



## A Physical Model for the Revised Blazar Sequence

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Synchrotron peak luminosity vs synchrotron peak frequency

 $48 \qquad @5 \text{ GHz} \\ 46 \qquad & & & & & \\ 44 \qquad & & & & & & \\ 42 \qquad & & & & & & \\ 42 \qquad & & & & & & \\ 42 \qquad & & & & & & \\ 40 \qquad & & & & & & \\ 12 \qquad 14 \qquad 16 \qquad 18 \\ \log(\nu_{\text{peak,sync}})$ 

5 GHz luminosity vs synchrotron peak frequency



#### **Previous work**



Peaks from fit to Swift / Planck data by Giommi et al. (2011).

"L" shape seen.

Note: y-axis is synchrotron peak + Compton peak.

But upper right may be filled in with BL Lacs with unknown redshift.

"L" (or "V") shape also seen by Nieppola et al. (2006), Meyer et al. (2011)









External radiation field for Compton scattering correlates with power injected into electrons.

As power increases, greater external radiation field leads to greater Compton scattering, and hence more Compton dominance.

At the same time, the greater scattering cools the electrons more, leading to a lower cooling break energy.

$$\gamma_c = \frac{3m_e c^2}{4c\sigma_{\rm T}(u'_B + u'_{sy} + u'_{ext})t_{esc}}$$

The peak frequency is directly related to this cooling break energy.



Ghisellini et al. (1998)



Abdo et al. (2010; CA: P. Giommi; M. Mazziotta; A. Tramacere) fit LBAS blazars to determine peak frequencies and luminosities:

$$\nu F_{\nu} = a \cdot \nu^3 + b \cdot \nu^2 + c \cdot \nu + d.$$

They provided empirical formulae for finding the peak frequency based on optical, radio, and X-ray data ( $\alpha_{ro}$ ,  $\alpha_{ox}$ ).

The 2LAC provides peak synchrotron frequency for sources with enough data using these formulae.

Can also use Abdo et al. (2010) empirical formula to calculate peak flux (normalized to 5 GHz flux density):

 $\operatorname{Log}(\nu_{peak_S}F(\nu_{peak_S})) = 0.5 \cdot \operatorname{Log}(\nu_{peak_S}) - 20.4 + 0.9 \cdot \operatorname{Log}(\operatorname{R}_{5\mathrm{GHz}}),$ 

Can use this to create the blazar sequence from the 2LAC.





Simple explanation: as synchrotron bump moves to higher frequencies, radio flux will decrease (e.g., Lister et al. 2011).



A physical effect, or a result of the way the peak frequency is calculated?



- ~ 350 sources
- "V" shape seen



#### **Correlations**



TABLE 1 Statistics of correlations involving  $\nu_{pk}^{sy}$ .

Sample	ρ	$PNC(\rho)$	au	$PNC(\tau)$
$L_{pk}^{sy}$ versus $\nu_{pk}^{sy}$				
BL Lacs	-0.19	0.019	-0.12	0.032
FSRQs	-0.12	0.073	-0.088	0.066
All sources with known $z$	-0.56	$2.1 \times 10^{-30}$	-0.37	0.00

Spearman and Kendall Rank Coefficients

Can objects with unknown z ruin this anti-correlation?





Fitting the high-energy comonents of blazars in the LBAS sample, Abdo et al. (2010) found a relationship between the LAT spectral index and peak frequency of the Compton component:

$$Log(\nu_{\rm peak}^{IC}) = -4.0 \cdot \Gamma + 31.6$$

This can be used to calculate the peak Compton frequency for the 2LAC sample.





## **High-Energy Component**



For the LBAS,  $L_{pk}^{C}$  from the fit (Abdo et al. 2010) and from the extrapolation:



Once the peak frequency is known, the power law can be extrapolated to find the peak Compton luminosity.

~10% of the 350 sources in our sample are also in the 58 month *Swift*-BAT catalog. For these objects their BAT spectra were also used to constrain the Compton peak.



#### **Compton Dominance**



Compton dominance: definitely an anticorrelation, and an "L" shape.

Compton dominance does not depend on redshift.

Sources with unknown redshifts are also plotted.





Objects with unknown redshift do not destroy the correlation!



# **A Little Theory**



Inject electrons as power law between two electron energies:

$$Q_e(\gamma) = Q_0 \gamma^{-q} H(\gamma; \gamma_1, \gamma_2)$$
.

Equilibrium solution, injection balanced with escape and injection slow cooling solution,  $\gamma_1 < \gamma_c$ :

$$N_e(\gamma) \approx Q_0 t_{esc} \begin{cases} (\gamma/\gamma_c)^{-q} & \gamma_1 < \gamma < \gamma_c \\ (\gamma/\gamma_c)^{-q-1} & \gamma_c < \gamma < \gamma_2 \end{cases}$$

$$\gamma_c = \frac{3m_ec^2}{4c\sigma_{\rm T}(u_B' + u_{sy}' + u_{ext}')t_{esc}}$$

fast cooling solution,  $\gamma_c < \gamma_1$  :

$$N_e(\gamma) \approx Q_0 t_{esc} \begin{cases} (\gamma/\gamma_1)^{-2} & \gamma_c < \gamma < \gamma_1 \\ (\gamma/\gamma_1)^{-q-1} & \gamma_1 < \gamma < \gamma_2 \end{cases}$$

Slow cooling, peak associated with cooling break



Fast cooling, peak associated with minimum injection break



See, e.g., Boettcher & Dermer (2002) 14





Peak frequency is associated with max( $\gamma_c,$   $\gamma_1).$ 

 $\gamma_{c}$  depends on power, but  $\gamma_{1}$  does not, presumably.

Once  $\gamma_c$  is less than  $\gamma_1$ , synchrotron peak luminosity will be roughly independent of peak frequency.

$$\epsilon_{pk} = \frac{h\nu_{pk}}{m_e c^2} = \delta_{\rm D} \epsilon_B \begin{cases} \gamma_c^2 & \gamma_1 < \gamma_c \\ \gamma_1^2 & \gamma_c < \gamma_1 \end{cases}$$

$$L_{pk}^{sy} = L_{\epsilon_{pk}}^{sy} = \frac{2\delta_{\rm D}^4}{3} c\sigma_{\rm T} u_B' Q_0 t_{esc} \begin{cases} \gamma_c^3 & \gamma_1 < \gamma_c \\ \gamma_1^3 & \gamma_c < \gamma_1 \end{cases}$$

Scale injected electrons:

 $Q_0 = Q_{00}\ell$ 

Scale injected Lorentz factor with power:

 $\Gamma = \Gamma_0 \ell^g$ 

Scale injected magnetic field with power:

$$B = B_0 \ell^b$$

Scale external radiation field with power:

$$u_{ext} = u_0 \ell^a$$

So we assume blazars are two parameter engines:  $\theta$ ,  $\ell$ 





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A clever choice of parameters can reproduce the "V" shape in the  $L_{pk}\mathchar`-\nu_{pk}$  diagram







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## **Reproducing Blazar Sequence**



It can also reproduce the "L" shape on the  $A_C^{-\nu_{pk}}$  diagram





## **Reproducing Blazar Sequence**



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- Blazar sequence generated from the 2LAC (L<sub>pk</sub>- $v_{pk}$  and A<sub>C</sub>- $v_{pk}$ ).
- Largest sample yet for Compton Dominance plot.
- Blazars with unknown z could ruin  $L_{pk}\text{-}\nu_{pk}$  anti-correlation, but not  $A_C\text{-}\nu_{pk}$  anti-correlation.
- Standard cooling scenario seems to explain "V" and "L" shapes on  $L_{pk} v_{pk}$  and  $A_C v_{pk}$  diagrams if  $v_{pk}$  is associated with the maximum of  $\gamma_c$  and  $\gamma_1$ .





### **Extra slides**



#### **Gamma-ray Component**



Fossati et al. (1998)



#### **Other works**





Meyer et al. (2011)

Meyer et al. (2011) "V" shape (or "L" shape)

Intermediate SED shapes don't appear common.

Meyer et al. explain this as due to viewing angle effects in stratified jets.

Other recent works have similar "V" or "L" shape: e.g., Nieppola et al. (2006), Giommi et al. (2011).

But upper right may be filled in with BL Lacs with unknown redshift (Giommi et al. 2011).





Nieppola et al. (2006) fit BL Lacs with logparabola to locate synchrotron peaks

They also found "V" or "L" shape.

Problems with fitting with log-parabola versus third-order polynomial?

Also, no FSRQs included.



Nieppola et al. (2006)





