## GeV-TeV emission from inverse-Compton scattering processes in GRBs

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## GRB-shock acceleration

### Protons-

shock acceleration  $\sim 10^{20} eV$  (e.g. Waxman 1995; Vietri 1995) photo-meson process:

(e.g. Waxman & Bahcall 1997; Bottcher & Dermer 1998)

 $\gamma + p \rightarrow \Delta^{+} \rightarrow \pi^{0} + p \qquad \qquad \pi^{0} \rightarrow \gamma + \gamma$ 

proton synchrotron radiation

(e.g. Vietri 1997; Totani 1998)

### Electrons-

Shock acceleration: ~10 TeV

 $t_{cool} \sim t_{acc} \Rightarrow \gamma_{e,Max} \le 10^7$ X-ray afterglows  $\rightarrow \gamma_{e,Max} > 10^5$ 

e.g. Li & Waxman 2006

### Leptonic inverse Compton scattering

# Inverse-Compton scattering in GRB shocks



Internal shock IC: e.g. Pilla & Loeb 1998; Razzaque et al. 2004

# External shock ICreverse shock IC:forward shock IC:e.g. Meszaros , et al. 94;Wang et al. 01;Granot & Guetta 03e.g. Meszaros & Rees 94;Dermer et al. 00;Zhang & Meszaros 01

# I will include...

- VHE photons from GRB early external shocks--very early afterglow phase
- VHE photons related to X-ray flares: X-ray flares photons are IC scattered by electrons in afterglow shocks

1. IC emission from very early external shocks

(Wang, Dai & Lu 2001 ApJ,556, 1010)

At deceleration radius, T\_obs~10-100 s Forward shock---Reverse shock structure is developed



## Energy spectra--- $\mathcal{V}f_{v} - \mathcal{V}$ (Wang, Dai & Lu 2001 ApJ,556, 1010)



At sub-GeV to GeV energies, the SSC of reverse shock is dominant; at higher energies, the Combined IC or SSC of forward shock becomes increasingly dominated

See also Wang, Dai & Lu 2001, ApJ,546, L33

### One interesting GeV burst



Gonzalez et al. 03: Hadronic model

Leptonic IC model:

Granot & Guetta 03 Pe'er & Waxman 04 Wang X Y et al. 05, A&A, 439,957

# 2. VHE photons from x-ray flare-blast wave interaction

# X-ray flares: late-time central engine activity

- ~30%-50% early afterglow have x-ray flares, Swift discovery
- Most striking flare GRB050502B discovered by Swift
- Flare light curves: rapid rise and decay
  - $\delta t/t$  <<1
- Afterglow decay consistent with a single power-law before and after the flare

X-ray flares occur inside the deceleration radius of the afterglow shock



IC between X-ray flare photons and afterglow electrons (Wang, Li & Meszaros 2006)

X-ray flare photons illuminate the afterglow shock electrons from inside



# Flare photons is the dominant cooling source of afterglow electrons

 Flare photons is the dominant cooling source for electrons in the afterglow jet

$$U'_X = D^2 F_X / (\Gamma^2 R^2 c) > B^2 / 8\pi$$

$$F_X > 10^{-10} \epsilon_{B,-2} E_{52} t_3^{-1} D_{28}^{-2} \,\mathrm{erg cm}^{-2} \mathrm{s}^{-1}$$

 Flare photons make the afterglow electrons fast-cooling

$$F_X > 10^{-10} \epsilon_{B,-2} E_{52} t_3^{-1} D_{28}^{-2} \,\mathrm{erg cm}^{-2} \mathrm{s}^{-1}$$
  
 $\gamma_m \gtrsim \gamma_c$ 

$$F_X > 3 \times 10^{-10} E_{52}^{1/2} \epsilon_{e,-1}^{-1} n_0^{-1/2} t_3^{-1/2} D_{28}^{-2} \mathrm{erg cm}^{-2} \mathrm{s}^{-1}$$



## IC GeV flare fluence-An estimate

So most energy of the newly shock electrons will be lost into IC emission

$$0.1 < \delta t/t < 1 \qquad \Psi = 10^{-7} - 10^{-6} \epsilon_{e,-1} E_{52} D_{28}^{-2} \mathrm{erg} \, \mathrm{cm}^{-2}$$
$$\varepsilon_{IC,p} \simeq 2\gamma_m^2 \varepsilon_X \simeq 3\epsilon_{e,-1}^2 E_{52}^{1/4} n_0^{-1/4} t_3^{-3/4} \left(\frac{\varepsilon_X}{1 \mathrm{keV}}\right) \mathrm{GeV}$$

X-ray flare peak energy

# Klein-Nishina suppression is unimportant below TeV

 $\nu F_{\nu} ~~1/2$  and -(p-2)/2 before and after the break at  $\varepsilon_{IC,p}$ 

Klein-Nishina cutoff  $\gamma_e < \gamma_{e,M} = \Gamma m_e c^2 / \varepsilon_X$   $\varepsilon_{IC,M} = 2\gamma_{e,M}^2 (\varepsilon_X/\Gamma)\Gamma = 0.4E_{52}^{1/4}n_0^{-1/4}t_3^{-3/4} \left(\frac{\varepsilon_X}{1 \text{ keV}}\right)^{-1} \text{ TeV}$  $\tau_{\gamma\gamma} \simeq 0.1\sigma_T \frac{U'_X}{\varepsilon_X/\Gamma} (\Gamma c \delta t) = 0.3F_{X,-9}n_0^{1/2}t_3^{1/2}E_{52}^{-1/2}D_{28}^2 \left(\frac{\delta t}{t}\right) \left(\frac{\varepsilon_X}{1 \text{ keV}}\right)^{-1}$ 

Magic, Milagro, HESS, etc.

But at high redshift, the infrared background absorption will be important

# Anisotropic IC scattering effect

Effect: decrease the power along the seed photon beam direction, but increase larger angle emission, about a half within 2/Γ (Wang, Li & Meszaros 2006)

e.g. Ghisellini 1978; Brunetti 2000; Fan & Piran 2006

- If  $\theta_j \gg 1/\Gamma$ , when integrated over angles, the received IC fluence will not be reduced.
- Suppose the sphere geometry: every direction has the same fluence.

### The IC fluence will not change

 $\theta_s$ 

## Temporal behavior of the IC emission

 Not correlated with the X-ray flare light curves. IC emission will be lengthened by the afterglow shock angular spreading time and the anisotropic IC effect



Delayed GeV emission ---GRB940217



### Summary

∫ IC in early external shocks ↓ IC of x-ray flares

Both are promising VHE photon sources for GLAST

- What could GLAST tell us?
- Origin of GeV photons (both prompt and delayed): Spectral and temporal properties
- > Maximum energy of the shock accelerated electrons :  $\mathcal{E}_{e,Max} > ?$
- Magnetic field in the shocks:  $L_{IC} / L_{syn} = U_{ph} / U_B$
- > ...