IDS Report GLAST SWG meeting Sept. 2, 2005

**Chuck Dermer (NRL)** 

GLAST grant supports costs of research with

- 1. Dr. Armen Atoyan (frequent visitor to NRL; travel to GLAST Symposium) UHECRs and v from GRBs and blazars, Chandra jets, Sgr A\*
- 2. Dr.Truong Le (NRL/NRC postdoctoral associate) Particle acceleration theory
- 3. Prof. Govind Menon (Summer Sabbatical Faculty and future vists) Black-hole research
- 4. Prof. Markus Böttcher (grad student support to Ohio U, visit to NRL) Microquasars
- 5. Jeremy Holmes (Summer hire) Cosmic ray propagation in the Galaxy

## Plans for Coming Year

- 1. GLAST Science Working Groups: GRBs (ISSS Course, L'Aquila, Italy) Solar System Blazars
- 2. New hire
- 3. Book: "High Energy Radiation from Black Holes: γ-rays, cosmic rays, and neutrinos"
- 4. HESS results "THE MULTI-MESSENGER APPROACH TO HIGH ENERGY GAMMA-RAY SOURCES" July 4-7, 2006 - Barcelona, Spain

## **Statistics of Gamma-Ray Blazars**



(C.D. and Stan Davis, 2001)
HESS obs. of high-z HBLs
GLAST studies of EBL;
internal absorption

HBLs, LBLs, and the EBL

• Peak of activity of cosmological  $\gamma$ -ray sources (blazars, gamma-ray bursts) at redshift  $z \sim 1$ 

• Population evolution is strongly non-Euclidean, so large number of sources near threshold for BL Lac objects and clusters of galaxies



#### **Energy Extraction from Rotating Black Holes**



Energy from Black Holes Rotating Black Holes Reducible vs. Irreducible mass Penrose process Blandford-Znajek process Force-free magnetosphere  $E \cdot J = 0,$  $\rho E + J \times B = 0.$ 

Maxwell's equation in 3+1 formalism (Komissarov 2004)

### Energy Extraction from Rotating Black Holes Analytic Solutions to the Constraint Equation for a Force-Free Magnetosphere around a Kerr Black Hole

Govind Menon<sup>1,2,3</sup> & Charles D. Dermer<sup>2</sup>

$$\frac{1}{2\Lambda} \frac{dH_{\varphi}^{2}}{d\Omega} = \frac{\alpha \gamma_{\varphi\varphi}}{\sqrt{\gamma}} [\Omega \partial_{r} (\frac{\Lambda}{\alpha \sqrt{\gamma}} (\gamma_{\varphi\varphi} \Omega + \beta_{\varphi}) \gamma_{\theta\theta} \Omega_{,r}) + \Omega \partial_{\theta} (\frac{\Lambda}{\alpha \sqrt{\gamma}} (\gamma_{\varphi\varphi} \Omega + \beta_{\varphi}) \gamma_{rr} \Omega_{\rho_{+}^{2}} + a^{2} \cos^{2} \theta + \partial_{r} (\frac{\Lambda}{\alpha \sqrt{\gamma}} (\beta^{2} - \alpha^{2} + \beta_{\varphi} \Omega) \gamma_{\theta\theta} \Omega_{,r}) + \partial_{\theta} (\frac{\Lambda}{\alpha \sqrt{\gamma}} (\beta^{2} - \alpha^{2} + \beta_{\varphi} \Omega) \gamma_{rr} \Omega_{,\theta})].$$

Constraint equation in 3+1 formalism

$$\Omega_{+} = \frac{a}{2Mr_{+} + \rho_{+}^{2}},$$

$$(r_{+} = M + \sqrt{M^{2} - a^{2}})$$

$$\rho_{+}^{2} = r_{+}^{2} + a^{2} \cos^{2} \theta$$

$$\Omega_{+} \rightarrow \frac{a}{8M^{2}} \quad \text{for a \ll M}$$

Generalizes monopole solution of BZ77 to  $a \rightarrow M$ 

 $\frac{d^2 \mathcal{E}}{dAdt} \approx \frac{a\Omega_H}{r^2} (\frac{B_0}{2})^2 \frac{\sin^2 \theta}{\rho_+^2}$ 

Microquasars as Gamma Ray Sources



# Gamma Rays from Jet Sources



Bread and butter physics

# Model for High Mass Microquasars







Fig. 1.— Geometry of the model. The direction of the radio jets defines the  $x_3$  axis. The orbital plane of the binary system is the  $(x_1, x_2)$  plane, defined in such a way that line of sight towards the observer lies in the  $(x_2, x_3)$  plane, where the azimuthal angle  $\phi = 0$ .

O6.5V, 23 M<sub>o</sub> primary T = 39000 K (3.5 eV) S =  $2.5 \times 10^{12}$  cm, i =  $25^{\circ}$ Period = 3.91 days

Claimed orbital variations of TeV radiation

Phase-dependent  $\gamma$ - $\gamma$  Opacity





Fig. 2.— Orbital modulation of the expected  $\gamma\gamma$  absorption trough, assuming a power spectrum with photon index  $\alpha_{ph} = 2.5$  and a photon production site at  $z_0 = 10^{12}$ The different curves represent the escaping photon spectrum at various orbital phases, i  $\phi_0 = 0$  (lowest curve) to  $\phi_0 = \pi$  (highest curve) in steps of  $\pi/10$ .

#### production region from the central compact object at phase $\phi_0 = 0$ . The figure illustrates that (1) VHE photons produced within a few $\times 10^{12}$ cm (i.e., of the order of the orbital separation of the binary system) would be subject to substantial $\gamma\gamma$ absorption; (2) the minimum of the absorption trough (maximum of $\tau_{\gamma\gamma}$ as a function of photon energy) is shifting towards higher energies for larger distances from the central source.

#### Spectral variations with phase

Opacity vs. location of gamma-ray production site

# Predictions for TeV Telescopes



Fig. 4.— Orbital modulation of the integrated photon number flux above energies  $E_0 = 250 \text{ GeV}$  (solid) and  $E_0 = 1 \text{ TeV}$  (short-dashed), the  $\gamma\gamma$  opacity at E = 250 GeV (dotted) and E = 1 TeV (dot-dashed), and the local photon spectral index  $\alpha_{500}$  at 500 GeV (long-dashed). As in Fig. 2, an underlying power-law of photon index  $\alpha_{\rm ph} = 2.5$  and a photon production site at  $z_0 = 10^{12}$  cm has been assumed. A periodic flux modulation is expected to be accompanied by positive spectral-index/flux correlation (spectral softening as the flux increases) at  $E_0 \gtrsim 300 \text{ GeV}$ .

Multiwavelength Spectrum of Microquasars



Spectral variations with phase both due to scattering kinematics and  $\gamma$ - $\gamma$  absorption

## Multiwavelength Emission from Sgr A\*



Very weak  $> 100 \text{ MeV} \gamma$ -ray emission

Second-order Fermi acceleration



With Peter Becker (GMU) and Truong Le





Predict GLAST detection of quasistationary Compton and bremsstrahlung fluxes from pc-scale plerion.

Propagation of GeV electrons power Sgr A West EGRET emission from young pulsar