

### The Crab Pulsar from HE to VHE

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# The Crab

- Remnant of historic supernova in 1054 A.D.
- Distance 2 kpc
- One of the most energetic pulsars 4.6 x 10<sup>38</sup> erg s<sup>-1</sup>
- One of the brightest gamma-ray pulsars
- Powers the brightest
  VHE gamma-ray source, the
  Crab Nebula





### Birth of a Pulsar



# Birth of a Pulsar

- Supernova remnant of a massive star (5-10 solar masses)
- Gravitational core collapse stopped by neutron degeneracy pressure
  - compact object of ~1.4 solar masses with radius of 10-20 km
- Conservation of magnetic energy and angular momentum during collapse
  - fast rotation ~100 Hz and strong magnetic dipole field ~ $10^{12}$  G

### → powerful generator of electromagnetic energy

### Gamma-Ray Production in the Vicinity of Pulsars

 $\gamma$ -ray emission (GeV and higher) is produced by relativistic particles

in the:

1. magnetosphere of the neutron star

-> pulsed emission

2. nebula powered by the pulsar  $\rightarrow$  steady emission





Radiation from a Pulsar-wind-nebula complex



### Gamma-Ray Emission from the Pulsar Magnetosphere



sketch from Alice Harding

- Stable vacuum gaps in the magnetosphere
- Potentials of  $\sim 10^{12} \, eV$ 
  - $\rightarrow$  particle acceleration
- Curved particle trajectory  $\rightarrow \gamma$ -ray emission
- Interaction with low energy photons
  - $\rightarrow$  inverse Compton emission
- Gamma-rays are subject to absorption

### Crab Pulsar Pulse Profile



8

### Pulsed Emission from the Crab



From Thompson et al 1999



- First good measurement of the HE part of spectrum
- No answer to how the cutoff looks like

# Two favored models

### Outer gap

 Emission region far out in the magnetosphere -> exponential cutoff (b=1)

Parameterization of the cutoff:

F(E)\*exp(-E/E<sub>c</sub>)<sup>b</sup>

### Polar cap

 Low emission region -> super exponential cutoff (b>2)



### Fermi & MAGIC (2008)





Fermi launched in summer 2008 MAGIC achieved energy threshold of 25 GeV in fall 2007

# MAGIC Pulsar Detection at 25 GeV



#### restricts the emission region to above 5 NS radii

# Unprecedented Spectroscopy with Fermi: Phase Averaged

- Cutoff energy 5.8 GeV
- Detection up to 20 GeV
  Constrains the emission region to > 4-5 NS radii
- Exclusion of the superexponential cutoff
- Cutoff consistent with exponential cutoff



Abdo et al. ApJ (708) 1254–1267, 2010

# Enough statistics to do phase resolved spectroscopy



15

Abdo et al. ApJ (708) 1254–1267, 2010



16

Abdo et al. ApJ (708) 1254–1267, 2010

# Is the Picture complete?

- Sensitivity limited above 10 GeV with Fermi
- Why bother, outer gap fits -> lets go home



# The highest Energies



Science, MAGIC Aliu et al. (2008)

Is the spectrum harder than described by an exponential cutoff?

#### **Assumption:**

The spectrum at the highest energies is described by an exponential cutoff

### Pulsed Emission above 60 GeV



Two independent data sets show pulsed emission above 60 GeV at the position of the interpulse

... at flux levels not predicted by model calculations!

### Detection of the Crab Pulsar above 100 GeV

Science, VERITAS Aliu et al, 2011



Detection alone rules out exponential cutoff



### The GeV – TeV Connection



Good description with smooth broken power law (does not exclude that two emission processes are at work) MAGIC data between 25 GeV and 100 GeV fill in blank space MAGIC stereo data released in 2011 confirm VERITAS measurements



Romani, 1996

### Curvature Radiation >100 GeV?

Lyutikov, Otte & McCann (2011) arXiv:1108.3824

Acceleration in gaps is radiation reaction limited:



For a magnetic dipole field it follows that the maximum Lorentz factor is:

$$\gamma_b = 9 \times 10^7 \eta^{1/4} \sqrt{\xi}$$

= radius of curvature of magnetic field lines in units of light cylinder radius = strength of the electric field  $E = \eta B$  typically between 0.01 and 0.1 due to screening

### **Curvature Radiation?**

Maximum curvature radiation (CR) gamma-ray energy is:

$$\epsilon_{\rm br} = \frac{3}{2}\hbar \frac{c}{R_{\rm c}} \gamma_b^3$$

Using the maximum Lorentz factor it follows:

$$\epsilon_{\rm br} = 150 \,{\rm GeV} \eta^{3/4} \sqrt{\xi}$$

#### >100 GeV emission only possible if: $\eta \sim 1$

CR not excluded but unlikely because of efficient pair production:

CR photons interact with photons coming from the cooling neutron star (Crab is just 1000 years old) and create electron positron pairs -> screening of electric field:  $\eta \ll 1$ 

### **Inverse Compton Scattering**



# **Inverse Compton (IC) Scattering**

Assumed seed photon fields for IC are the observed ones UV photon field (luminosity 10<sup>34</sup> erg s<sup>-1</sup> and typical photon energy 1 eV) and X-ray fields (luminosity 10<sup>36</sup> erg s<sup>-1</sup> and typical photon energy 1 keV)

Minimum Lorentz factor to be in Klein-Nishina regime:

$$\gamma_{KN} = \frac{1}{4} \frac{m_e c^2}{\epsilon_{\text{soft}}} \approx 1.2 \times 10^5 \, \epsilon_{UV,0}^{-1} \approx 1.2 \times 10^2 \, \epsilon_{X,3}^{-1}$$

Energy loss/gain equation:

$$\dot{\epsilon} = ec\eta B - \frac{2}{3}\frac{e^2}{c}\gamma^4 \left(\frac{c}{R_c}\right)^2 - \frac{4}{3}\left(\frac{m_e c^2}{\epsilon_{\text{soft}}}\right)^2 U_{\text{soft}}\sigma_T c,$$

### Inverse Compton Scattering: The Primaries

Primaries = Particles with maximum Lorentz factors

Use curvature radiation limited Lorentz factor  $\sim 10^7$ 

-> IC takes place in KN regime

IC losses are comparable to CR losses if

$$L_{\text{soft,crit}} = \eta \frac{B_{NS} R_{NS}^3 \Omega \epsilon_{UV}^2}{e^3} = \begin{cases} 10^{35} \, \text{ergs}^{-1} \, \epsilon_{UV,0}^2 \eta_{-2} \\ 10^{41} \, \text{ergs}^{-1} \epsilon_{X,3}^2 \eta_{-2} \end{cases}$$

-> IC losses are comparable to CR for IC on UV photons

#### **Energetics:**

Maximum IC photon energy is  $\epsilon_{\gamma} \approx \gamma_b m_e c^2$  ~ 15 TeV

But IC luminosity not enough to explain VHE observations

### Inverse Compton Scattering: The Secondaries

Secondaries = pair created electrons and positrons from CR photons

Multiplicity factor of secondaries lambda is about 100

Wang & Hirotani 2011

 $\gamma_p \approx \gamma_b / \lambda = 3 \times 10^5 \eta_{-2}^{1/4} \sqrt{\xi} \, \lambda_2^{-1}$ 

-> IC takes also place in the KN regime

#### **Energetics:**

Maximum IC photon energies  $\epsilon_{\gamma,p} \approx \gamma_p m_e c^2 = 150 \,\text{GeV}\,\eta_{-2}^{1/4}\sqrt{\xi}\,\lambda_2^{-1}$ 

Estimated VHE IC Luminosity matches observed one

### Synchrotron Radiation from Secondaries

Use doppler boosted cyclotron emission

Secondaries also emit synchrotron radiation:

Luminosity peaks at about 3 keV (observed 100 keV)

Calculated luminosity matches observed one and requires multiplicity of secondaries of ~100

Secondaries are responsible for bulk of X-Ray and VHE emission

-> Expect similar pulse profiles in X-Ray and VHE



# A much more detailed Modelling



Outer gap modeling by K. Hirotani In MAGIC Collaboration, ArXiv:1108:5391

#### Changes with respect to older models

Emission from further out in the gap: 0.7 to 1 light cylinder radii Emission from secondaries are taken into account

#### Or maybe something completely different? Observer X-ray photons arget photons γ-rays δR, $\Gamma_{\rm w} \approx 10^6$ <sup>\_\_</sup>w[1 – cos(*θ*)] $R_{\rm w} \approx 30 R_{\rm I}$ (this work) C Poynting-flux-dominated wind $\Gamma_{\rm w} \approx 10^3$ Kinetic-energy-dominated wind $R_{\rm I} \approx 10^6 \, {\rm m}$ m e+e-1.2 acceleration zone Magnetosphere MAGIC stereo . e+e-VERITAS . P1 MM m X-ray photons + $\gamma$ -ray $R_{\rm w} = R_{\rm I}$ 0.8 $\gamma$ -ray $R_w = 30R_1$ P2 Normalized flux e+e-Wind γ-rav anisot, wind 0.6 Mmx e+e-MM> 0.4 $R_{\rm sh} \approx 3 \times 10^{15} \, {\rm m}$ 0.2 0 0.2 0.4 0.6 0.8 1.2

F. A. Aharonian, S. V. Bogovalov & D. Khangulyan Nature (2012)

Phase



#### Prediction of cutoff at 400 GeV

# Can be tested with future observations

F. A. Aharonian, S. V. Bogovalov & D. Khangulyan Nature (2012)

# **Lorentz Invariance Violation**

### **Search for new physics**

#### For example:

Microscopic structure of space time



#### In photon sector:

$$c'(E) = c + a \cdot \frac{E}{E_{\rm LIV}} + b \cdot \left(\frac{E}{E_{\rm LIV}}\right)^2$$

$$\Delta t_1 = \frac{d}{c} \cdot \frac{E_{\rm h} - E_{\rm l}}{E_{\rm LIV}} \qquad \Delta t_2 = \frac{d}{c} \cdot \frac{3}{2} \cdot \frac{E_{\rm h}^2 - E_{\rm l}^2}{E_{\rm LIV}^2}$$
  
Linear term Quadratic term

### A closer Look at the Pulses



Peak positions aligned with peak positions in radio and Fermi-LAT.

Pulses above 120 GeV 2-3 times narrower than in Fermi-LAT data

### **Testing Lorentz Invariance**

Otte, Proceedings 32<sup>nd</sup> ICRC 2011 Beijing

 Peaks at 100 MeV (Fermi) and 120 GeV (VERITAS) line up

$$\Delta t_{95\%} < 1.65 \cdot \delta \cdot P / \sqrt{2} < 100 \,\mu s$$

- Linear:  $E_{IV} > 3 \times 10^{17} \text{ GeV}$
- Quadratic:  $E_{IIV} > 7x10^9 \text{ GeV}$



#### Linear term:



# Distinguishing between Source and Propagation Effects

- Delay observed in corotating frame
- Pulsar slows down:

 $P = P_0 + t \cdot \dot{P} + \dots$ 

• 10<sup>-4</sup> per year for the Crab

Monitor how delay  $\Delta \phi$  changes over time (years):

No change : Delay is intrinsic to source

Becomes smaller : Delay is due to propagation

$$\Delta\phi(t) = \Delta t / \left(P + t \cdot \dot{P}\right)$$

