

Time-Dependent, Multi-Wavelength Models for Active Flares of Fermi Blazars

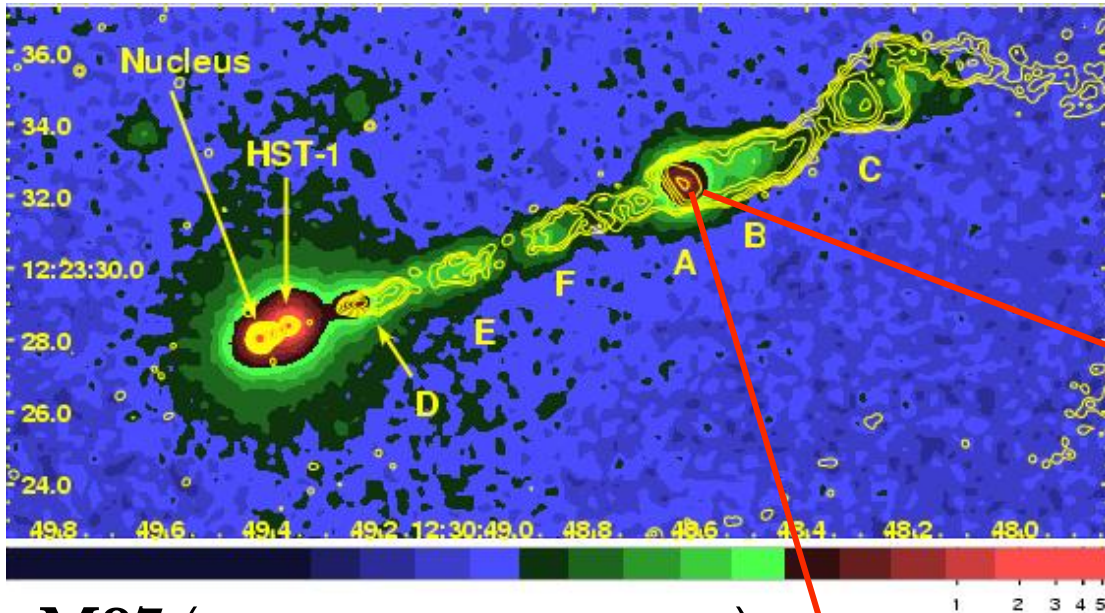
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Fermi Symposium, Baltimore, MD, October 16th 2018

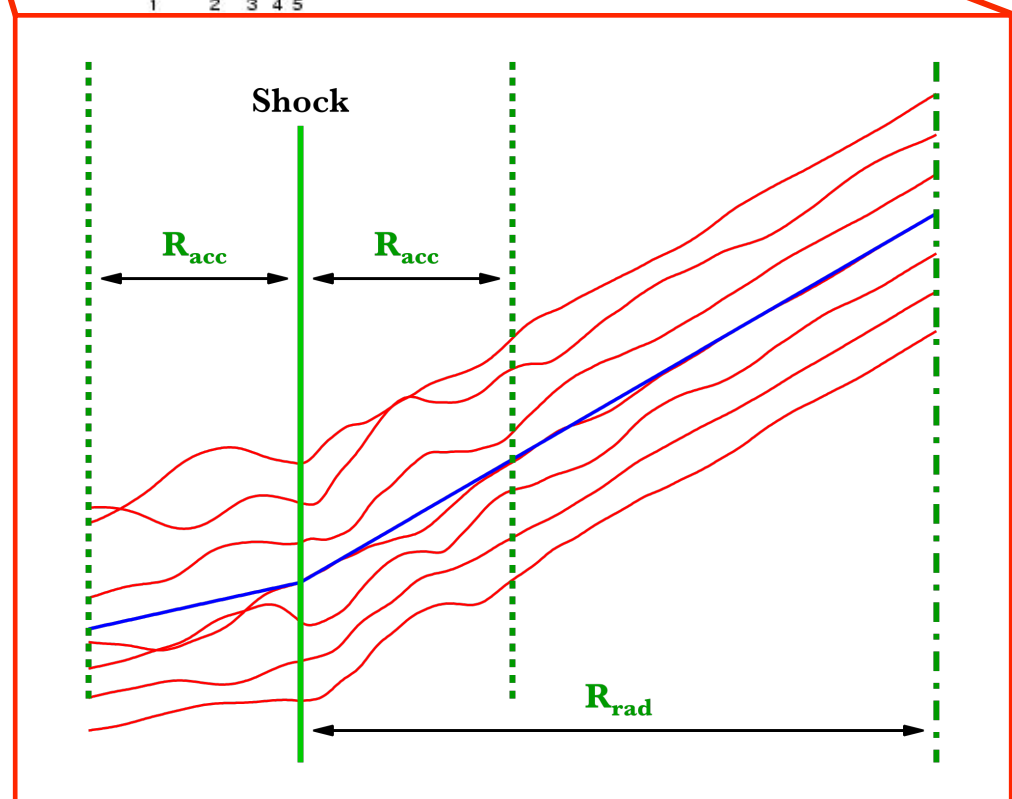


Confined Acceleration + Radiation Zones

Two-zone

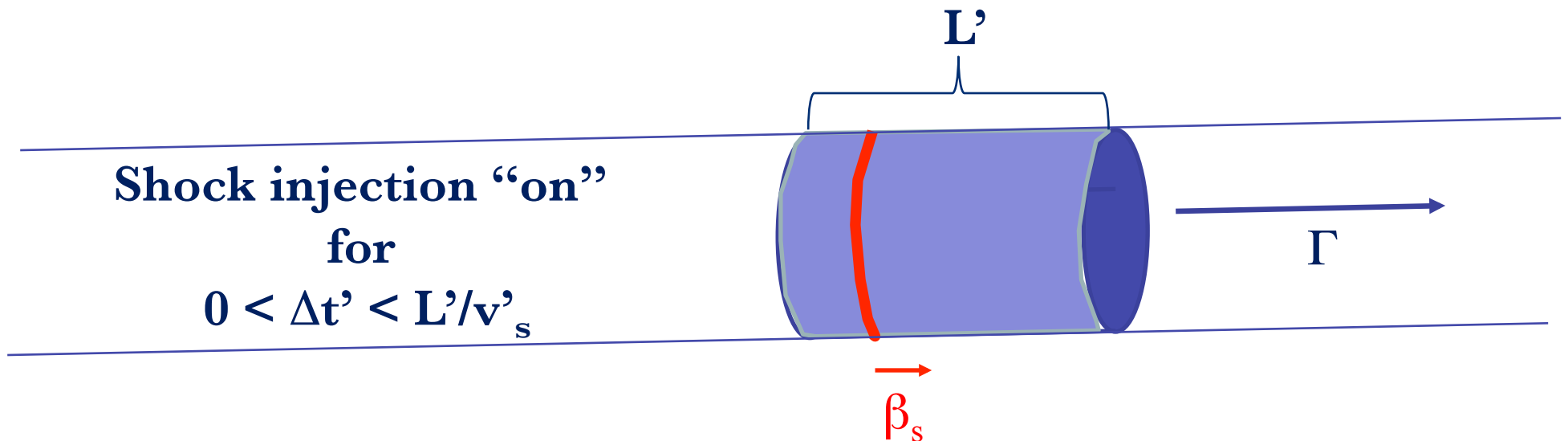
M87 (Harris & Krawczynski 2006)

- *Right:* Schematic of our blazar model geometry, consisting of a region proximate to the shock that is the **acceleration (injection) zone** which is embedded in a **much larger radiation zone**.
- Also depicted is a turbulent field, signified by the **red field line** projections.



Numerical Scheme

- Injection spectra $Q_e(\gamma, t')$ from turbulence characteristics + MC simulations of DSA
- Injection from small acceleration zone (shock) into larger radiation zone
- Time-dependent leptonic code based on Böttcher & Chiang (2002)
- Radiative processes:
 - Synchrotron
 - Synchrotron self-Compton (SSC)
 - External Compton (EC: dust torus + BLR + direct accretion disk)



$Q_e(\gamma, t') = Q_e(\gamma) H(t'; 0, Dt')$, primes denoting comoving jet frame.

Time-Dependent Electron Evolution with Radiative Energy Losses

The **acceleration timescale for electrons is short**: it scales as the cyclotron period (for shock drift and diffusive acceleration):

$$t_{\text{acc}} \sim \eta t_{\text{cyc}} = \eta \frac{2\pi\gamma\beta m_e c}{eB} \ll t_{\text{cool}}, t_{\text{dyn}}$$

for most electrons, i.e. those below around 300 MeV in blazars.

- For time evolution of blazar flares, we therefore use a shock-acceleration electron distribution (from Summerlin & Baring 2012) as an **instantaneous injection** $Q_e(\gamma)$.

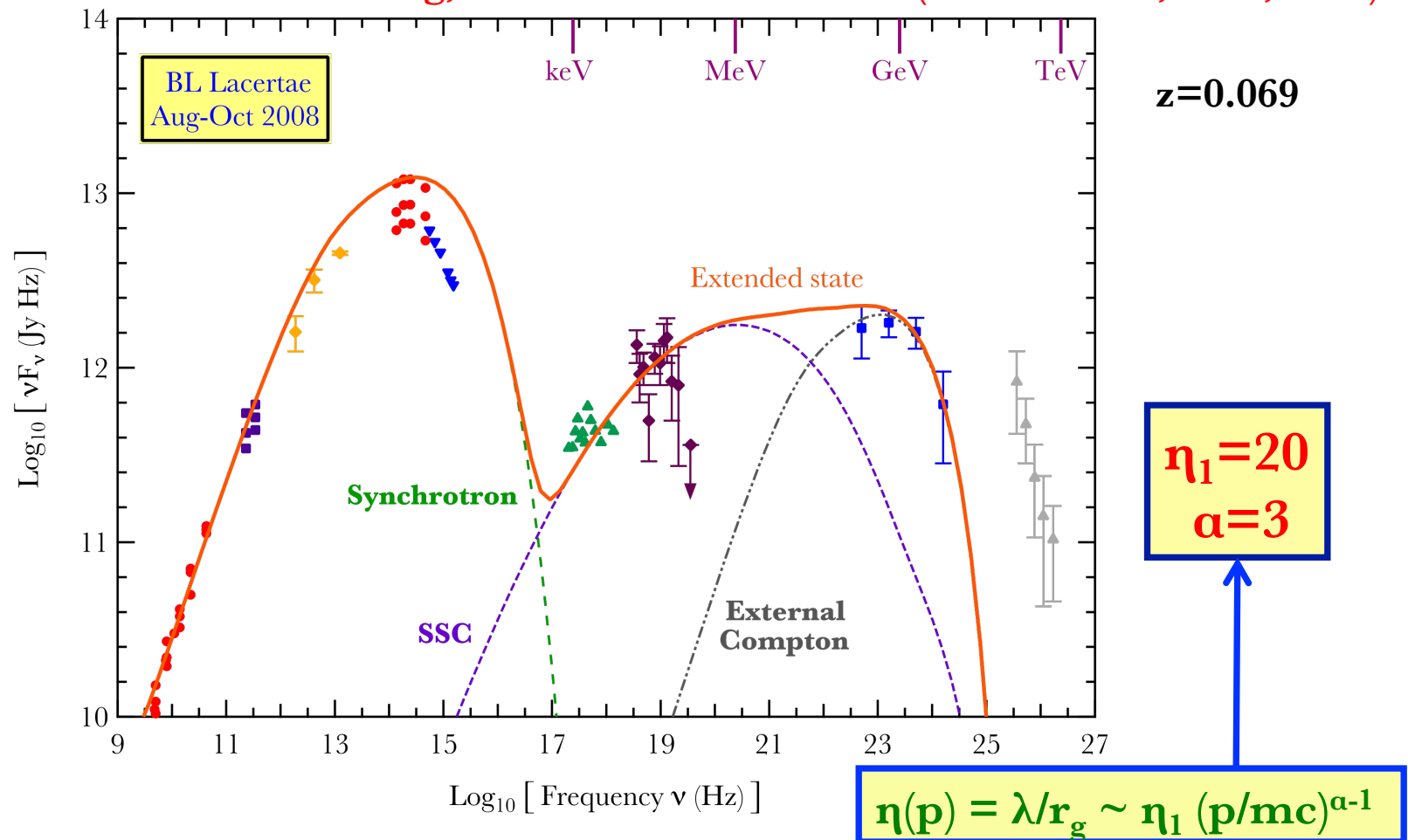
- Then **solve a Fokker-Planck equation** for electron distribution evolution in the jet frame:

$$\frac{\partial n_e}{\partial t'} = -\frac{\partial}{\partial \gamma} (\dot{\gamma} n_e) + Q_e(\gamma, t') - \frac{n_e}{t_{\text{esc}}} \quad \text{for } n_e \equiv n_e(\gamma, t').$$

This includes competition between acceleration and cooling.

One-zone Multiwavelength SSC fits to BL Lacertae

Baring, Böttcher & Summerlin (MNRAS 464, 4875, 2017)



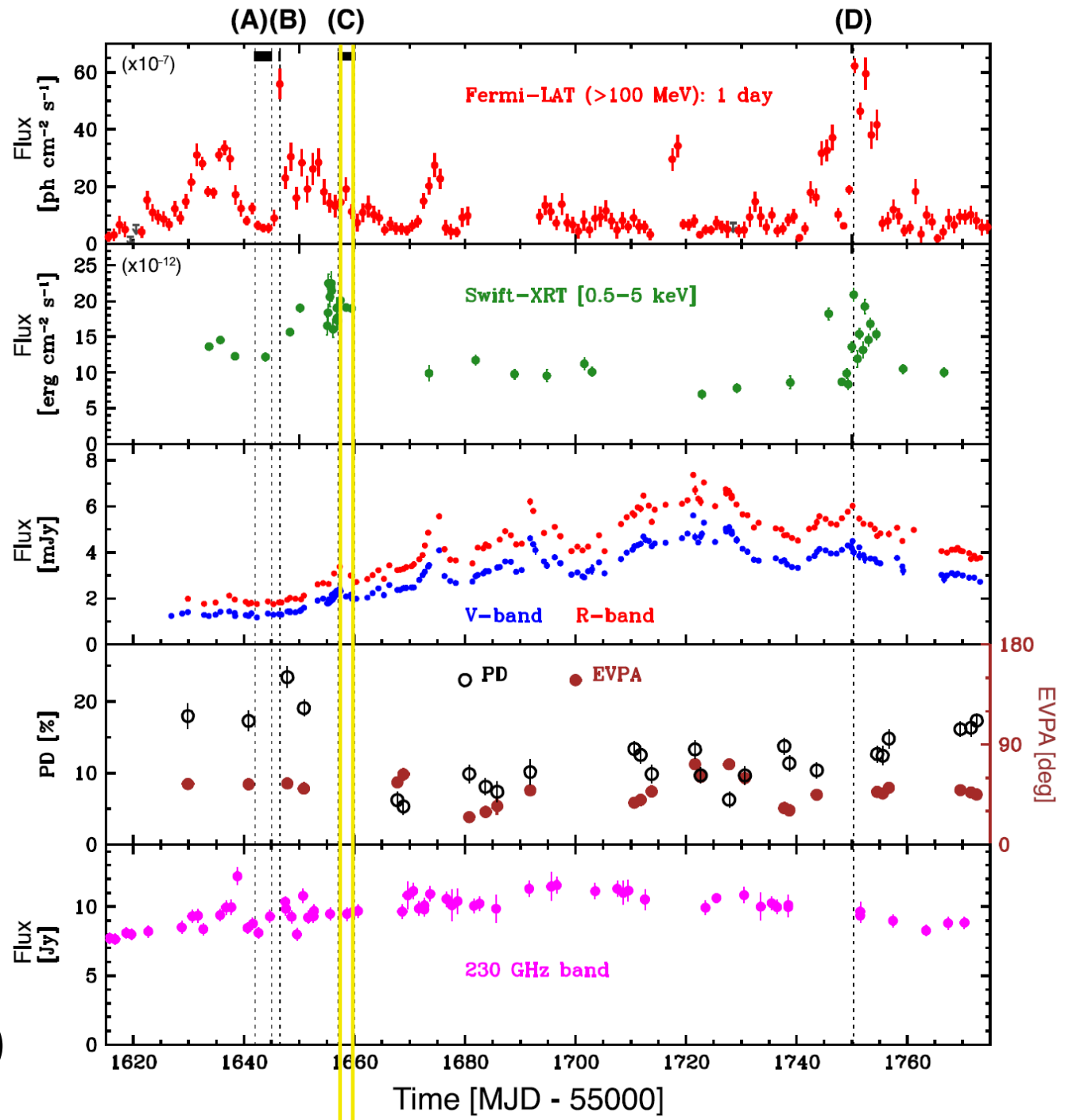
- SSC explains X-rays but cannot fit gamma-rays; EC component added.
- Large η ($\sim 10^7$) needed to move synchrotron peak into optical (for LBLs).

FSRQ 3C279

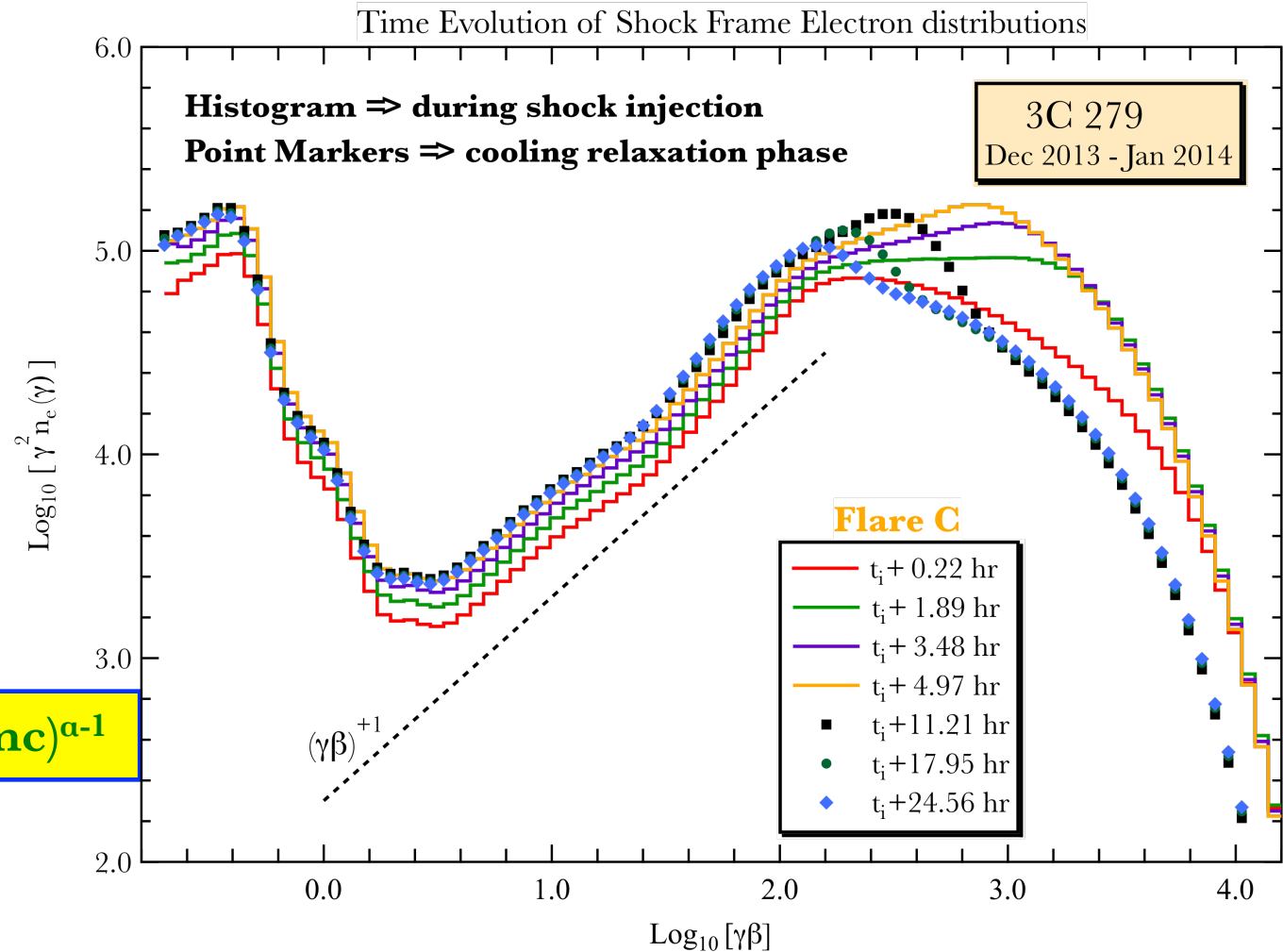
Extended
flaring
period
2013 – 2014

Variability
time scale
 ~ 1 day

(Hayashida et al. 2015)



Shock Injected and Cooled Electrons

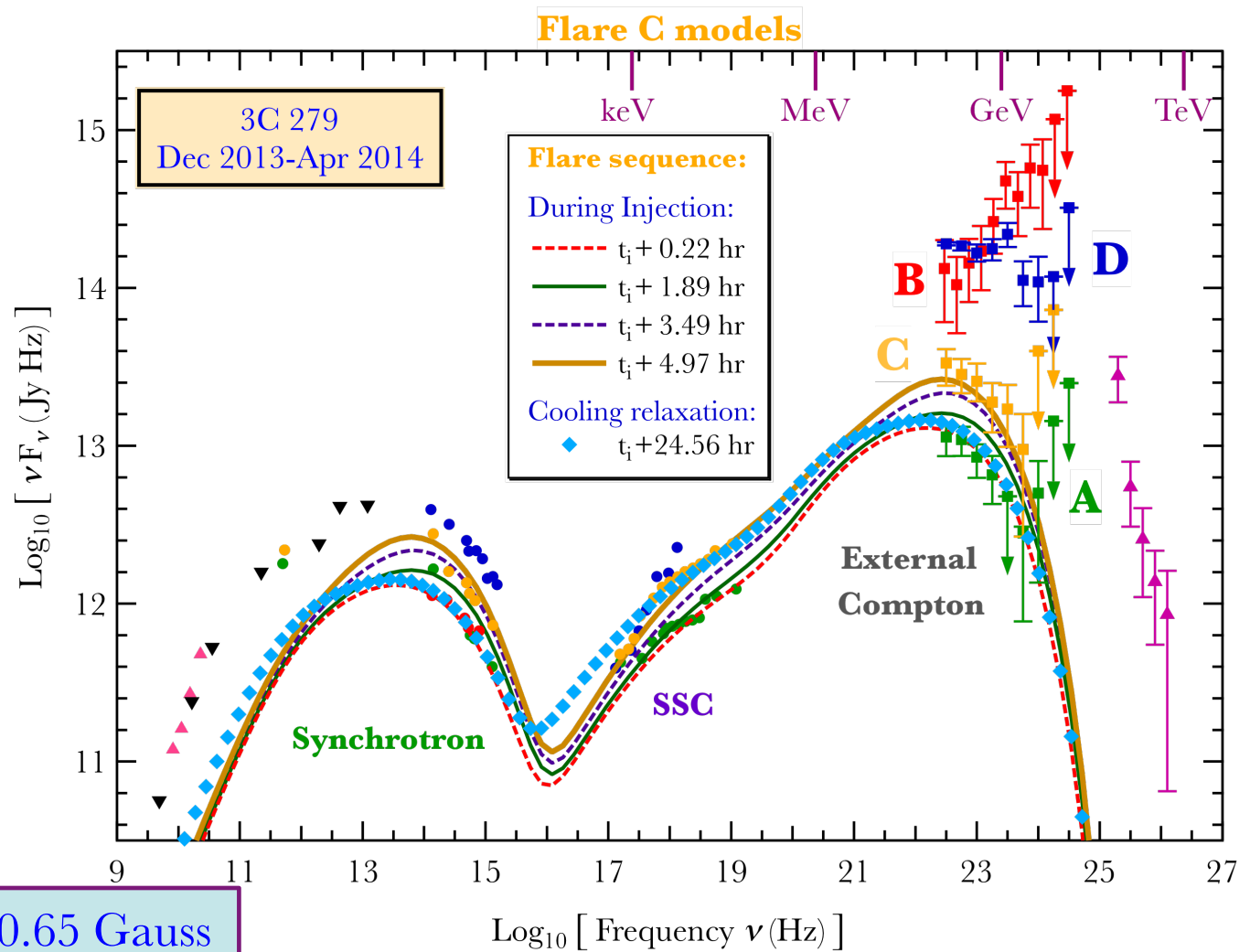


$$\eta(p) = \lambda/r_g \sim \eta_1 (p/mc)^{\alpha-1}$$

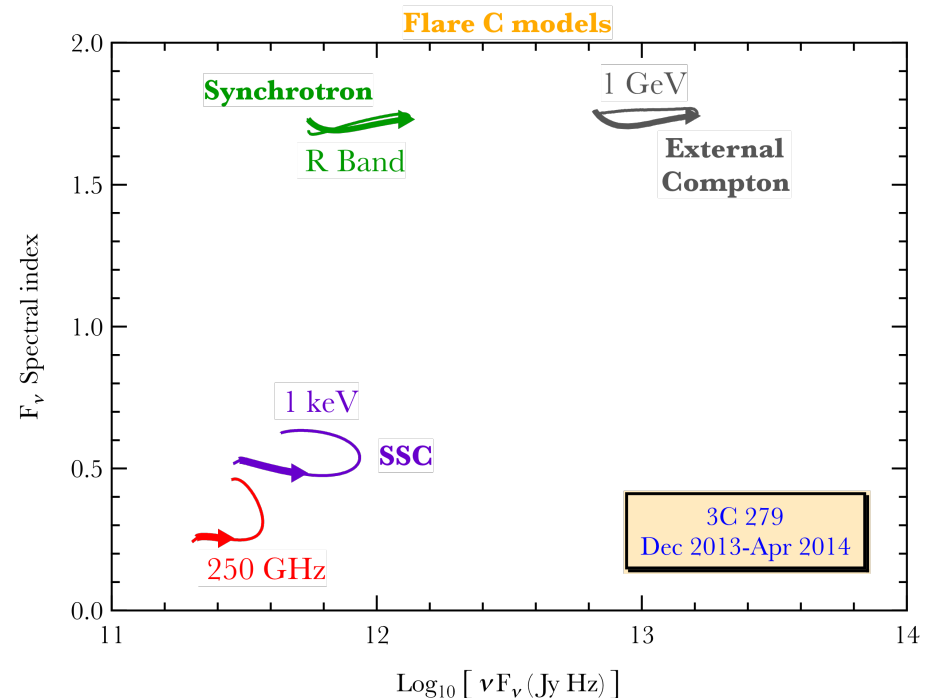
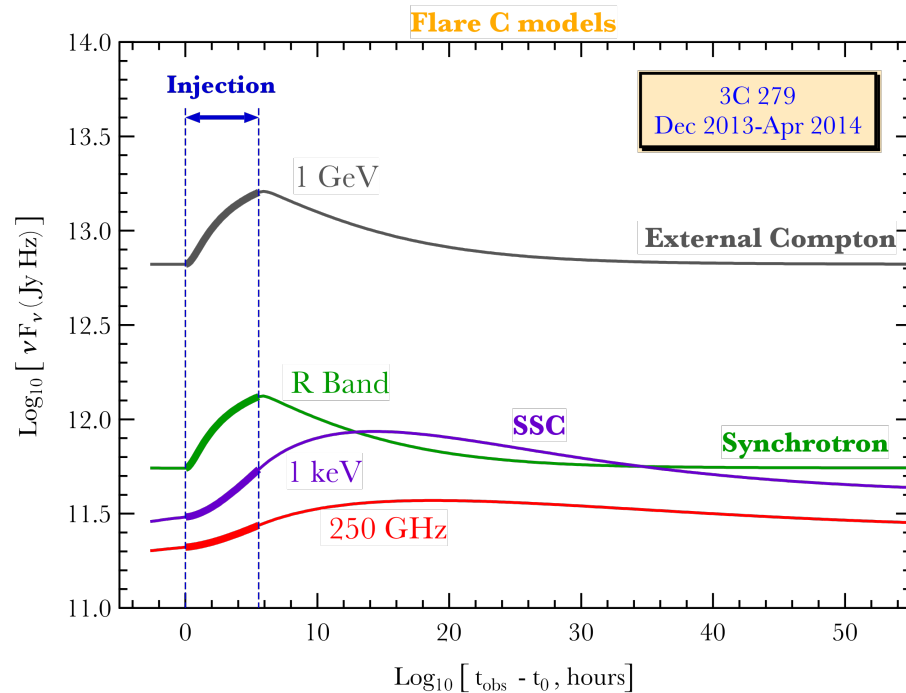
- The non-thermal particle spectral index and thermal-to-non-thermal normalization from Monte Carlo acceleration simulations are strongly dependent on η_1 (=100 in the figure) and α (=3 in figure).

Evolving MW Spectra during 3C 279 flares – Dec 2013-Apr 2014

- Time evolution of MW spectra for 3C 279 during strong flares in January 2014.
- Model curves are derived for **flare C** only, with different times during a flare corresponding to different colors and linestyles.



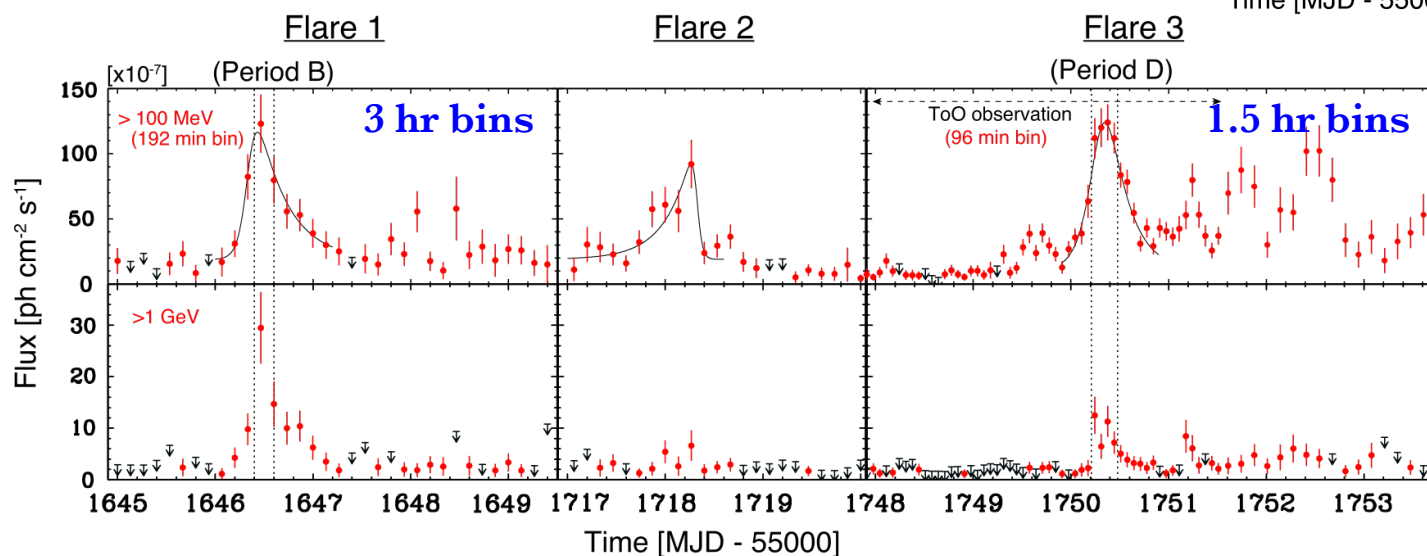
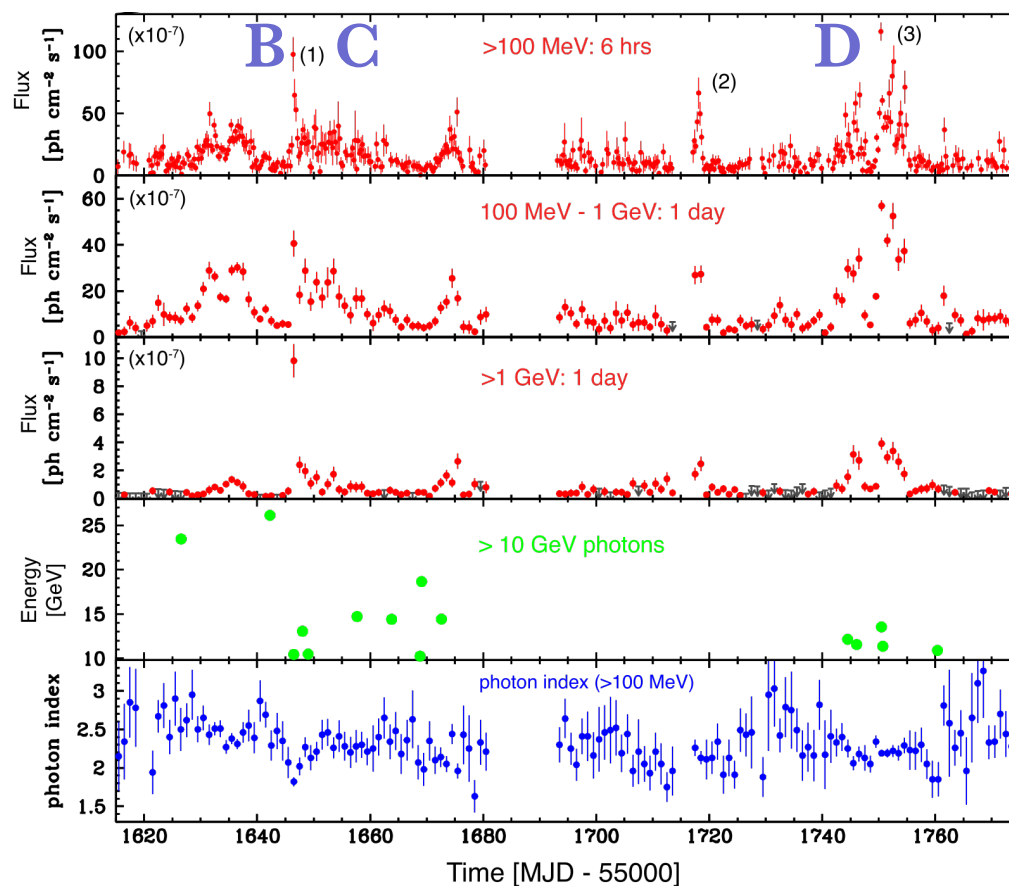
3C 279: Time Evolution Models for Flare C



- *Left panel:* Model time traces of νF_ν fluxes displaying various delays in radio to gamma-ray wavebands during Flare C. Onset of injection is at $t=0$.
- *Right panel:* Evolutionary traces in the spectral index/flux plane with time directions as indicated. The different wavebands are labeled by their components: synchrotron, SSC and external Compton. Spectral hysteresis assumes a characteristic counter-clockwise form.

3C 279: *Fermi*-LAT light curves

- Periods B and D suggest possibly asymmetric injection of duration close to GeV cooling timescales (1-3 hours)
- Hayashida et al. (2015, *ApJ* 807:79)

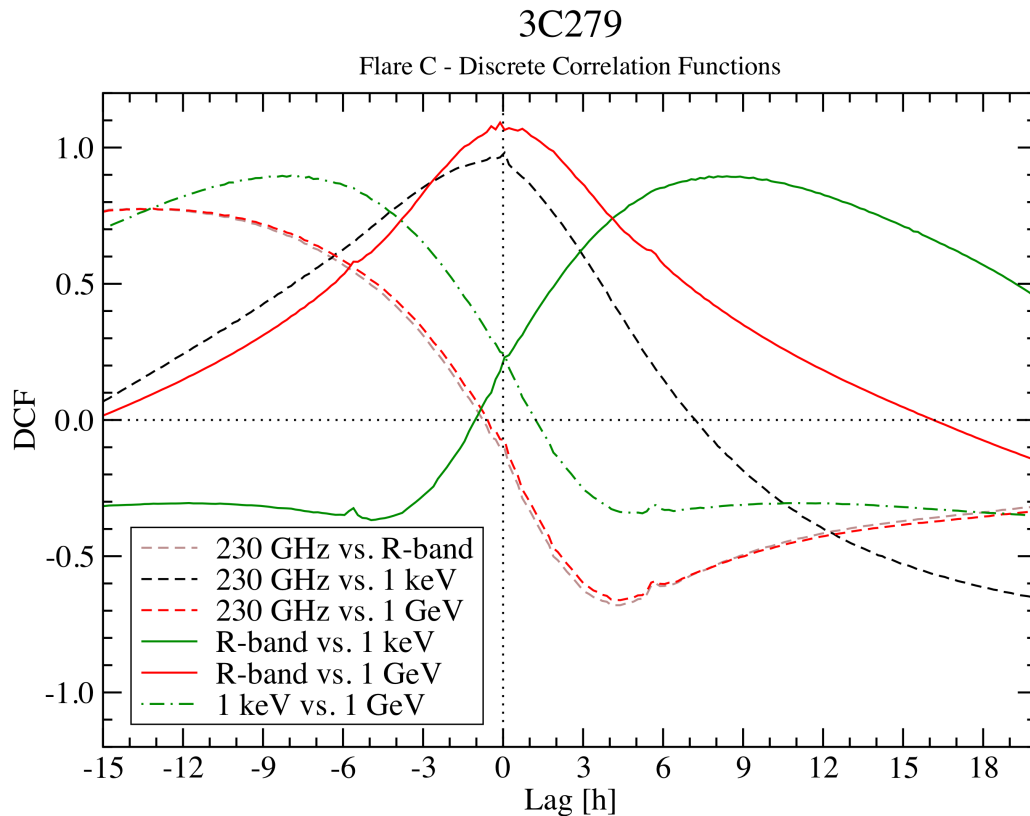


Conclusions

- **Broadband blazar spectra: X-ray/ γ -ray diagnostics on turbulence power spectra and particle diffusion.**
 - Details in Baring, Böttcher & Summerlin (MNRAS 464, 4875, 2017)
- **Coupled Monte Carlo Simulations of diffusive shock acceleration and radiation transport reveal strongly energy-dependent mean-free-path to scattering.**
 - MW fits demand $\eta = \lambda / r_g$ to be an **increasing function of p** as scales sample greater distances from the shock.
- **3C 279 (an LBL) presents a particular temporal case:**
 - Hardness/flux evolution displays **characteristic counter-clockwise spectral hysteresis** in all bands.
 - **X-rays and radio are well-correlated**, but **lag optical and γ -rays by about 5-9 hours**. => **comparative cooling times**.
 - Different lags/hysteresis are expected for other blazar types.
- **Spectroscopy and evolution present consistency in model parameter determination.**
- **Next test: modeling hard *Fermi*-LAT Flare B.**

3C279 – Flare C

Discrete Correlation Functions



The discrete correlation function is

$$C_{a,b}(\tau) \equiv \frac{1}{\mathcal{N}_a \mathcal{N}_b} \int_{-\infty}^{\infty} \mathcal{F}_a(t) \mathcal{F}_b(\tau - t) dt$$

for fluxes $\mathcal{F}_{a,b}$ in wavebands a,b, with

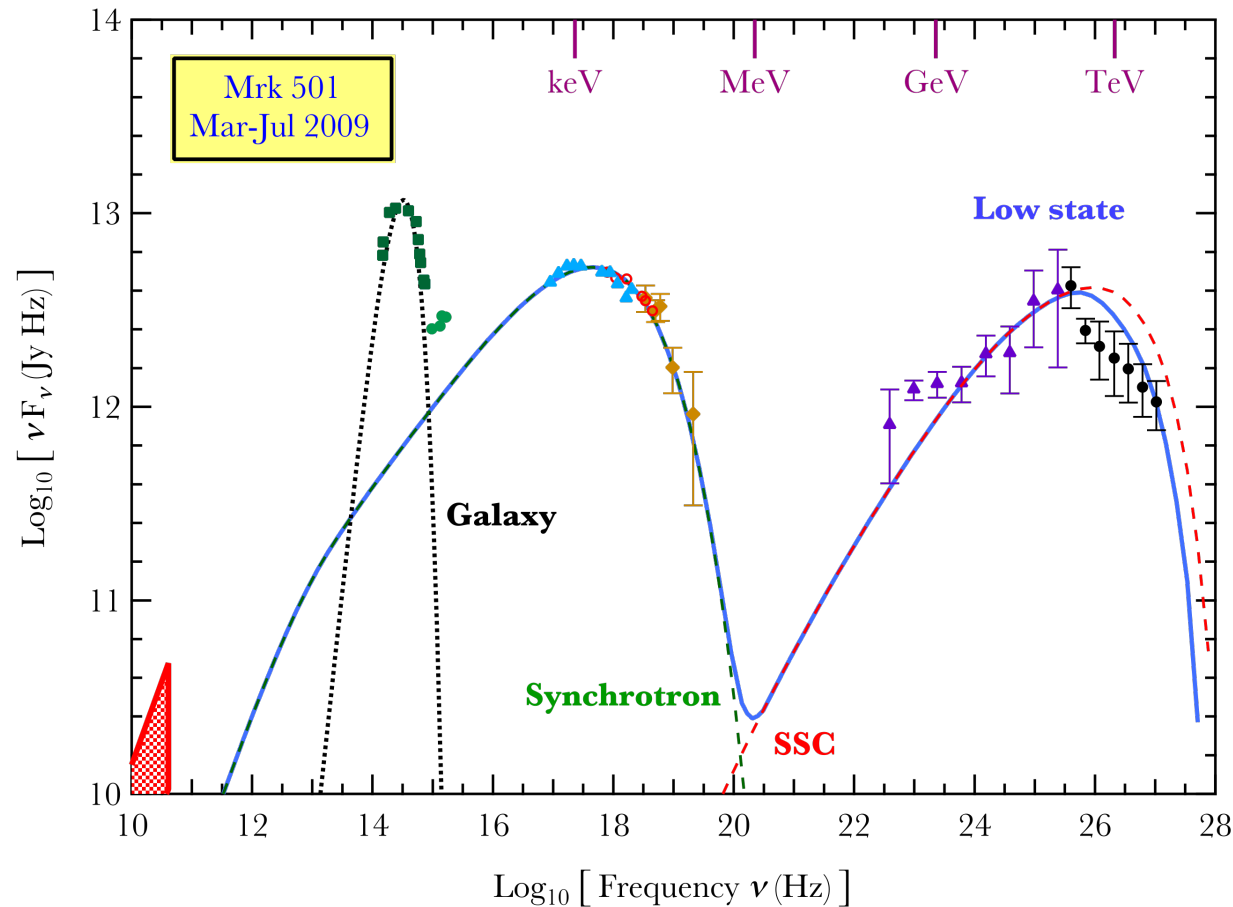
$$\mathcal{N}_a = \int_{-\infty}^{\infty} \mathcal{F}_a(t) dt$$

defining the normalizations.

After Edelson & Krolik (1988, ApJ **333**, 646).

- **DCF** measures relative fluxes between different bands.
- Optical and γ -rays **well correlated** (0 lag). X-rays and radio **well correlated** (0 lag).
- X-rays and radio **lag** optical + γ -rays by **5-9 hrs**, latter being generated by higher energy electrons, and therefore possessing a shorter response (cooling) time.

One-zone Multiwavelength SSC fits to Mrk 501



Baring, Boettcher
& Summerlin (2017,
MNRAS): steady-
state models only

$z=0.034$

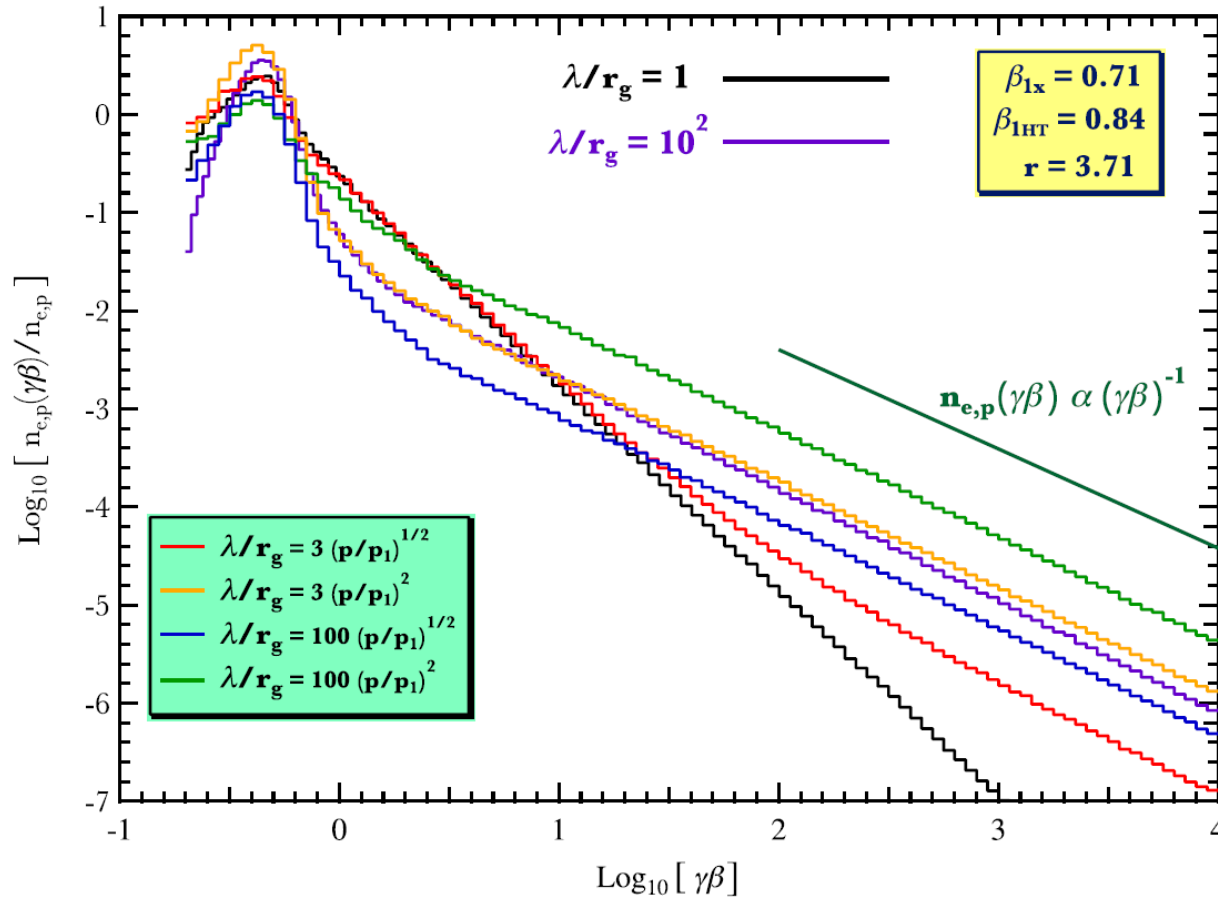
$\eta_1 = 100$
 $\alpha = 3/2$

- Synchrotron explains X-rays but cannot fit optical/UV; galaxy component added.
- Large η ($\sim 10^4$) needed to move synchrotron peak $E_{\text{max}} \sim m_e c^2 / (\eta \alpha_f)$ into X-rays.
- Need for large $\eta = \lambda / r_g$ in blazars identified by Inoue & Takahara (1996).

Shock Acceleration Injection Efficiencies

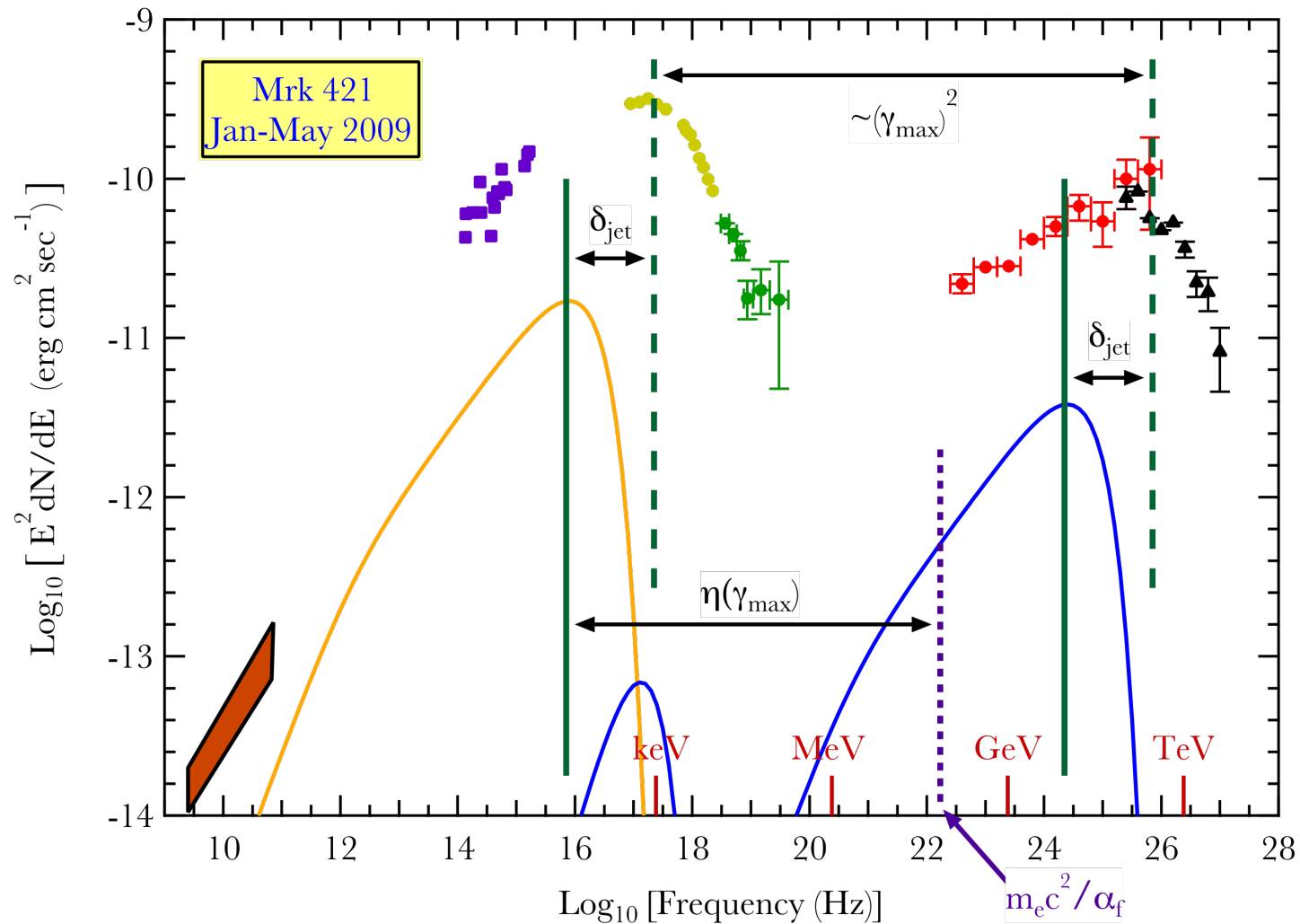
$$\eta(p) = \lambda/r_g \sim \eta_1 (p/mc)^{\alpha-1}$$

Baring et al. (MNRAS, 2017)



- The non-thermal particle spectral index and thermal-to-non-thermal normalization from Monte Carlo acceleration simulations are strongly dependent on η_1 ($=3, 100$ in the figure) and α ($=1/2, 2$ in figure), and the B-field obliquity to the shock normal!

Constraining SSC fit Parameter $\eta = \lambda_{||} / r_g$



- Large η needed to move synchrotron peak $E_{\text{max}} \sim m_e c^2 / (\eta \alpha_f)$ into X-rays.

Canonical Turbulence Power Spectrum

- Inertial range can span 1-5 orders of magnitude.
- Doppler gyro-resonance condition $\omega = \Omega/\gamma$ may not be satisfied by charges with large gyroradii;
- \Rightarrow increase of diffusive mean free path parameter $\eta = \lambda/r_g$ at large momenta.
- Expect $\lambda \propto p^2$ at long wavelengths, below stirring scale (QLT).

