

*Micaela Caragiulo*

*University and INFN of Bari*

**Development  
of a simple  
modeling for SNR**

- **Fit of the spectrum** of some SNR published by Fermi-LAT collaboration
- **Systematic study** of the SED of SNRs

# PROCESSES INVOLVED

## Synchrotron Radiation

$$Q_\gamma(\omega) = \frac{\sqrt{3}Be^3}{2\pi m_e c^2} \frac{4\pi}{\beta c} \int \frac{dN_e}{dE_e} R\left(\frac{\omega}{\omega_c}\right) dE_e$$

[Zirakashvili, V.N. Aharonian, F., 2007, AA, 465, 695]

## Inverse Compton Radiation

$$Q_\gamma(E_\gamma) = \int \frac{dN_e}{dE_e} dE_e \int n(E_s) \sigma_{K-N}(E_s, E_e, E_\gamma) dE_s$$

[Sturmer, S.J., Skibo, J.G., Dermer C.D., Mattox J.R., 1997, ApJ, 490, 619]

## Bremsstrahlung

$$Q_\gamma(\varepsilon) = 4\pi n_H \int \frac{dN_e}{dE_e} \frac{d\sigma_{B-H}}{d\varepsilon} dE_e$$

[Baring, M.G., Ellison, D.C., Reynold, S.P., Grenier, I.A., Goret, P., 1999, ApJ, 513, 311]

## Hadronic production

$$Q_\gamma(E_\gamma) = \frac{4\pi}{\beta c} n_H \int \frac{dN_p}{dE_p} \frac{d\sigma(E_p / E_\gamma)}{dE_\gamma} dE_p$$

[Kamae, T., et al. 2006, ApJ, 647, 692]

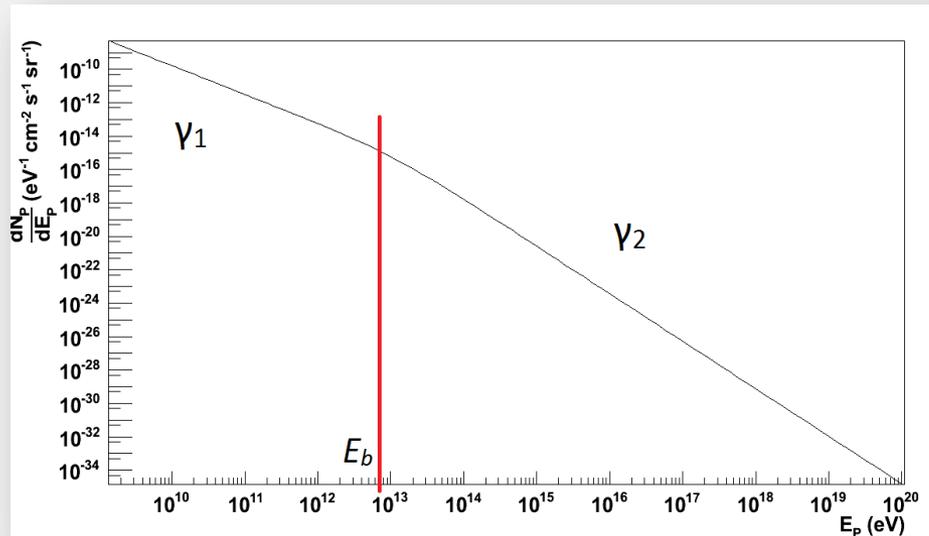


# INJECTION OF PARTICLES

We assumed the same broken power law to describe the population of protons and electrons accelerated:

$$\frac{dN}{dE} = A \left( \frac{E}{E_0} \right)^{-\gamma_1} \left[ 1 + \left( \frac{E}{E_b} \right)^{\frac{(\gamma_2 - \gamma_1)}{\gamma_3}} \right]^{-\gamma_3}$$

- $E_0 = 1 \text{ GeV}$
- $\gamma_3 = 1$



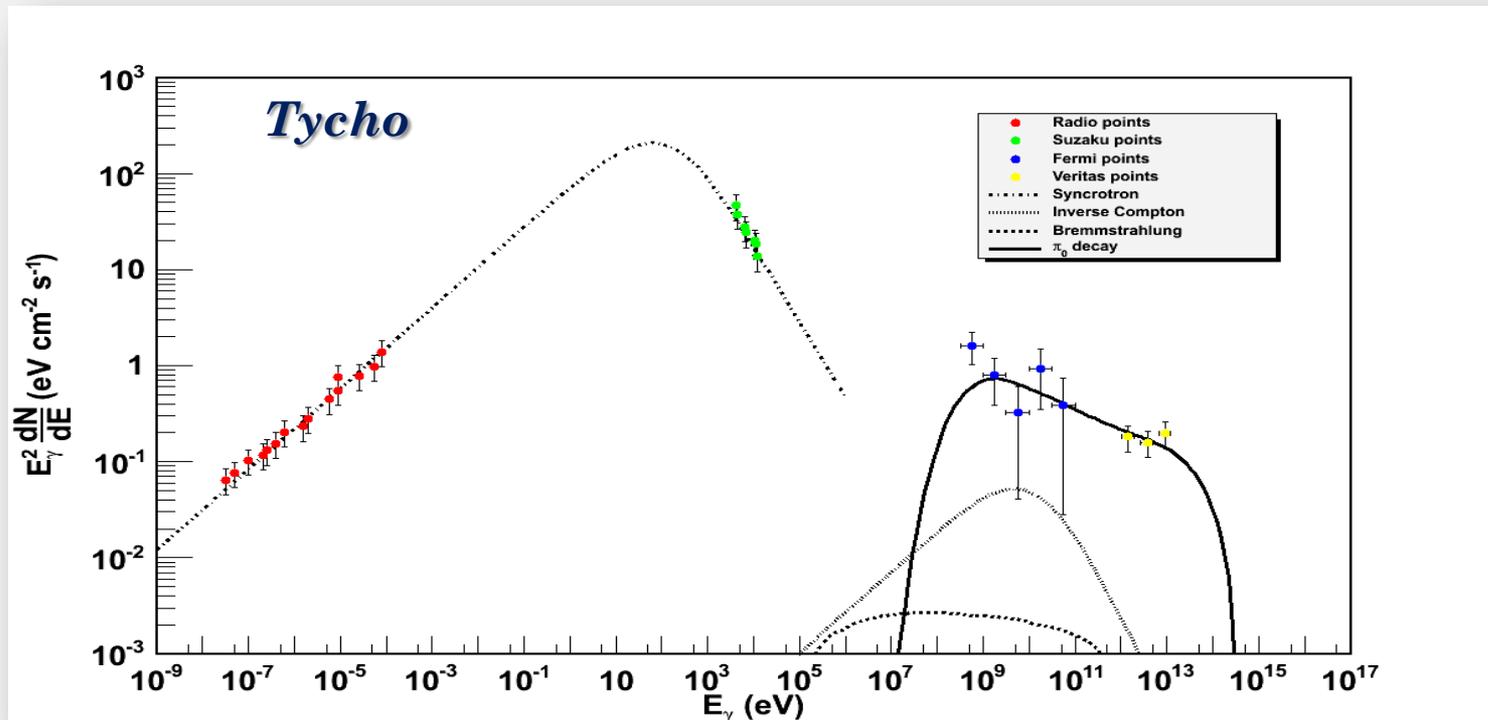
The parameters to be determined thanks to fit are the normalization constant, the two indices and the energy break.



# SED OF SNR

$$F_{\gamma,SNR}(E_{\gamma}) = \frac{Q_{\gamma}(E_{\gamma})V_{SNR}}{4\pi d^2}$$

- $d$  is the distance of the SNR from the Earth,
- $V_{SNR}$  is the volume of the sphere that approximate the expanding SNR.



# *IMPROVEMENTS*

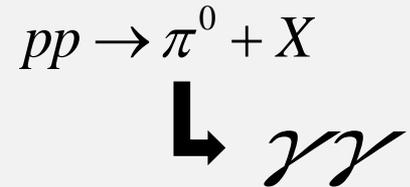
- ✓ use of the Monte Carlo code Fluka for the determination of the proton- proton cross section;
- ✓ development of a code for the temporal evolution of SNR.





# HADRONIC SCENARIO

The process we are interested is described by the following:



We used **Kamae et al. (2006)** for parametrizing the p-p cross section, thanks to which we have determined the emissivity of the photons and thus the SED of the SNRs studied.

$$Q_{\pi^0}(E_\gamma) = \frac{4\pi}{\beta c} n_H \int \frac{dN_p}{dE_p} \frac{1}{E_\gamma} \frac{d\sigma(E_p / E_\gamma)}{d \log(E_\gamma)} dE_p \quad \longrightarrow \quad F_{\gamma,SNR}(E_\gamma) = \frac{Q_\gamma(E_\gamma) V_{SNR}}{4\pi d^2}$$

where  $\mathbf{d}$  is the distance of the SNR from the Earth and  $V_{SNR}$  is the volume of the sphere that approximate the expanding SNR.



# ***FIT OF EXPERIMENTAL DATA***

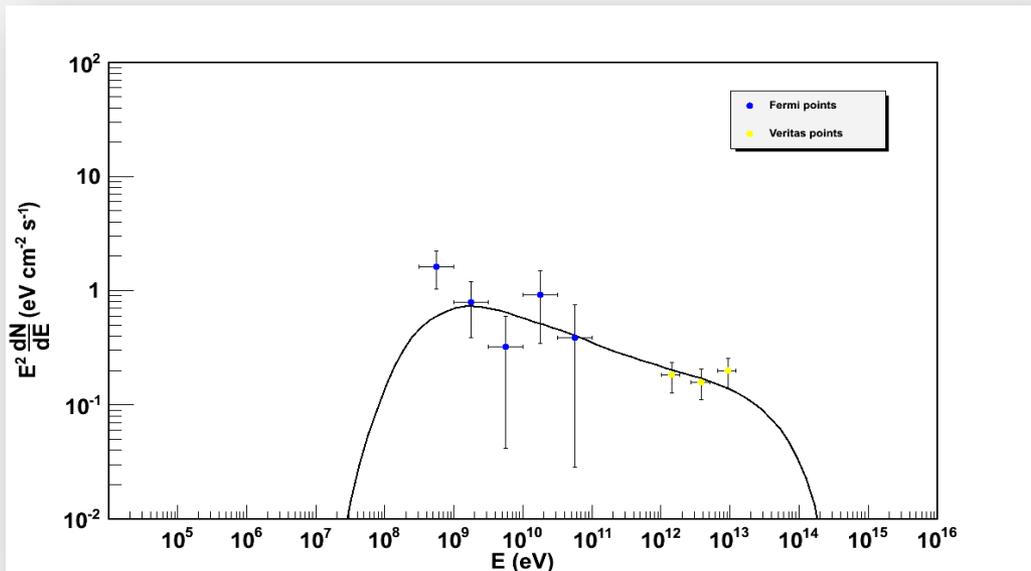
We fit the public GeV-TeV experimental data of several SNRs. We choose **young SNRs** (e.g. Tycho and Vela jr.) that show different spectral shapes in GeV band, and **molecular cloud SNRs** (e. g. IC 443) that instead interacts with the dense environment.

In our model we fixed the values of SNR *radius, the density of the environment and the distance of the SNR from the Earth.*



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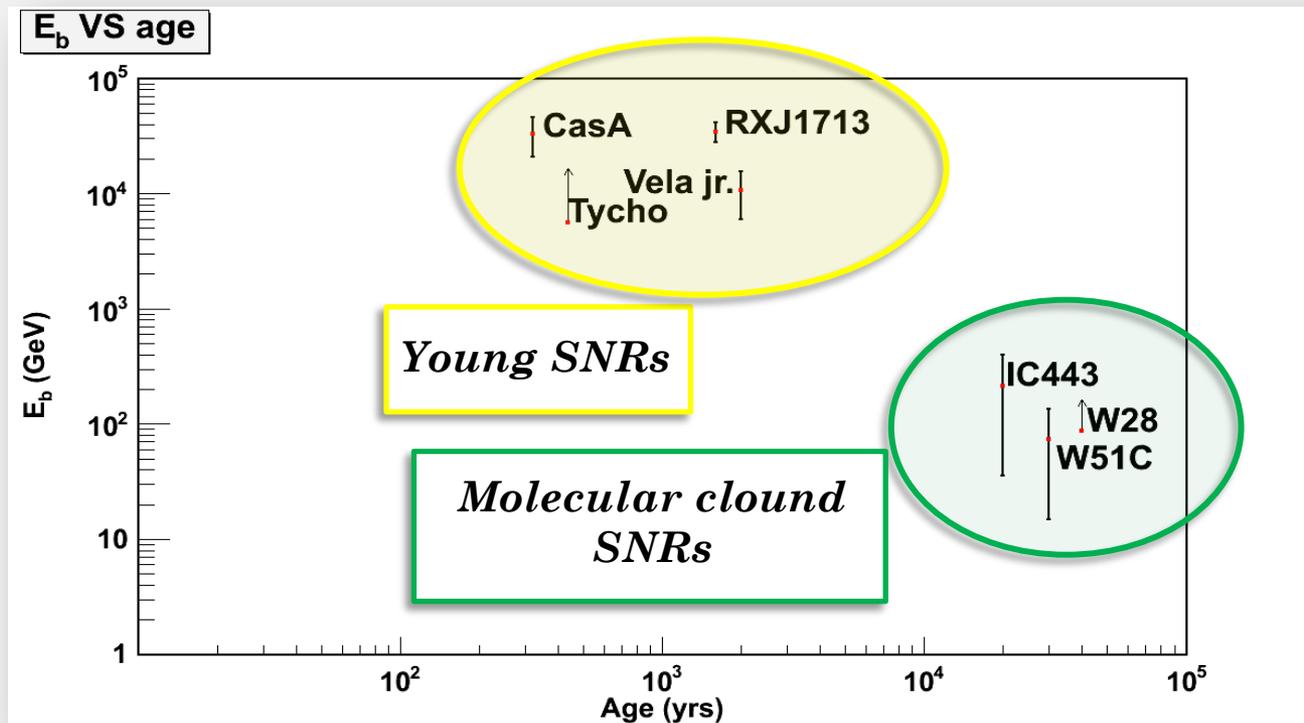
## **TYCHO SNR**

- ***SN 1572***
- ***SN type: Ia***
- ***Age: 349 anni***
- ***Distance: ~ 3.5 kpc***
- ***Radius: ~ 3.7 pc***
- ***$n_H = 0.24 \text{ cm}^{-3}$***

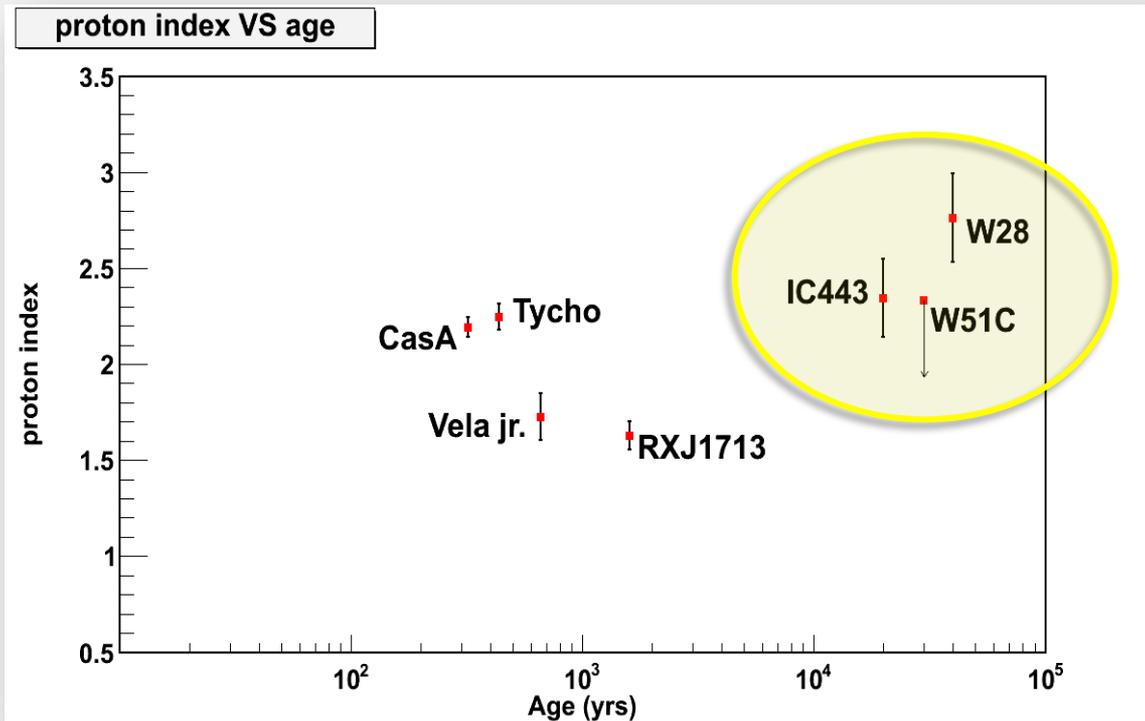


# ***FIT OF EXPERIMENTAL DATA***

This plot shows the behavior of the energy break as a function of the SNRs age. We can distinguish the young SNRs from the molecular cloud ones.



# ***FIT OF EXPERIMENTAL DATA***

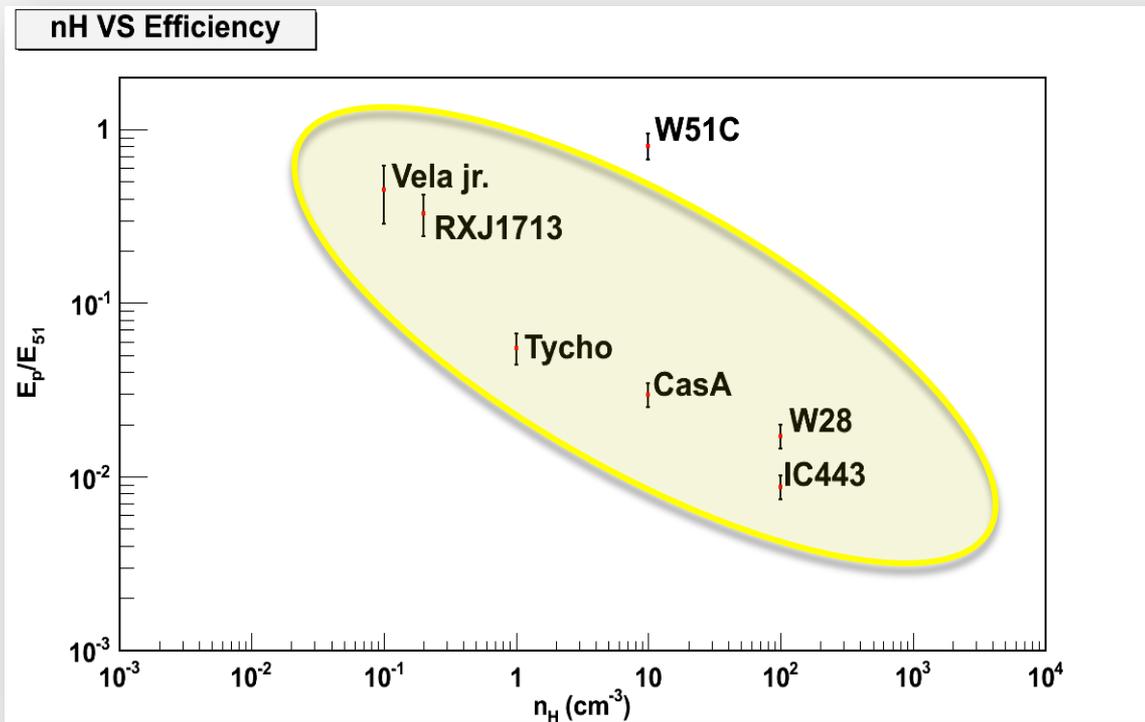


We relate the first index of the broken-power law and the environment density.

The greater the density softer is the index, indicating a possible different propagation of CRs in different environments.

# ***FIT OF EXPERIMENTAL DATA***

In this plot we show the behavior of the efficiency (normalized to total kinetic energy of  $10^{51}$  erg.) as a function of the density of the environment.



The role of the density is to enhance the GeV brightness even keeping at a lower level the acceleration efficiency.

