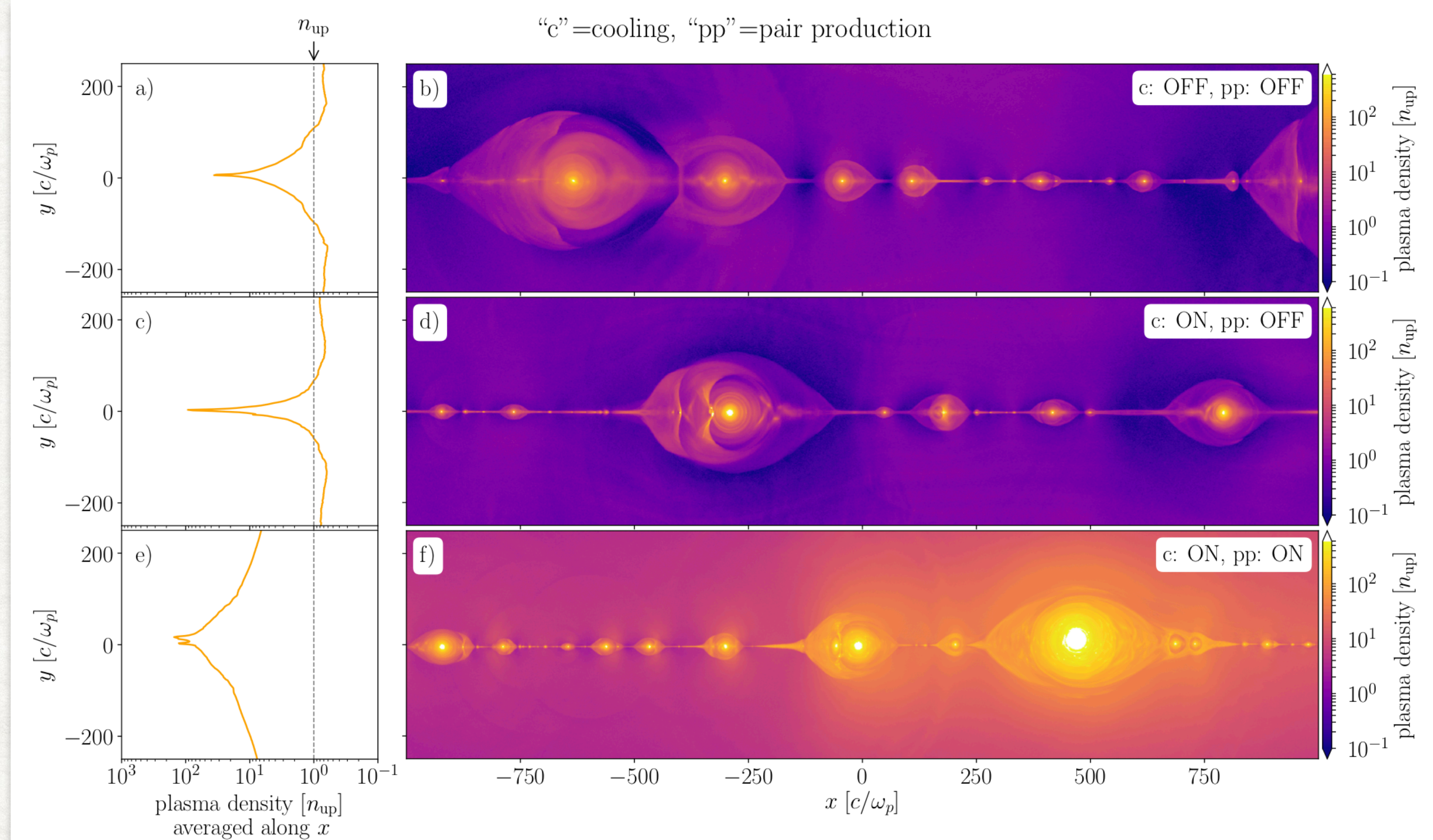
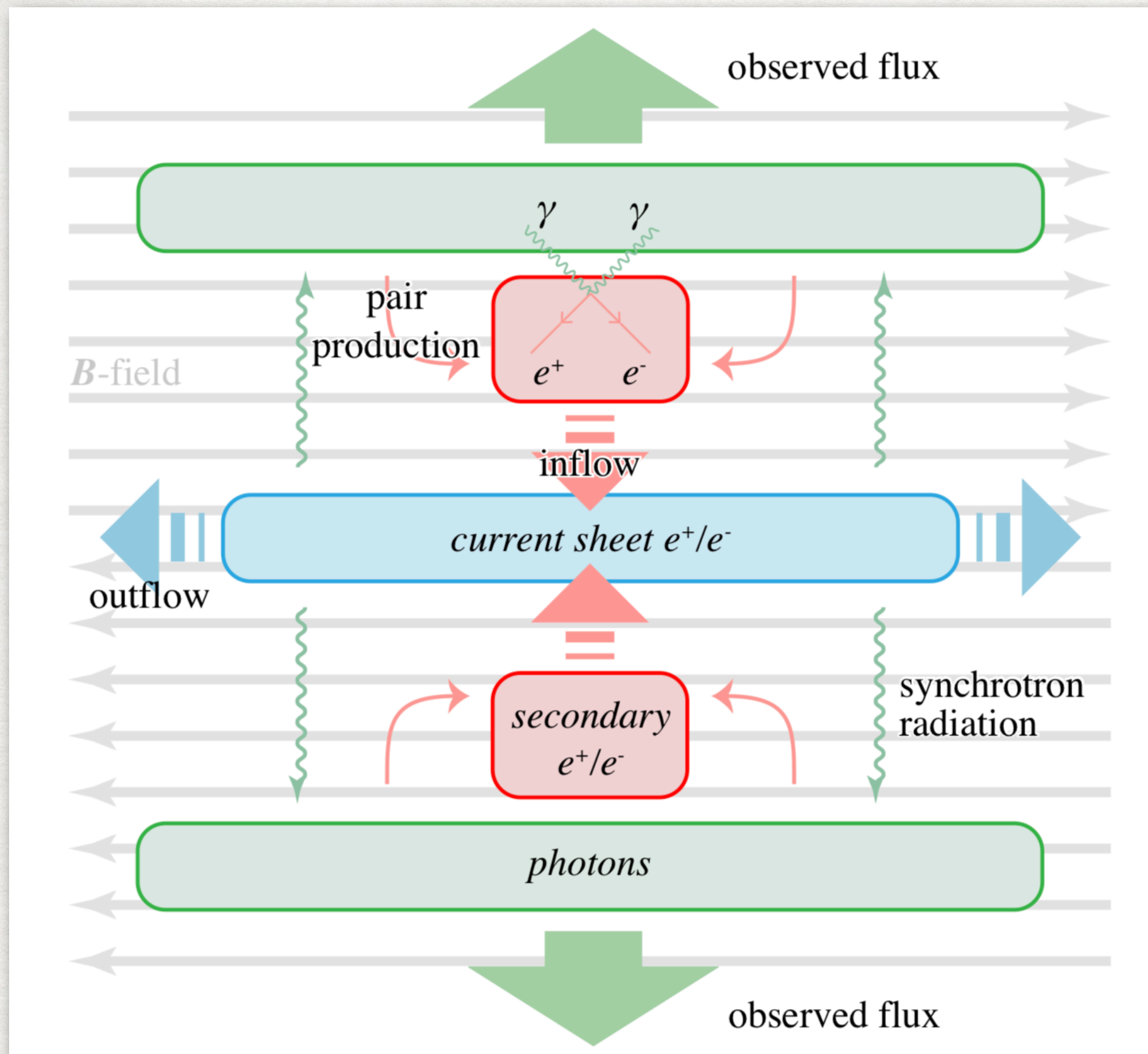
$\sigma_{LC}$ :

$$\nu_{\max} \approx 3e (0.1 B_{LC}) \sigma_{LC}^2 / 4\pi m_e c$$

Pair production in sheet sets the sigma parameter. Ions gain good fraction of  $\Phi_{pc}$



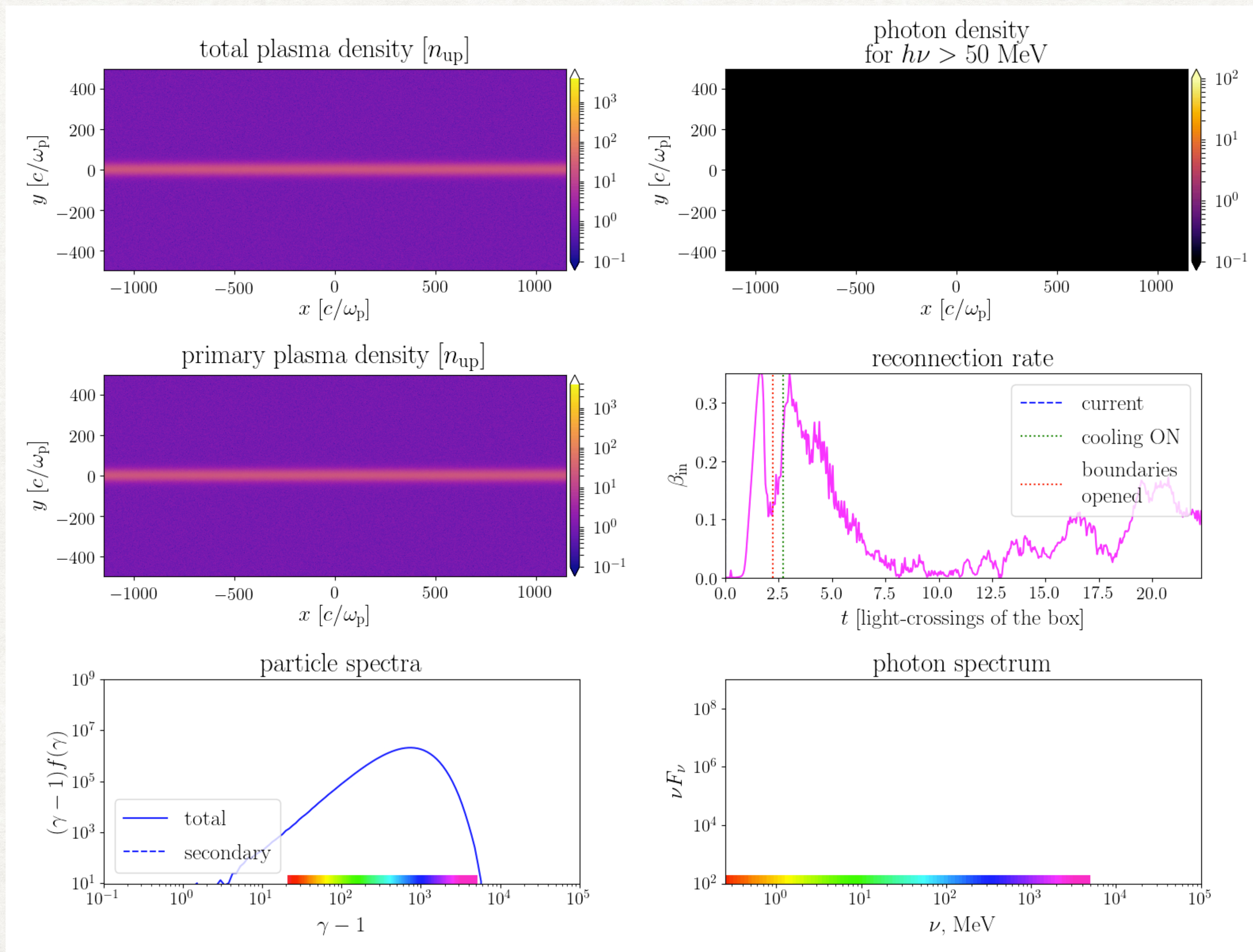
# THE ROLE OF RECONNECTION WITH PAIR PRODUCTION IN SETTING CUTOFF ENERGY



Reconnection in the current sheet is main particle accelerator. Gamma-gamma pair formation can start. Pair formation increases the pair loading above the sheet, and lowers effective magnetization in the sheet. Particle acceleration follows magnetization, max particle energy is reduced.

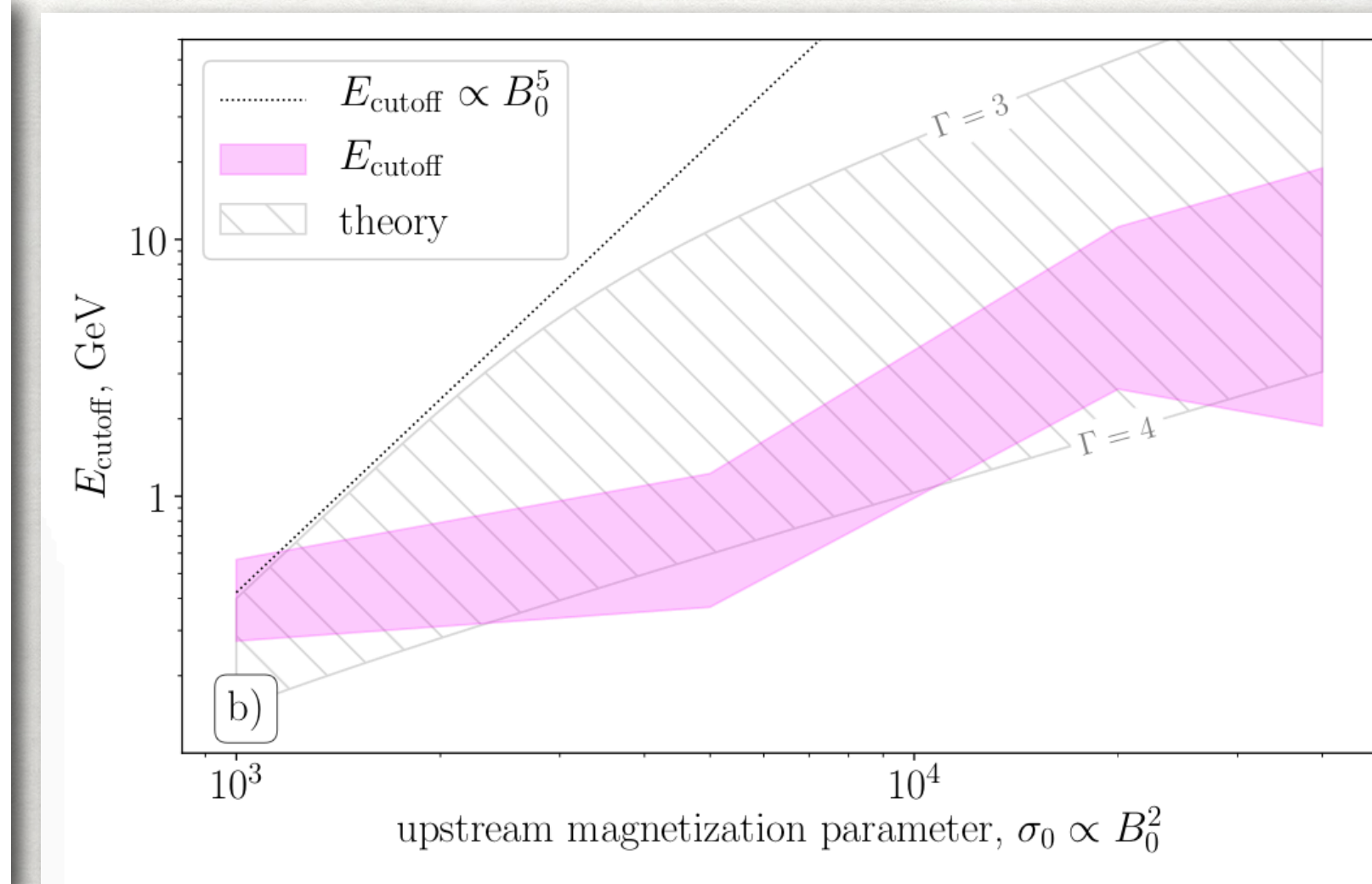
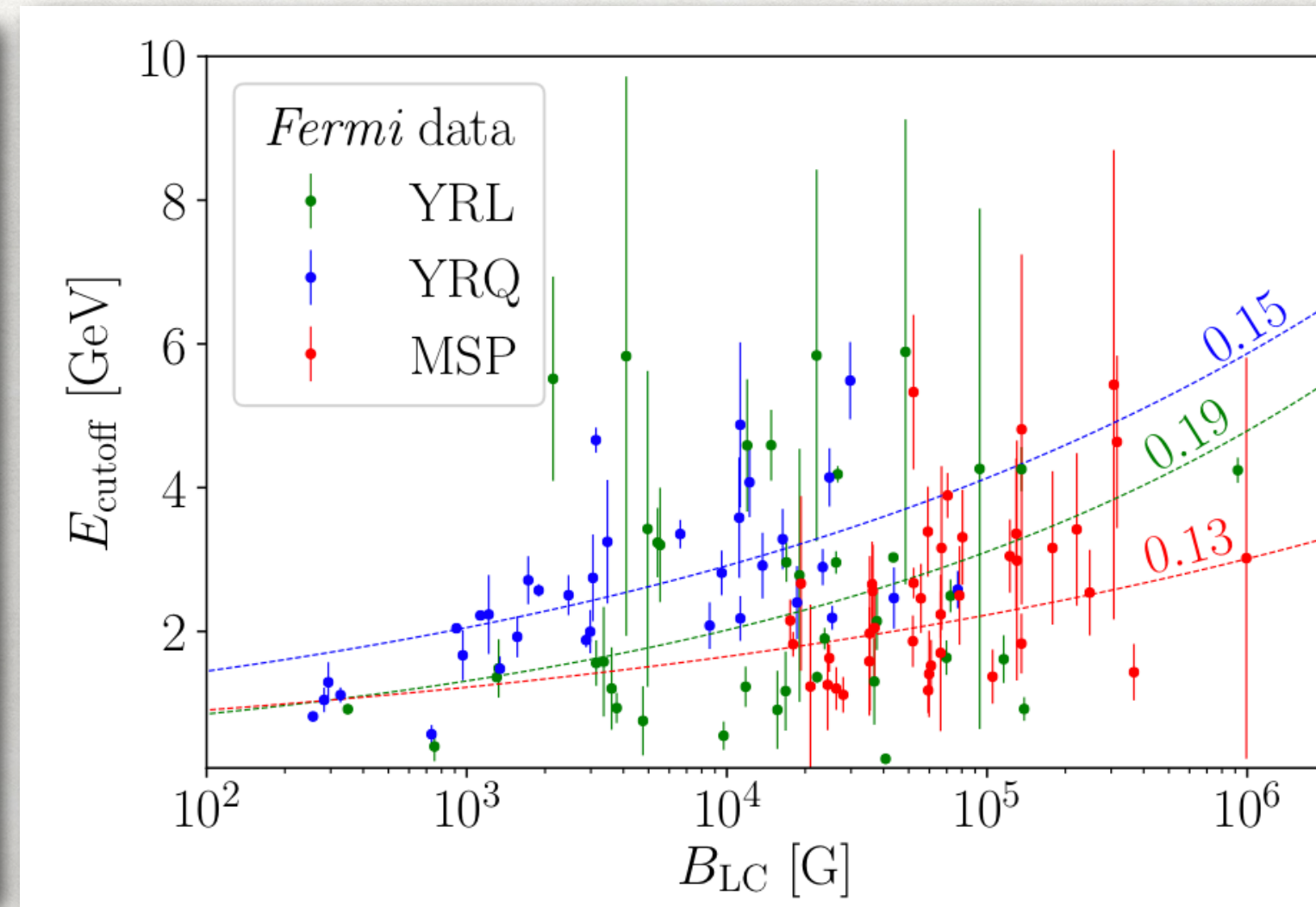
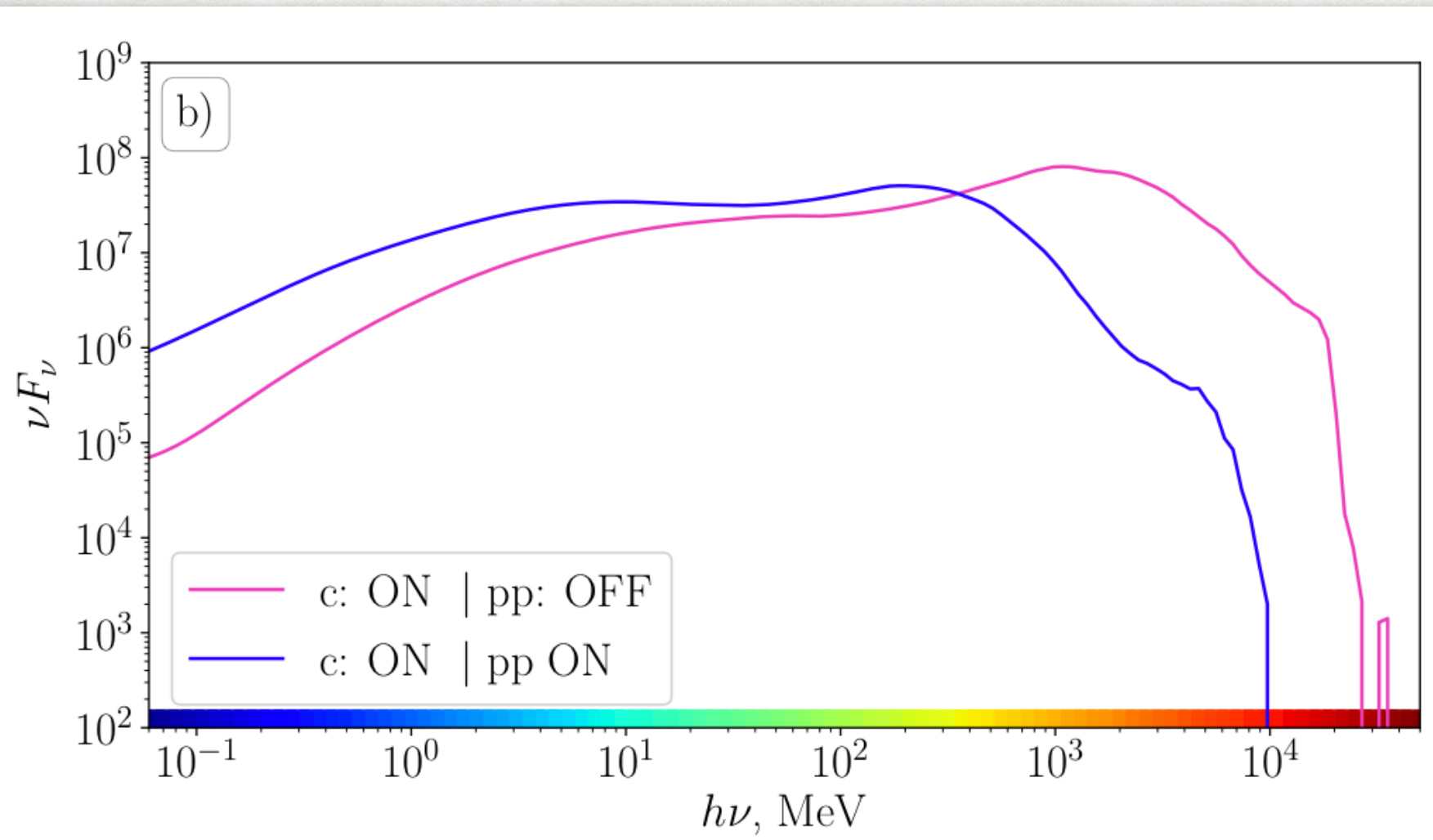


# THE ROLE OF RECONNECTION WITH PAIR PRODUCTION IN SETTING CUTOFF ENERGY





# THE ROLE OF RECONNECTION WITH PAIR PRODUCTION IN SETTING CUTOFF ENERGY



Pair formation increases the pair loading above the sheet, and lowers effective magnetization in the sheet. Particle acceleration follows magnetization, max particle energy is reduced. Naively, cutoff energy should be a strong function of B at the LC.

$$\gamma_{\text{cutoff}} \propto \sigma_0 \propto B_0^2, \quad E_{\text{cutoff}} \propto \gamma_{\text{cutoff}}^2 B_0 \propto B_0^5$$

Pair loading softens the dependence

$$\gamma_{\text{cutoff}} \sim \sigma_{\text{LC}} \propto B_{\text{LC}}^2 / \eta n_{\text{GJ}}$$

Expect cutoff energy dependence to be between  $E_{\text{cutoff}} \propto B_{\text{LC}}^{1.2} - B_{\text{LC}}^{1.8}$  and  $E_{\text{cutoff}} \propto B_{\text{LC}}^{-0.8} - B_{\text{LC}}^{-0.2}$

Observed dependence:

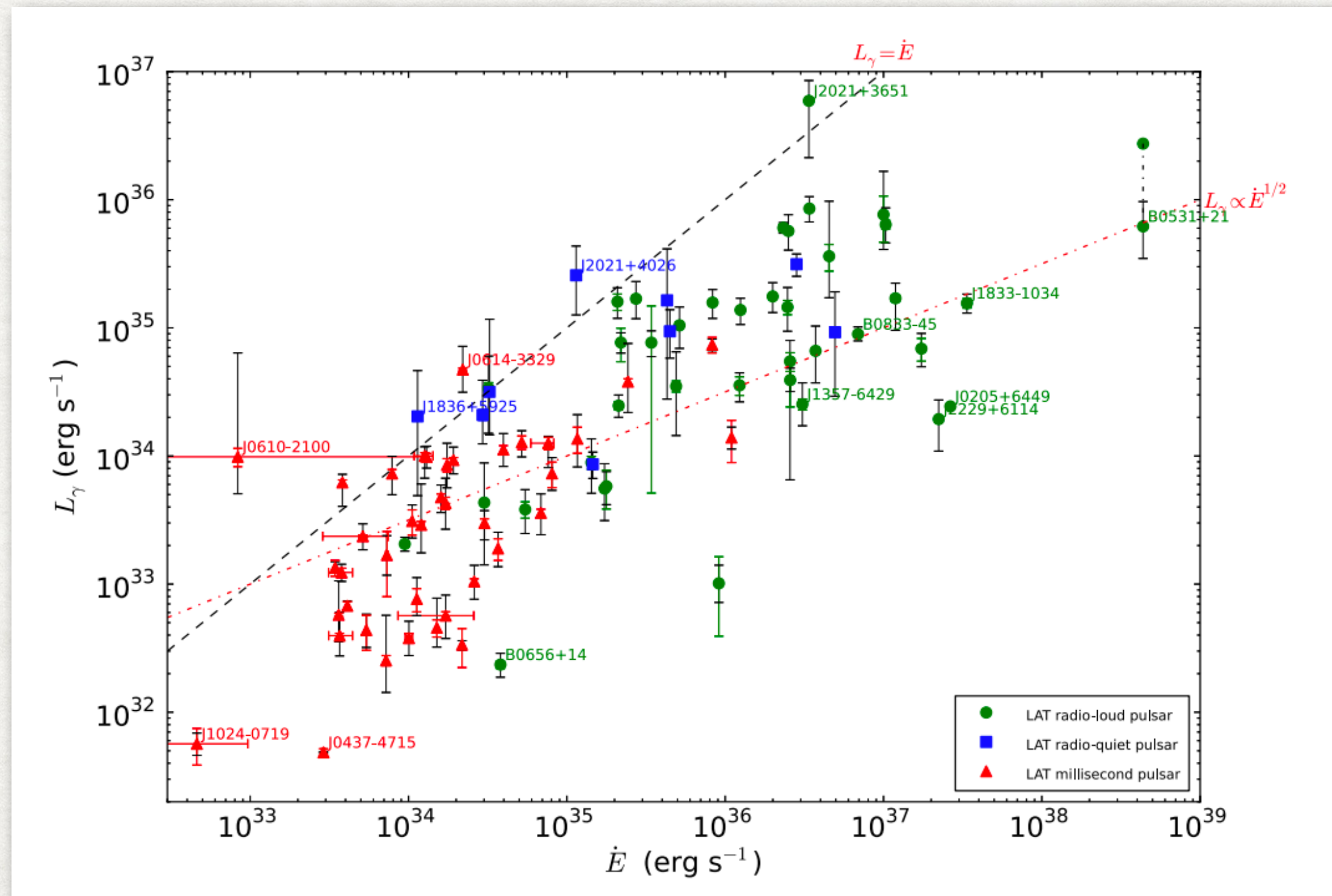
$$E_{\text{cutoff}} \propto B_{\text{LC}}^{0.1} - B_{\text{LC}}^{0.2}$$



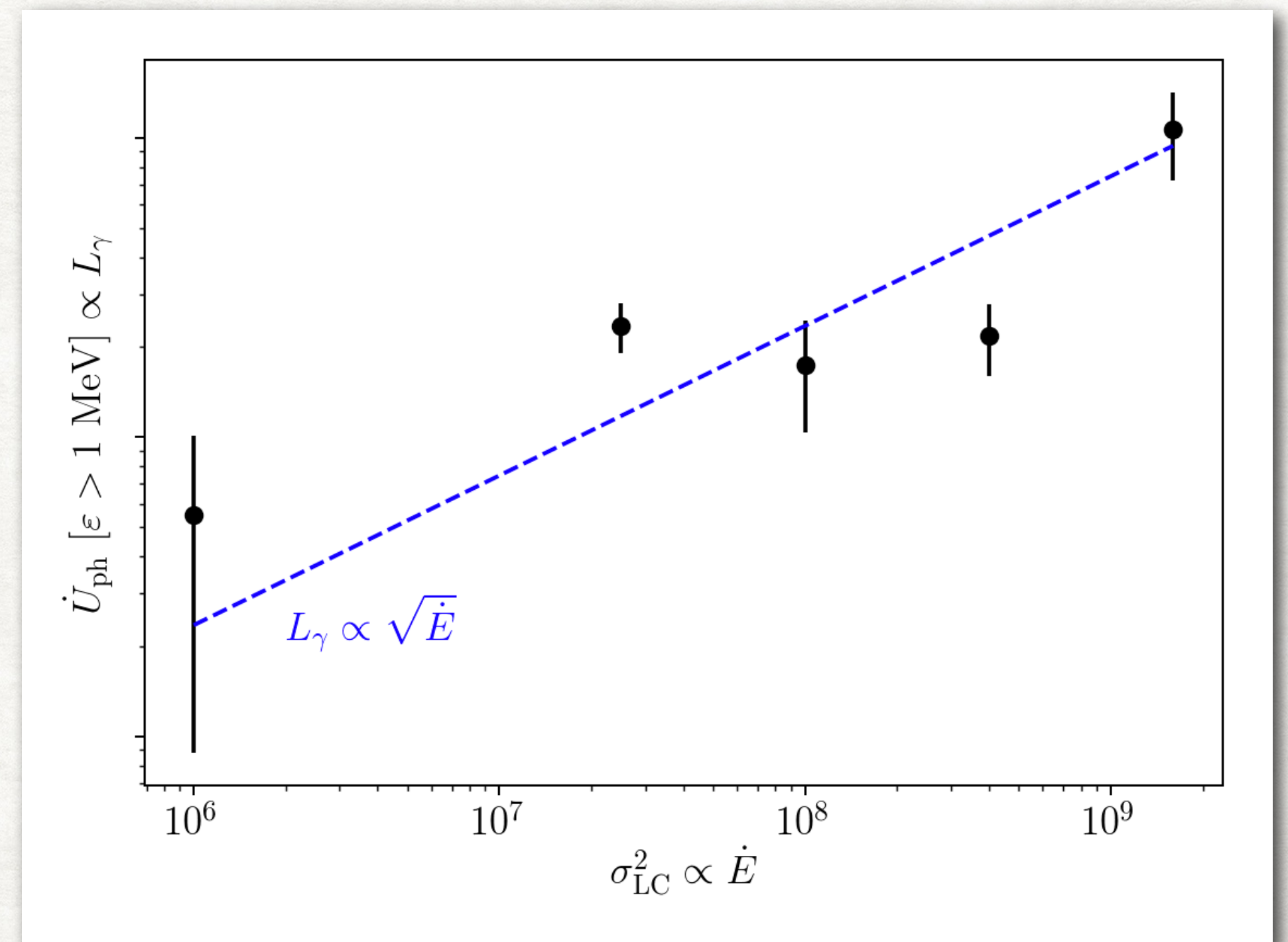
# THE ROLE OF RECONNECTION WITH PAIR PRODUCTION IN SETTING $L_\gamma$

Gamma luminosity is larger for aligned rotators than for oblique ones.  $L_\gamma/\dot{E}$  varies from 1% for orthogonal rotator to 10% for near aligned. Obliqueness effects can explain the spread in observed values of  $L_\gamma$ . In this regime  $L_\gamma \propto \dot{E}$ .

Pair formation in the current sheet decreases magnetization and lowers maximum particle energy, and radiative efficiency decreases. Also, reconnection slows down. This leads to slower dependence.



Abdo et al 2013





# CONCLUSIONS

- Electrodynamically self-consistent, working magnetospheric models with pair formation and emission are now available using PIC simulations. GR frame-dragging is essential for polar cascades!
- Paradigm change — current sheet beyond LC is effective particle accelerator and the site of majority of high-energy emission. Need global simulations for modeling the right field shape! Forced reconnection with self-consistent pair formation needs to be studied.
- Light curves and spectra are consistent with synchrotron radiation for gamma-ray and below. IC is needed for TeV.
- Pair creation in the current sheet beyond the LC makes the synchrotron cutoff energy to depend weakly on  $B_{LC}$ . When pair creation is weak,  $L_{\gamma}$  is proportional to  $\dot{E}$ . When pair loading is strong,  $L_{\gamma}$  scales as  $\sqrt{\dot{E}}$ .
- Radio emission is likely caused by the non-stationary discharge at the polar cap — first signatures of this are seen in global simulations. More on radio and multiwavelength predictions to come!