

Maximum Likelihood:

Statistics in photon counting experiments

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Questions in γ -ray astronomy

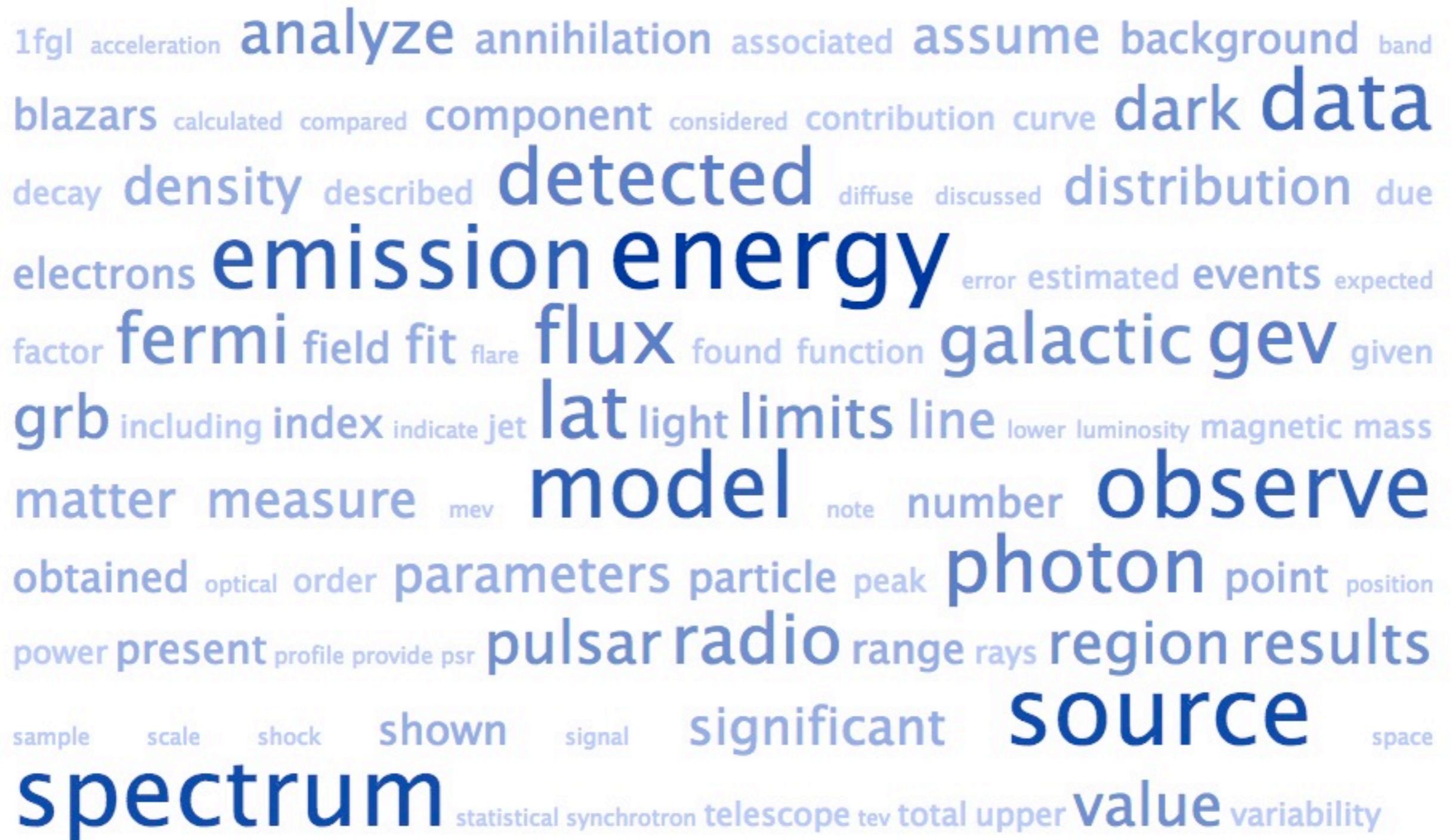
- Is a source significantly detected?
- If so, what is its flux?
- If not, what is upper limit on the flux?
- What kind of spectrum does it have?
- What is its spectral index?
- What is its location in the sky?
- What are the errors on these values?
- Is the source variable?

Questions in DM astrophysics

- Does Fermi detect γ -ray line emission from DM particle annihilation?
- With what significance?
- What is the energy of the line?
- What is the measurement error?
- What is the spatial distribution?
- What kind of systematic errors may be present?

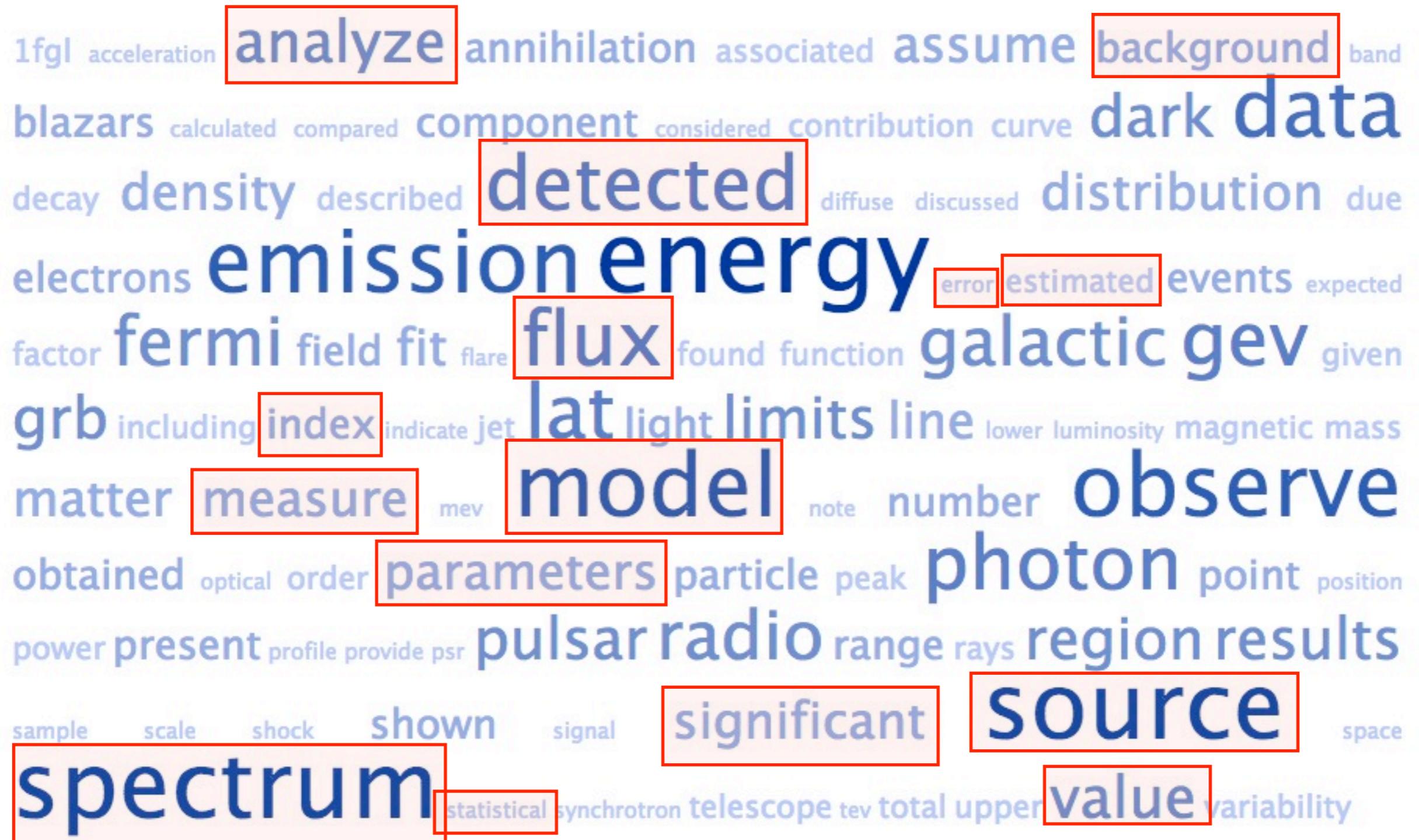
LAT paper tag cloud

100 most frequently used words from 875 papers mentioning LAT and γ -ray on arXiv



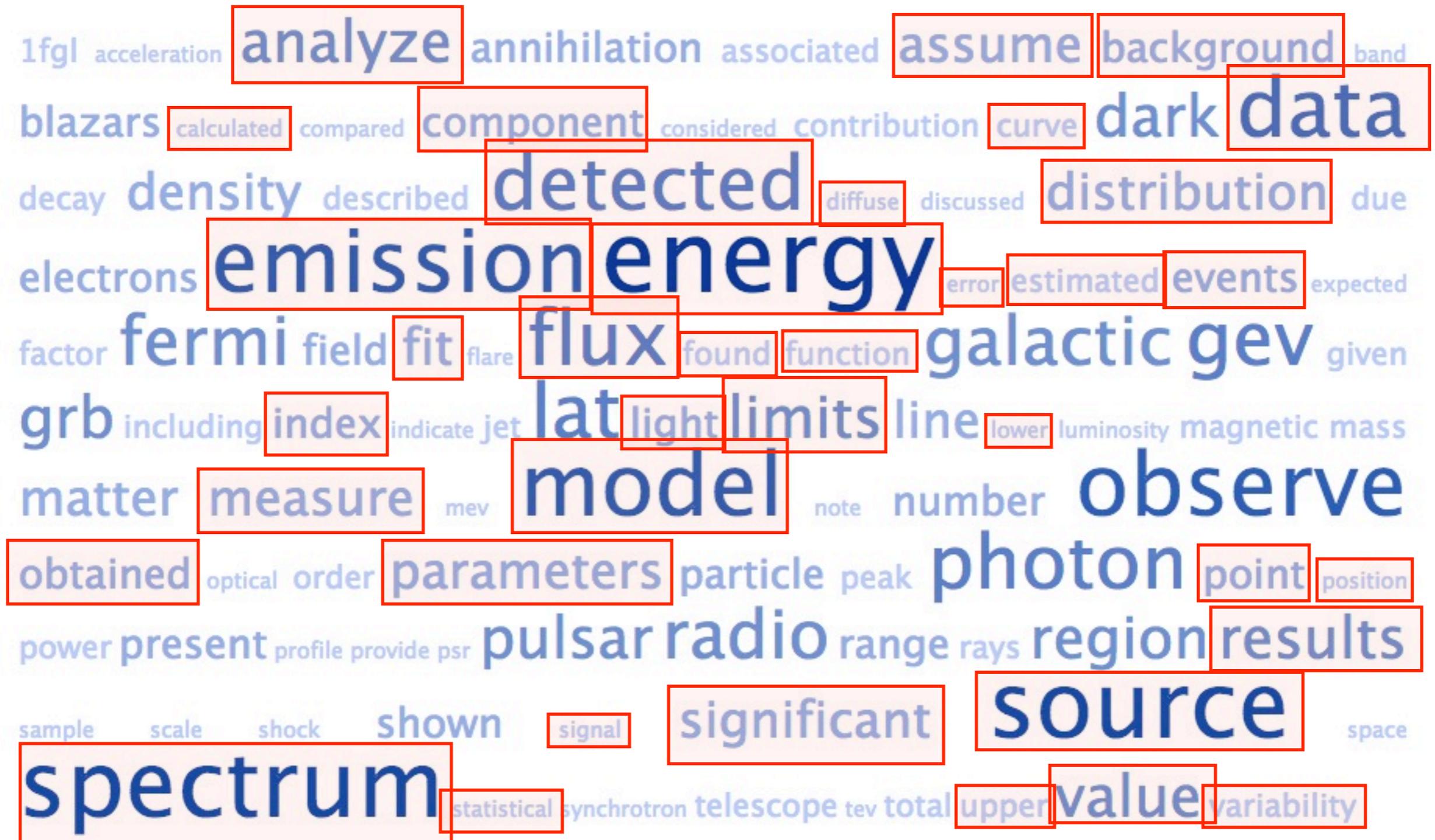
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Questions in γ -ray astronomy

Hypothesis testing

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Maximum likelihood technique

- Given a set of observed data:
- ... produce a model that *accurately* describes the data, including parameters that we wish to estimate,
- ... derive the probability (density) for the data given the model (PDF),
- ... treat this as a function of the model parameters (likelihood function), and
- ... maximize the likelihood with respect to the parameters - ML estimation.

Maximum likelihood basics

- Data: $X = \{x_i\} = \{x_1, x_2, \dots, x_N\}$
- Model parameters: $\Theta = \{\theta_j\} = \{\theta_1, \theta_2, \dots, \theta_M\}$
- Likelihood: $\mathcal{L}(\Theta|X) = P(X|\Theta)$
- Conditional probability rule for independent events: $P(A, B) = P(A)P(B|A) = P(A)P(B)$
CPR Independence
- For independent data:
$$P(X|\Theta) = P(\{x_i\}|\Theta) = P(x_1|\Theta)P(x_2, \dots, x_N|\Theta) = \dots$$
$$= P(x_1|\Theta)P(x_2|\Theta) \dots P(x_N|\Theta) = \prod_i P(x_i|\Theta)$$

$$\mathcal{L}(\Theta|X) = \prod_i P(x_i|\Theta)$$

ML estimation (MLE)

- Parameters can be estimated by maximizing likelihood. Easier to work with log-likelihood:

$$\ln \mathcal{L}(\Theta) = \ln \mathcal{L}(\Theta|X) = \sum_i \ln P(x_i|\Theta)$$

- Estimates of $\{\hat{\theta}_k\}$ from solving simultaneous equations:

$$\left. \frac{\partial \ln \mathcal{L}}{\partial \theta_j} \right|_{\{\hat{\theta}_k\}} = 0$$

- For one parameter, if we have: $\mathcal{L}(\theta) \sim e^{-\frac{(\theta-\hat{\theta})^2}{2\sigma_\theta^2}}$

then: $\left. \frac{\partial^2 \ln \mathcal{L}}{\partial \theta^2} \right|_{\hat{\theta}} = -\frac{1}{\sigma_\theta^2}$

Gaussian approximation

so 2nd derivative is related to “errors”

Why maximum likelihood...

...rather than some ad-hoc estimation method?

- ML framework provides a “cookbook” through which problems can be solved.
In other methods ad-hoc choices may have to be made.
- ML provides unbiased, minimum variance estimate as sample size increases.
Same may not be case for ad-hoc methods.
- Asymptotically Gaussian: evaluation of confidence bounds & hypothesis testing.
- Well studied in the literature.
- Starting point for Bayesian analysis.

χ^2 fit of constant - I

- Data: independent measurements of flux of some source with errors - (x_i, σ_i)
- Model: all measurements are of a constant flux F with Gaussian errors.

- Probabilities: $P(x_i|F) = \frac{1}{\sqrt{2\pi}\sigma_i} e^{-\frac{(x_i-F)^2}{2\sigma_i^2}}$

- Log likelihood:

$$\ln \mathcal{L}(F) = - \sum \frac{(x_i - F)^2}{2\sigma_i^2} - \sum \ln \sigma_i - \frac{N}{2} \ln 2\pi$$

χ^2 fit of constant - II

- Log likelihood:

$$\ln \mathcal{L}(F) = - \sum \frac{(x_i - F)^2}{2\sigma_i^2} \quad \text{Constant with respect to } F$$

Constant with respect to F

$$- \sum \ln \sigma_i - \frac{N}{2} \ln 2\pi$$

- Maximize for MLE of F :

$$\frac{\partial \ln \mathcal{L}}{\partial F} = \sum \frac{x_i - F}{\sigma_i^2} = 0 \implies \hat{F} = \frac{\sum x_i / \sigma_i^2}{\sum 1 / \sigma_i^2}$$

- Curvature gives “error” on F :

$$\frac{1}{\sigma_F^2} = - \left. \frac{\partial^2 \ln \mathcal{L}}{\partial F^2} \right|_{\hat{F}} = \sum \frac{1}{\sigma_i^2} \implies \sigma_F = \frac{1}{\sqrt{\sum 1 / \sigma_i^2}}$$

Example: combining results

Combined EBL measurements from Sanchez et al. (2013)

The spectral break predicted using the model of Fra08 is in good agreement with the data; the mean scaling factor is $\langle\alpha\rangle = 0.85 \pm 0.10$

Similar results were found by Ackermann et al. (2012) and Abramowski et al. (2013), who modeled the EBL-absorbed spectra of AGNs detected in the HE and VHE regimes respectively, and found scaling factors of $\alpha_{Fermi} = 1.02 \pm 0.23$ and $\alpha_{HESS} = 1.27^{+0.18}_{-0.15}$.

Taking only the statistical errors, a χ^2 fit to the HESS, *Fermi* and our results gives a mean value of $\alpha_{combined} \approx 0.98$ and a value of $\chi^2 = 5.45$ for two degrees of freedom, compatible with the hypothesis that the values are consistent at the 1.85

Event counting experiment

- Experiment detects n events (e.g. γ rays)
- Model: Poisson process with mean of λ :

$$P(x|\theta) \rightarrow P(n|\lambda) = \frac{\lambda^n e^{-\lambda}}{n!}$$

- Log likelihood: $\ln \mathcal{L}(\lambda) = n \ln \lambda - \lambda - \ln n!$
- ML estimate and error in Gaussian regime:

$$\frac{\partial \ln \mathcal{L}}{\partial \lambda} = \frac{n}{\lambda} - 1 \implies \hat{\lambda} = n$$

$$\frac{1}{\sigma_\lambda^2} = - \left. \frac{\partial^2 \ln \mathcal{L}}{\partial \lambda^2} \right|_{\hat{\lambda}} = \frac{n}{\hat{\lambda}^2} \implies \sigma_\lambda^2 = n$$

Gaussian approximation

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Constant WRT λ

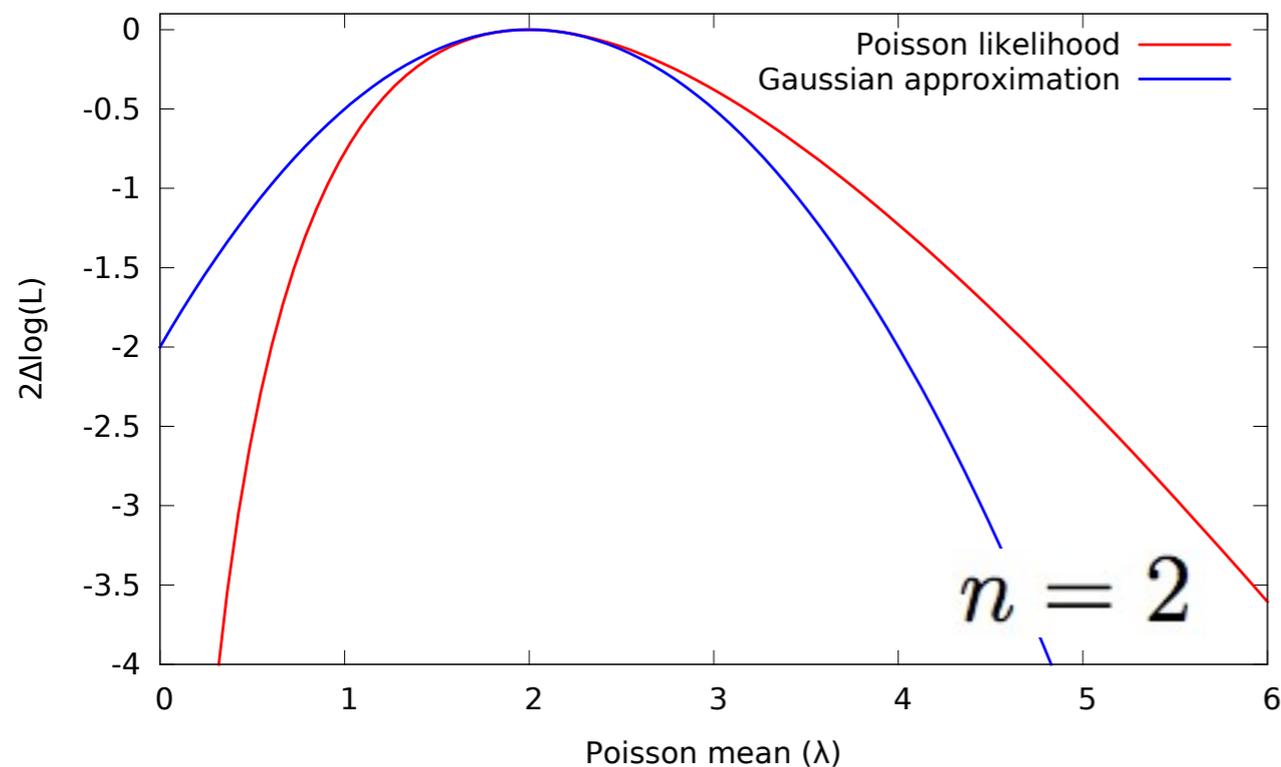
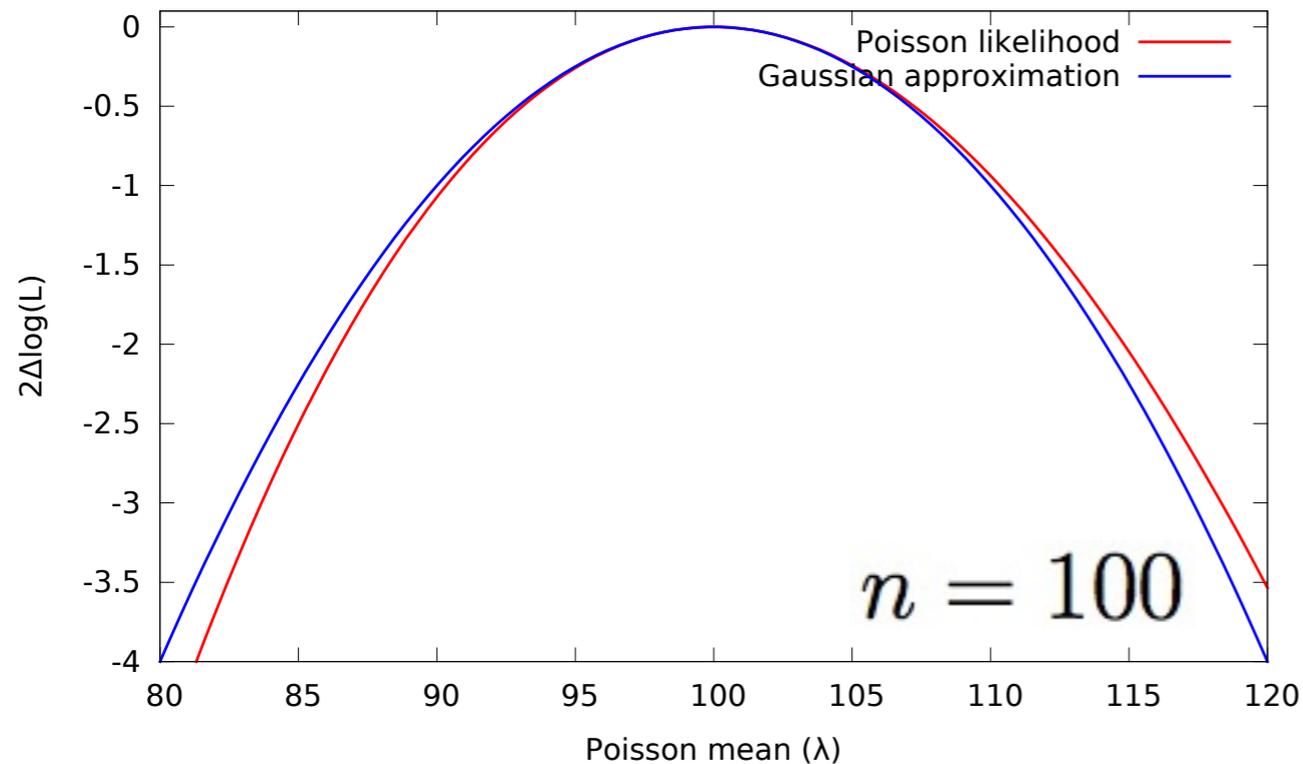
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Data cpt Npred
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Gaussian approximation

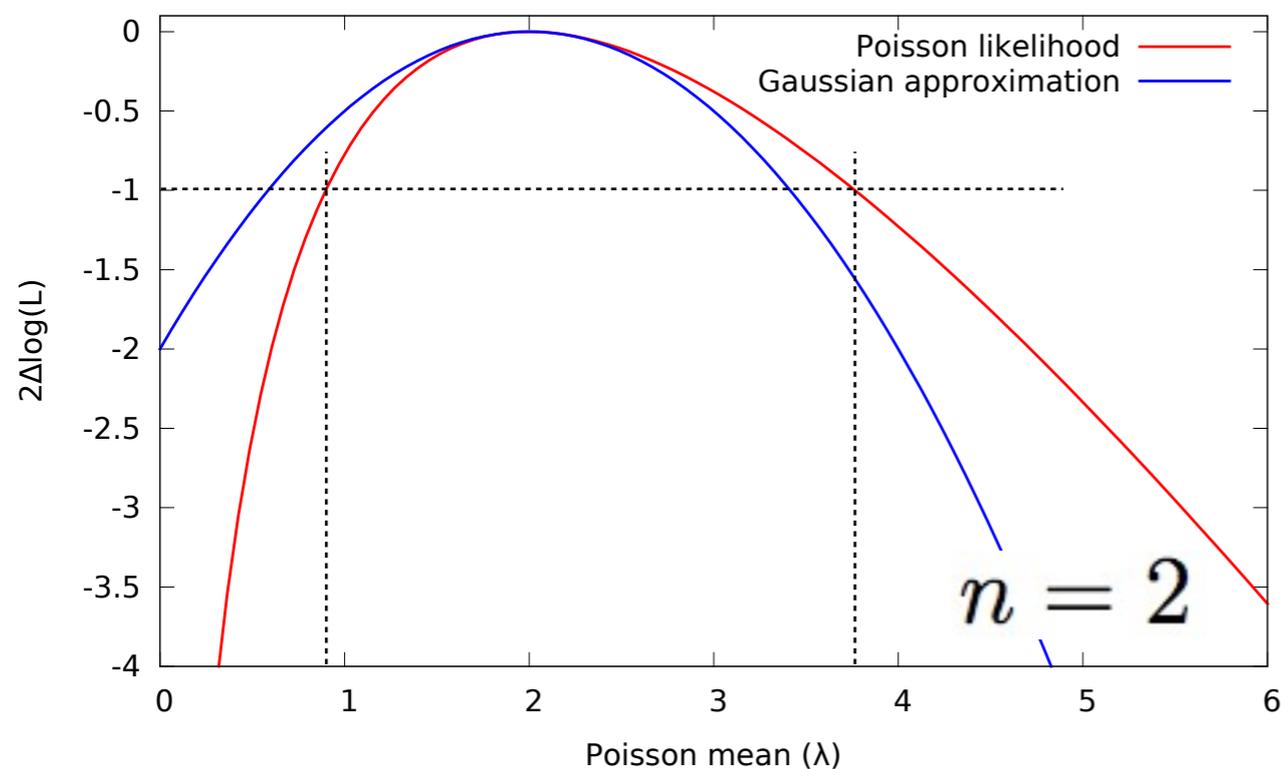
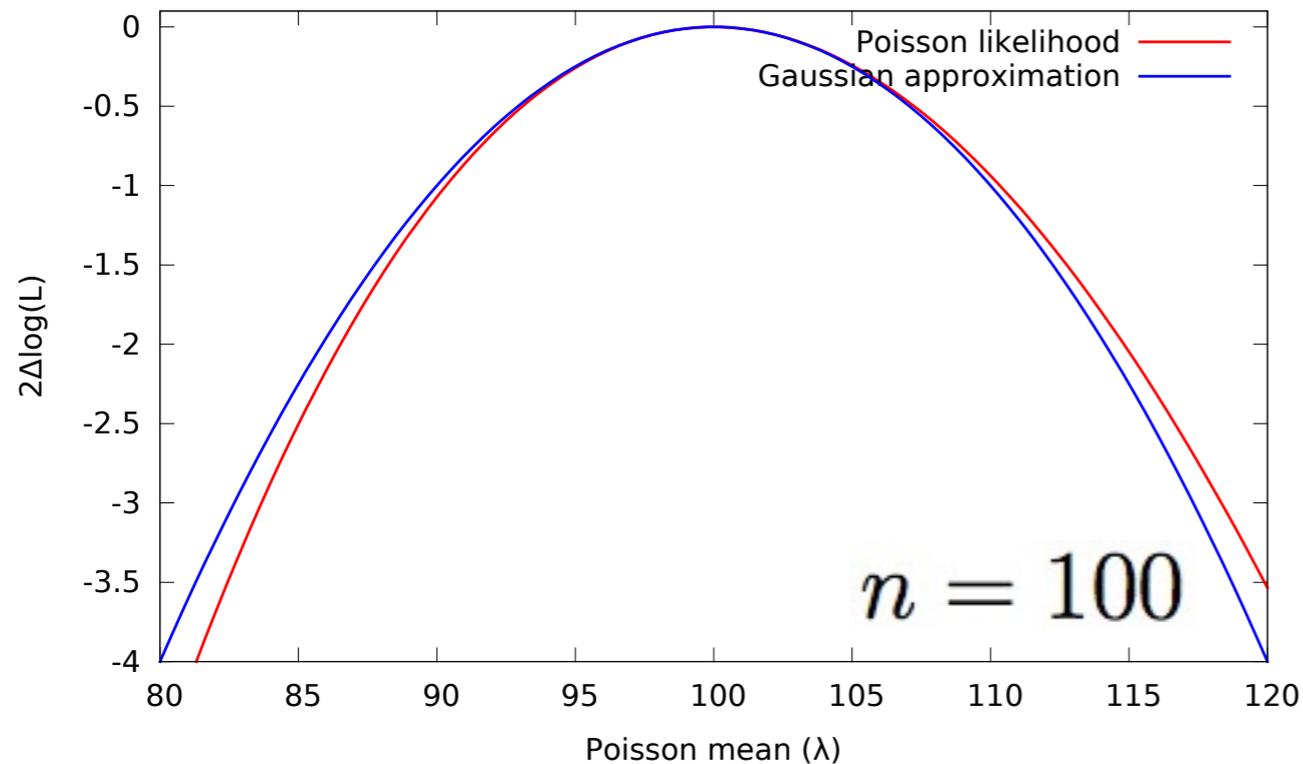
Log-likelihood profiles



- Gaussian approximation is reasonable when n is “large enough”. In this case $\sigma_\lambda^2 = n$ is a good estimate of the “error”.
- If not, estimate errors by finding points where $2 \ln \mathcal{L}(\lambda)$ decreases by 1.0 from maximum, i.e.,

$$2 \ln \mathcal{L}(\lambda) = 2 \ln \mathcal{L}(\hat{\lambda}) - 1$$
- $n=100$: $\hat{\lambda} = 100.0^{+10.33}_{-9.67}$
- $n=2$: $\hat{\lambda} = 2.0^{+1.77}_{-1.10}$

Log-likelihood profiles



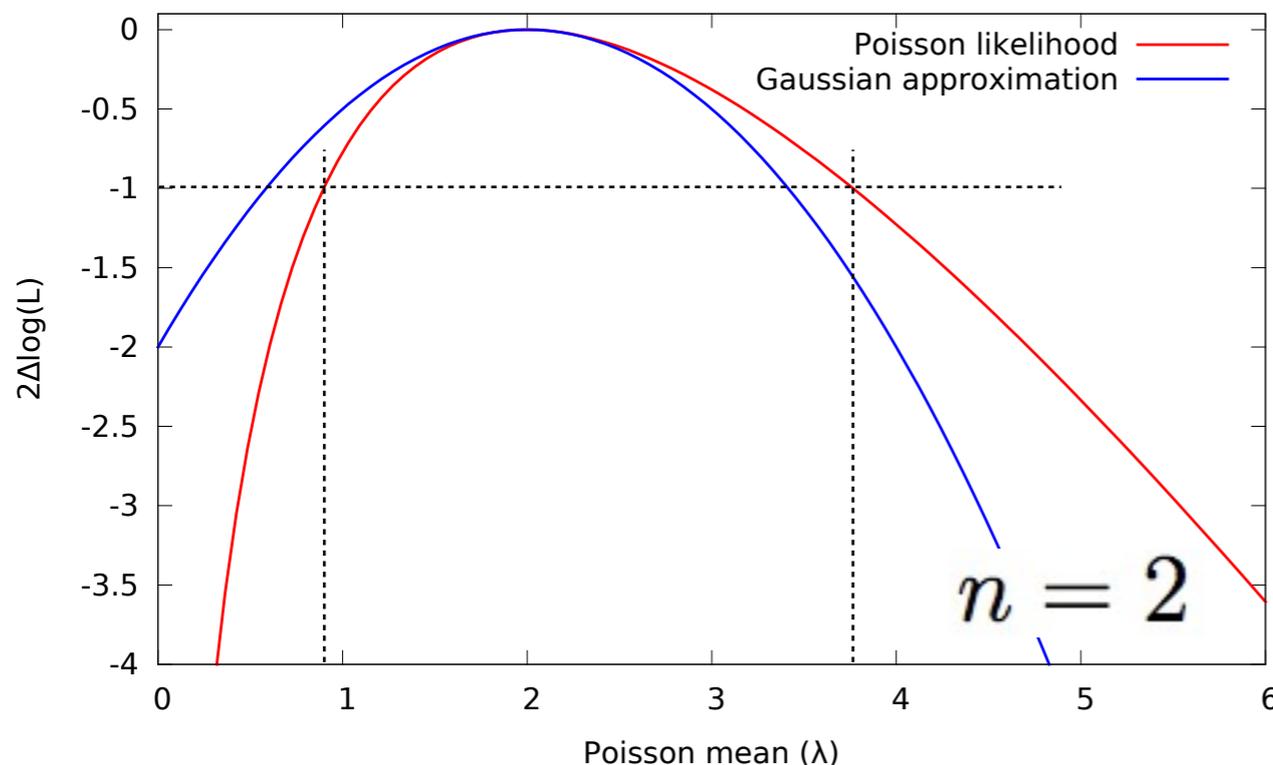
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MLE example 2:

Log-likelihood profiles

```
# errors_poisson.py - 2013-05-07 SJF
# Evaluate the errors on the Poisson mean
import math, scipy.optimize
n_meas      = 2
logL        = lambda lam: n_meas*math.log(lam) - lam
opt_fn      = lambda lam: -logL(lam)
opt_res     = scipy.optimize.minimize(opt_fn, 1e-8)
lam_est     = opt_res.x[0]
logL_max    = logL(lam_est)
root_fn     = lambda lam: 2.0*(logL(lam) - logL_max) + 1.0
lam_lo      = scipy.optimize.brentq(root_fn, 1e-8, lam_est)
lam_hi      = scipy.optimize.brentq(root_fn, lam_est, 1e8)
print lam_est, lam_lo - lam_est, lam_hi - lam_est
```



$2 \ln \mathcal{L}(\lambda)$ decreases by 1.0 from maximum, i.e.,

$$2 \ln \mathcal{L}(\lambda) = 2 \ln \mathcal{L}(\hat{\lambda}) - 1$$

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Hypothesis testing

- Compare likelihoods of two hypotheses to see which is better supported by the data.
- Likelihood-ratio test (LRT) & Wilks' theorem.

- Given a model with $N+M$ parameters:

$$\Theta = \{\theta_1, \dots, \theta_N, \theta_{N+1}, \dots, \theta_{N+M}\}$$

where N have true values: $\theta_1^T, \dots, \theta_N^T$

- Values of likelihood under two hypotheses:

$$\mathcal{L}_1 = \mathcal{L}(\hat{\theta}_1, \dots, \hat{\theta}_N, \hat{\theta}_{N+1}, \dots, \hat{\theta}_{N+M})$$

$$\mathcal{L}_0 = \mathcal{L}(\theta_1^T, \dots, \theta_N^T, \hat{\theta}_{N+1}, \dots, \hat{\theta}_{N+M})$$

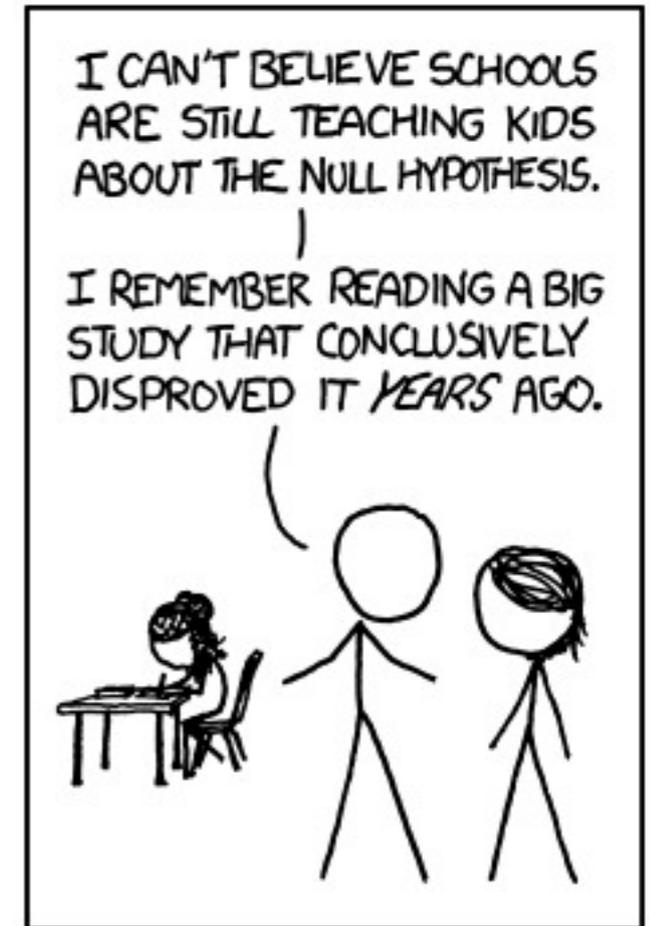
- “Ratio” distributed as: $2(\ln \mathcal{L}_1 - \ln \mathcal{L}_0) \sim \chi^2(N)$

Terms and conditions apply

Why is that useful?

(We don't know the true values of any parameters!)

- We make an assumption about the model (*the null hypothesis*), in which the parameters have some presumed “true” values.
- Compute \mathcal{L}_0 from these values and \mathcal{L}_1 using MLE for all params.
- Hope to show that $2(\ln \mathcal{L}_1 - \ln \mathcal{L}_0)$ is so large that it is improbable from $\chi^2(N)$,
- and, hence, reject the null hypothesis. Usually cannot say hypothesis is true!



<http://xkcd.com/892/>

Source & Background

- Data: events detected in two independent “channels”: $X = \{n_1, n_2\}$

- Model: Poisson process with...

- Unknown “source” and “background”:

$$\Theta = \{\theta_1, \theta_2\} = \{S, B\} \quad \vec{\Theta} = \begin{pmatrix} \theta_1 \\ \theta_2 \end{pmatrix} = \begin{pmatrix} S \\ B \end{pmatrix}$$

- Response matrix (presumed known)

$$\mathbf{R} = \begin{pmatrix} r_{11} & r_{12} \\ r_{21} & r_{22} \end{pmatrix}$$

- Poisson means:

$$\vec{\lambda} = \mathbf{R}\vec{\Theta} \quad \begin{pmatrix} \lambda_1 \\ \lambda_2 \end{pmatrix} = \begin{pmatrix} r_{11} & r_{12} \\ r_{21} & r_{22} \end{pmatrix} \begin{pmatrix} S \\ B \end{pmatrix}$$

MLE

- Log likelihood:

$$\ln \mathcal{L}(S, B) = n_1 \ln(r_{11}S + r_{12}B) + n_2 \ln(r_{21}S + r_{22}B) \\ - (r_{11} + r_{21})S - (r_{12} + r_{22})B + \text{const}$$

- MLE: $\frac{\partial \ln \mathcal{L}}{\partial S} = \frac{\partial \ln \mathcal{L}}{\partial B} = 0 \implies \hat{\vec{\Theta}} = \mathbf{R}^{-1} \vec{n}$

$$\begin{pmatrix} \hat{S} \\ \hat{B} \end{pmatrix} = \frac{1}{r_{11}r_{22} - r_{12}r_{21}} \begin{pmatrix} r_{22} & -r_{12} \\ -r_{21} & r_{11} \end{pmatrix} \begin{pmatrix} n_1 \\ n_2 \end{pmatrix}$$

$$\ln \mathcal{L}_1 = \ln \mathcal{L}(\hat{S}, \hat{B}) = n_1 \ln n_1 + n_2 \ln n_2 - (n_1 + n_2)$$

- If likelihood: $\mathcal{L}(\vec{\Theta}) \sim e^{-\frac{1}{2}(\vec{\Theta} - \hat{\vec{\Theta}})^T \mathbf{\Sigma}^{-1}(\vec{\Theta} - \hat{\vec{\Theta}})}$

Gaussian approximation

“errors” are: $\left. \frac{\partial^2 \ln \mathcal{L}}{\partial \theta_i \partial \theta_j} \right|_{\hat{\vec{\Theta}}} = -(\mathbf{\Sigma}^{-1})_{ij} = -\mathcal{I}_{ij}$

MLE

- Log likelihood:

$$\ln \mathcal{L}(S, B) = \overbrace{n_1 \ln(r_{11}S + r_{12}B) + n_2 \ln(r_{21}S + r_{22}B)}^{\text{Data component}} - \underbrace{(r_{11} + r_{21})S - (r_{12} + r_{22})B}_{\text{Npred}} + \cancel{\text{const}}$$

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↑
↑

Covariance matrix

Fisher information matrix

Covariances and errors

- Calculate Fisher information matrix and invert:

$$\mathcal{I}_{ij} = - \left. \frac{\partial^2 \ln \mathcal{L}}{\partial \theta_i \partial \theta_j} \right|_{\hat{\Theta}} \rightarrow \Sigma = \begin{pmatrix} \sigma_S^2 & \text{cov}(S, B) \\ \text{cov}(S, B) & \sigma_B^2 \end{pmatrix} = \mathcal{I}^{-1}$$

- For our example we get:

$$\mathcal{I} = \frac{1}{n_1 n_2} \begin{pmatrix} r_{21}^2 n_1 + r_{11}^2 n_2 & r_{21} r_{22} n_1 + r_{11} r_{12} n_2 \\ r_{21} r_{22} n_1 + r_{11} r_{12} n_2 & r_{22}^2 n_1 + r_{12}^2 n_2 \end{pmatrix}$$

$$\Sigma = \frac{1}{\det(\mathbf{R})^2} \begin{pmatrix} r_{22}^2 n_1 + r_{12}^2 n_2 & -r_{21} r_{22} n_1 - r_{11} r_{12} n_2 \\ -r_{21} r_{22} n_1 - r_{11} r_{12} n_2 & r_{21}^2 n_1 + r_{11}^2 n_2 \end{pmatrix}$$

- In general parameters are correlated, but can choose set that is uncorrelated. Here they are $\{\lambda_1, \lambda_2\}$ giving $\hat{\lambda}_1 = n_1, \hat{\lambda}_2 = n_2, \Sigma_\lambda = \text{diag}(n_1, n_2)$

Source significance

- Null hypothesis: suppose $S = 0$, then:

$$\begin{aligned}\ln \mathcal{L}_0(B) &= \ln \mathcal{L}(S = 0, B) \\ &= n_1 \ln r_{12}B + n_2 \ln r_{22}B - (r_{12} + r_{22})B\end{aligned}$$

- MLE for B gives: $\frac{\partial \ln \mathcal{L}_0}{\partial B} = 0 \implies \hat{B}_0 = \frac{n_1 + n_2}{r_{12} + r_{22}}$

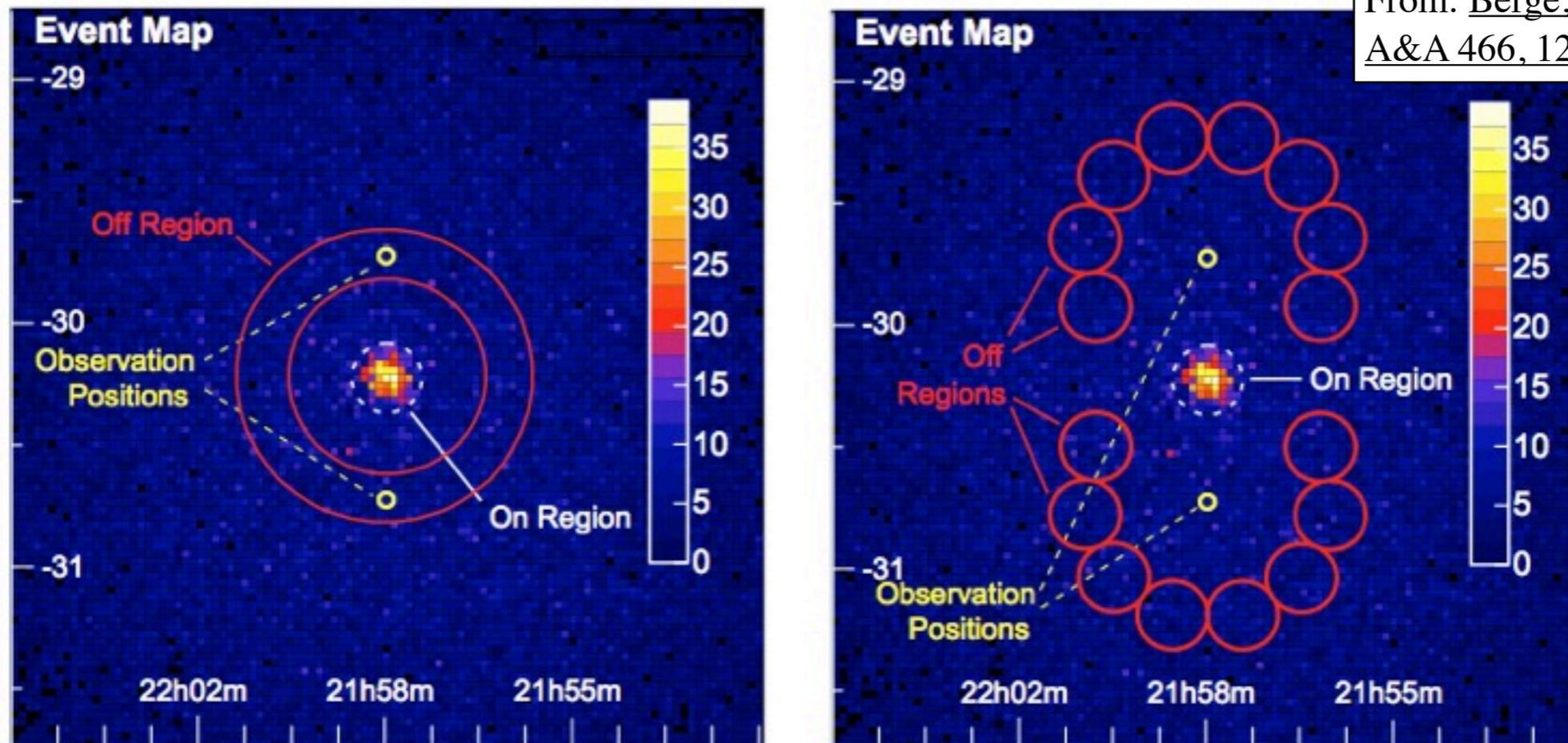
$$\begin{aligned}\ln \mathcal{L}_0 &= \ln \mathcal{L}_0(\hat{B}_0) \\ &= n_1 \ln \frac{r_{12}(n_1 + n_2)}{r_{12} + r_{22}} + n_2 \ln \frac{r_{22}(n_1 + n_2)}{r_{12} + r_{22}} - (n_1 + n_2)\end{aligned}$$

- Test statistic: $TS = 2(\ln \mathcal{L}_1 - \ln \mathcal{L}_0) \sim \chi^2(1)$

$$TS = 2 \left[n_1 \ln \frac{(r_{12} + r_{22})n_1}{r_{12}(n_1 + n_2)} + n_2 \ln \frac{(r_{12} + r_{22})n_2}{r_{22}(n_1 + n_2)} \right]$$

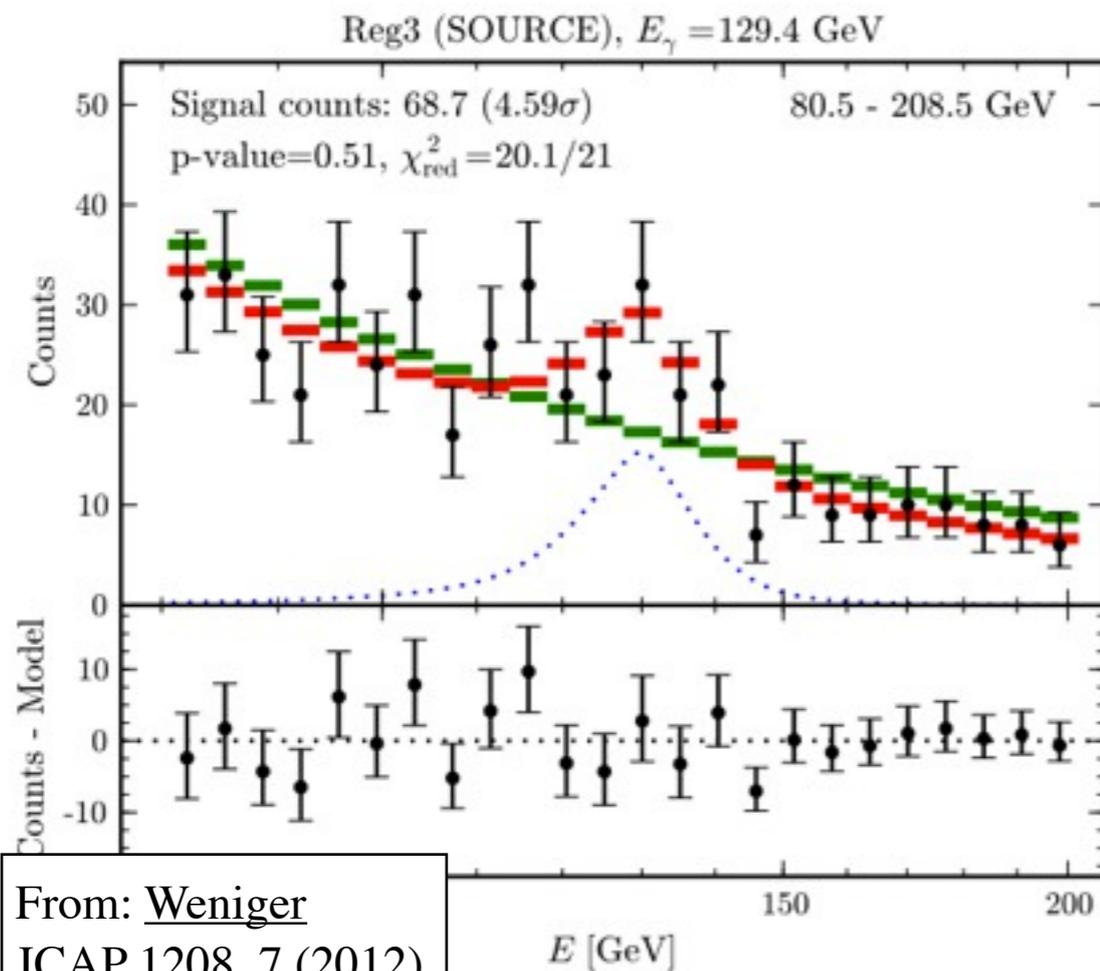
On/Off problems

From: [Berge, Funk, Hinton](#)
A&A 466, 1219–1229 (2007)

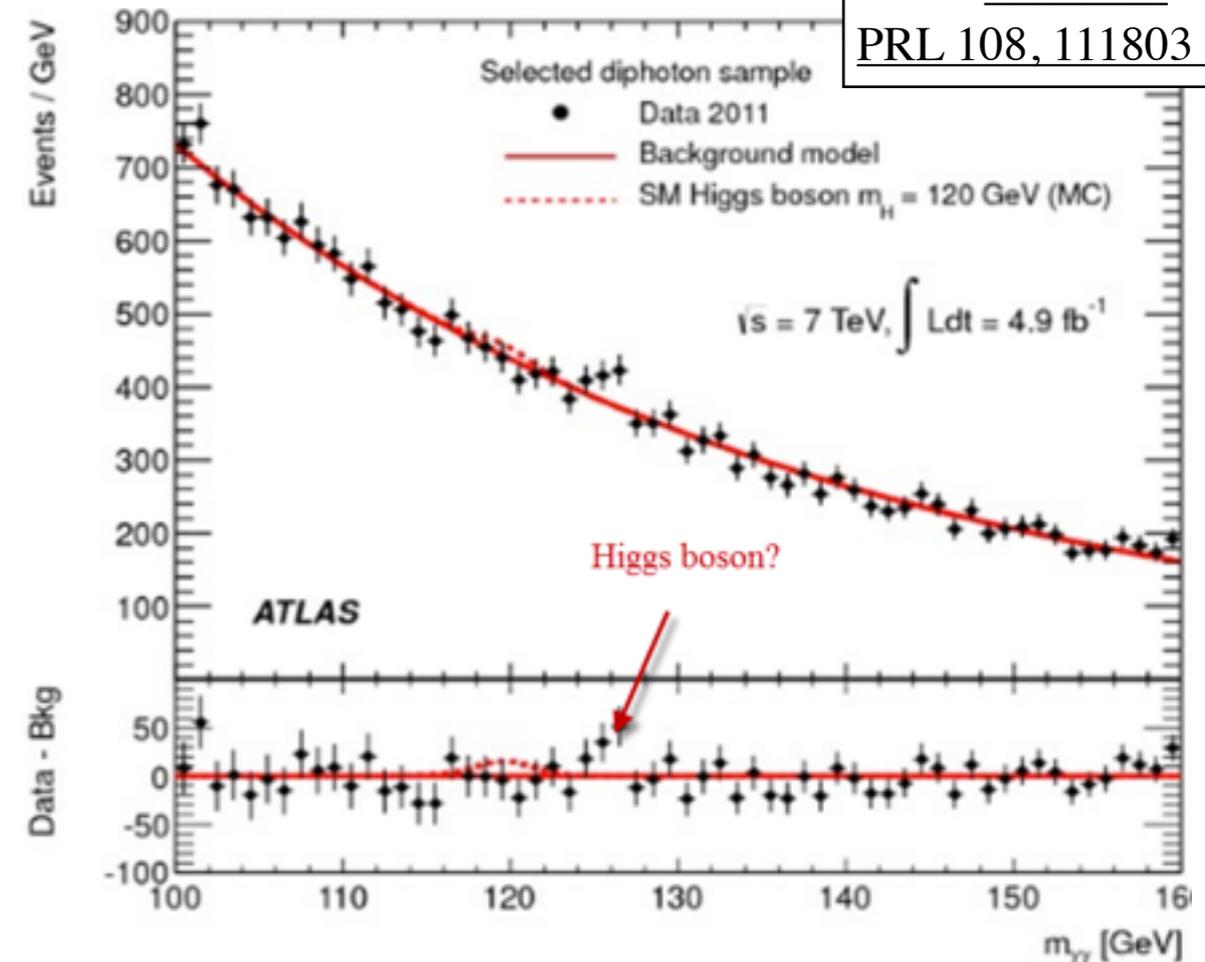


- VHE astronomy - gamma-ray sources and a background of cosmic rays.
- Problem - to evaluate flux of source and its statistical significance. Define on-source (source+background) and off-source (background) channels.

On/Off problems



From: Weniger
 JCAP 1208, 7 (2012)



From: Aad et al.
 PRL 108, 111803 (2012)

- Line searches - DM with Fermi, or Higgs with ATLAS.
- Problem - detect line signal on top of spectrum of background events. Define “on-source” and “off-source” regions. Must assume that spectrum of background is known or calculable.

On/Off problems

- General set of problems where:

$$n_2 \rightarrow n_{off}$$

$$n_1 \rightarrow n_{on}$$

$$\lambda_2 \rightarrow \lambda_{off} = BT$$

$$\lambda_1 \rightarrow \lambda_{on} = (S + \alpha B)T$$

- and where these are assumed to be known:
 - α - on to off-source background ratio
 - T - observation time (or other detector factors)

MLE for On/Off problems

- Then: $\mathbf{R} = T \begin{pmatrix} 1 & \alpha \\ 0 & 1 \end{pmatrix}$ $\mathbf{R}^{-1} = \frac{1}{T} \begin{pmatrix} 1 & -\alpha \\ 0 & 1 \end{pmatrix}$

$$\ln \mathcal{L}(S, B) = n_{on} \ln[(S + \alpha B)T] + n_{off} \ln BT - (S + (1 + \alpha)B)T$$

- MLE & (co)variances of S and B are:

$$\hat{B} = \frac{1}{T} n_{off}$$

$$\sigma_B^2 = \frac{1}{T^2} n_{off}$$

$$\hat{S} = \frac{1}{T} (n_{on} - \alpha n_{off})$$

$$\sigma_S^2 = \frac{1}{T^2} (n_{on} + \alpha^2 n_{off})$$

This is what you would expect!

$$\text{cov}(\hat{S}, \hat{B}) = -\frac{1}{T^2} \alpha n_{off}$$

TS for On/Off problems

- Test statistic for source detection in On/Off problems is:

$$TS = 2 \left[n_{on} \ln \frac{(1 + \alpha)n_{on}}{\alpha(n_{on} + n_{off})} + n_{off} \ln \frac{(1 + \alpha)n_{off}}{(n_{on} + n_{off})} \right]$$

- Significance is: $\sigma = \sqrt{TS}$
- This is the famous “Li & Ma” formula from: ApJ 272, 317 (1983) - 493 citations on ADS
- Probably, you wouldn’t arrive at this formula using ad hoc estimation methods
- P-values: `scipy.stats.chi2.sf(TS, 1)`

Eg: 1ES1218+304 w/VERITAS

Discovery of Variability in the Very High Energy γ -Ray Emission of 1ES 1218+304 with VERITAS

Acciari, et al., ApJ, 709, 163 (2010)

Table 1 summarizes the results of the VERITAS observations of 1ES 1218+304. For the spectral analysis, we report an excess of 1155 events with a statistical significance of 21.8 standard deviations, σ , from the direction of 1ES 1218+304 during the 2008-2009 campaign (2808 signal events, 4959 background events with a normalization of 0.33). Figure 2 shows the corresponding time-averaged differential energy spectrum. The spectrum extends from 200 GeV to 1.8 TeV and is well described ($\chi^2/\text{dof} = 8.2/7$) by a power law,

$$n_{off} = 4959$$

$$n_{on} = 2808$$

$$\alpha = 1/3$$

$$T = 27.2 \text{ hr}$$

$$\hat{S} = 42.5 \text{ hr}^{-1}$$

$$\sigma_S = 2.1 \text{ hr}^{-1}$$

$$TS = 474.9$$

$$\sigma = 21.8$$

$$P - \text{value} = 2.8 \times 10^{-105}$$

Table 1. Summary of observations and analysis of 1ES 1218+304^a.

	Live Time [hours]	Zenith [$^\circ$]	Significance [σ]	$\Phi(> 200 \text{ GeV})$ [$10^{-12} \text{ cm}^{-2} \text{ s}^{-1}$]	Units of Crab Nebula flux ($E > 200 \text{ GeV}$)
2006-2007 ^b	17.4	2-35	10.4	$12.2 \pm 2.6_{stat}$	0.05 ± 0.011
2008-2009	27.2	2-30	21.8	$18.4 \pm 0.9_{stat}$	0.07 ± 0.004

$$\sigma_{POE} = \frac{\hat{S}}{\sigma_S} = 19.9 \approx \frac{18.4}{0.9}$$

Ratio of value to error - used as "significance" before Li&Ma

Eg: 1ES1218+304 w/VERITAS

```
# lima.py - 2013-05-15 SJF
# Example of Li & Ma significance calculation
import math, scipy.stats

def ts_lima(non,noff,alpha):
    opa = 1.0+alpha
    ntot = non+noff
    return 2.0*(non*math.log(opa*non/alpha/ntot) \
               + noff*math.log(opa*noff/ntot))

non      = 2808
noff     = 4959
alpha    = 1.0/3
T        = 27.2

S_hat    = (non - noff*alpha)/T
sig2_S   = (non + noff*alpha**2)/T**2
ts       = ts_lima(non,noff,alpha)
signif   = math.sqrt(ts)
Pval     = scipy.stats.chi2.sf(ts,1)

print S, math.sqrt(sig2_S), ts, signif, Pval
```

Ratio of value to error - used as significance before Li&Ma

Confidence regions

In problems with multiple parameters.

- Saw earlier that we can calculate “asymmetric errors” by finding points where $2\ln\mathcal{L}$ decreases by 1.0: 2-sided 1σ confidence interval (68%)
- Actually this comes from LRT (Wilks’ theorem). This is region where null hypothesis that parameter value has some value cannot be rejected at given confidence level.
- But what to do if likelihood depends on more than our parameter of interest?
- It depends...

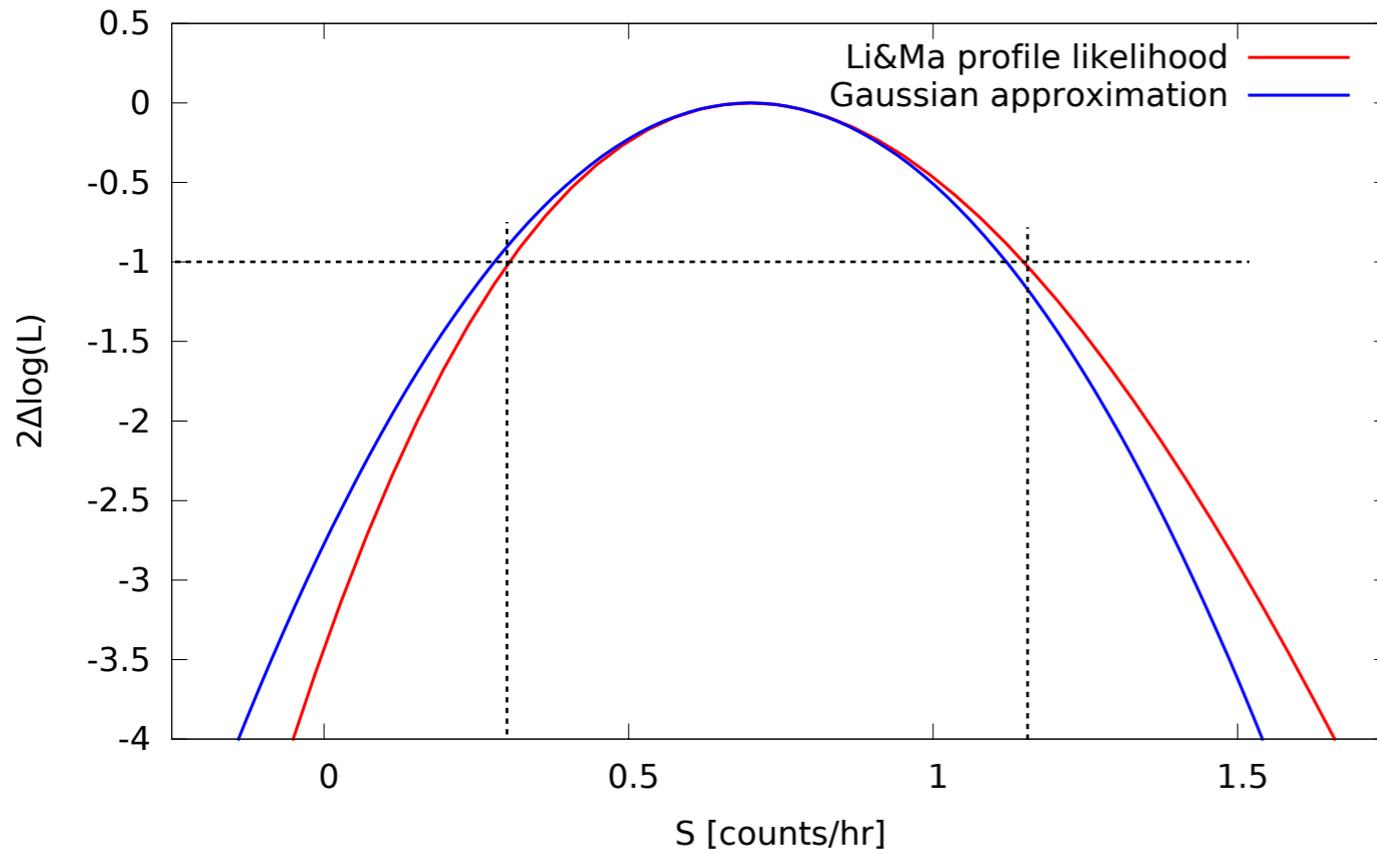
Profile likelihood

Confidence regions with nuisance parameters

Rolke, et al., NIM A, 551, 493 (2005)

- Often we are either concerned only with the one parameter, or wish to treat the multiple parameters separately (ignore covariance).
- Produce “profile log-likelihood” curve, a function of only one parameter (at a time), maximized over all others.
- LRT says this should behave as $\chi^2(1)$.
- Define confidence region using this function exactly as before.

Example of profile likelihood



$$\hat{S} = 0.7_{-0.39}^{+0.45} \text{ hr}^{-1}$$

- Our 1ES1218 example isn't very enlightening here, so take:

$$n_{off} = 24$$

$$n_{on} = 15$$

$$\alpha = 1/3$$

$$T = 10.0 \text{ hr}$$

- Giving:

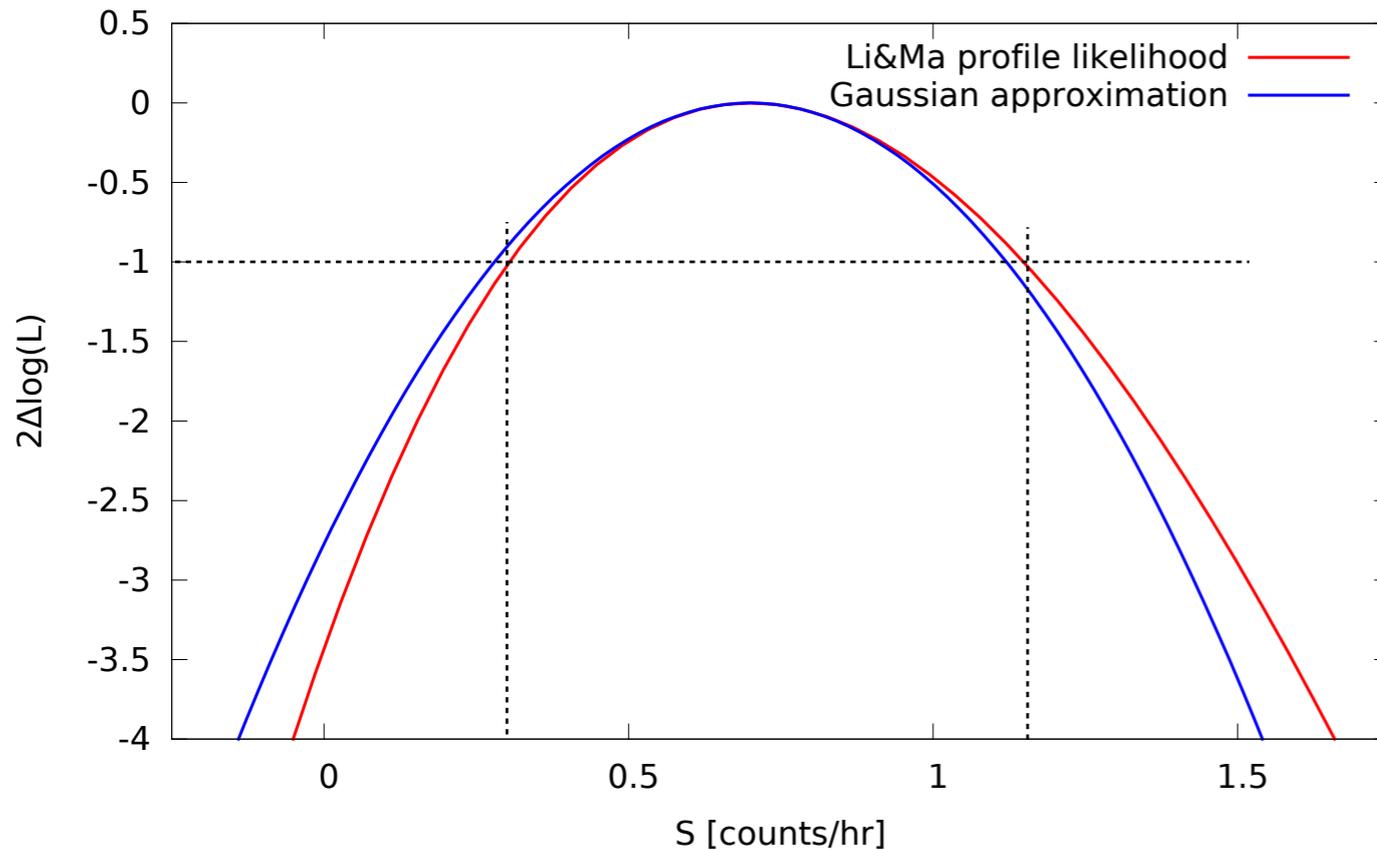
$$\hat{S} = 0.7 \text{ hr}^{-1}$$

$$\sigma_S = 0.42 \text{ hr}^{-1}$$

$$TS = 3.43$$

$$\sigma = 1.85$$

Example of profile likelihood



$$\hat{S} = 0.7_{-0.39}^{+0.45} \text{ hr}^{-1}$$

This is not a significant result, so we would usually not claim a detection. Provide an upper limit instead.

- Our 1ES1218 example isn't very enlightening here, so take:

$$n_{off} = 24$$

$$n_{on} = 15$$

$$\alpha = 1/3$$

$$T = 10.0 \text{ hr}$$

- Giving:

$$\hat{S} = 0.7 \text{ hr}^{-1}$$

$$\sigma_S = 0.42 \text{ hr}^{-1}$$

$$TS = 3.43$$

$$\sigma = 1.85$$

Example of profile likelihood

```
# conf_lima_1d.py - 2013-05-25 SJF
# 1-D 2-sided confidence interval in Li & Ma problem
from math import *
import scipy.stats, scipy.optimize, sys
# non, noff, alpha, T = (2808, 4959, 1.0/3, 27.2)
non, noff, alpha, T = (15, 24, 1.0/3, 10.0)
 $2\Delta\log(L)$  C = 0.68; # Use 1-sigma confidence region
d2logL = scipy.stats.chi2.ppf(C, 1)
def logL(S, B):
    return non*log(max((S+alpha*B)*T, sys.float_info.min)) + \
    noff*log(max(B*T, sys.float_info.min)) - (S+(1+alpha)*B)*T
def profileLogL(S):
    opt_fn = lambda B: -logL(S, B)
    opt_res = scipy.optimize.minimize(opt_fn, 1)
    return -opt_res.fun
S_hat = (non-noff*alpha)/T
B_hat = noff/T
logL_max = logL(S_hat, B_hat)
sig_S = sqrt(non+noff*alpha**2)/T
TS = -2.0*(profileLogL(0)-logL_max)
root_fn = lambda S: 2.0*(profileLogL(S)-logL_max)+d2logL
S_lo = scipy.optimize.brentq(root_fn, 1e-8, S_hat)
S_hi = scipy.optimize.brentq(root_fn, S_hat, 1e8)
print S_hat, S_lo-S_hat, S_hi-S_hat, sig_S, TS, sqrt(TS)
```

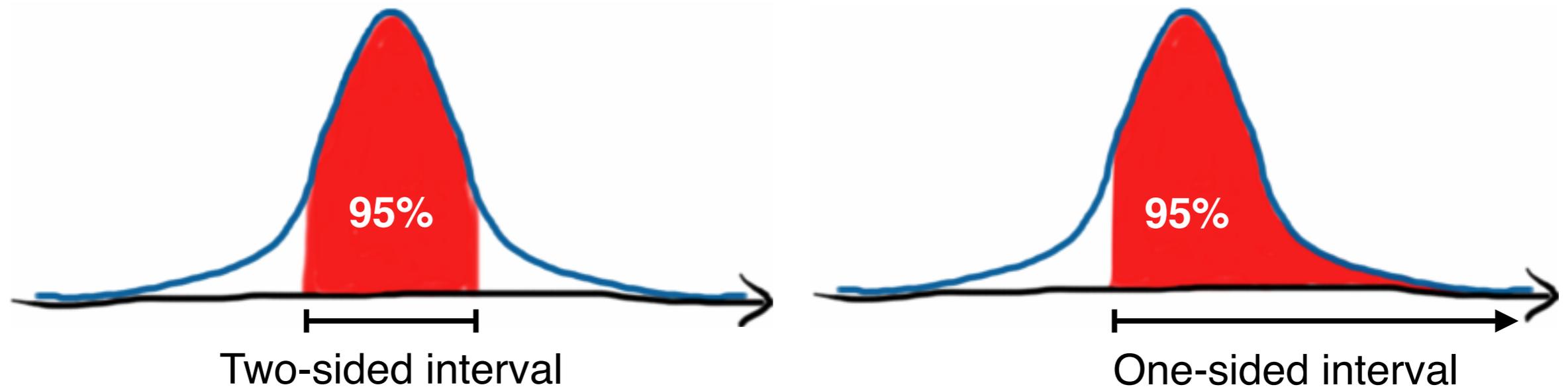
Example of profile likelihood

```
# conf_lima_1d.py - 2013-05-25 SJF
# 1-D 2-sided confidence interval in Li & Ma problem
from math import *
import scipy.stats, scipy.optimize, sys
# non, noff, alpha, T = (2808, 4959, 1.0/3, 27.2)
non, noff, alpha, T = (15, 24, 1.0/3, 10.0)
 $2\Delta\log(L)$  C = 0.68; # Use 1-sigma confidence region
d2logL = scipy.stats.chi2.ppf(C, 1)
def logL(S, B):
    return non*log(max((S+alpha*B)*T, sys.float_info.min)) + \
    noff*log(max(B*T, sys.float_info.min)) - (S+(1+alpha)*B)*T
def profileLogL(S):
    opt_fn = lambda B: -logL(S, B)
    opt_res = scipy.optimize.minimize(opt_fn, 1)
    return -opt_res.fun
S_hat = (non-noff*alpha)/T
B_hat = noff/T
logL_max = logL(S_hat, B_hat)
sig_S = sqrt(non+noff*alpha**2)/T
TS = -2.0*(profileLogL(0)-logL_max)
root_fn = lambda S: 2.0*(profileLogL(S)-logL_max)+d2logL
S_lo = scipy.optimize.brentq(root_fn, 1e-8, S_hat)
S_hi = scipy.optimize.brentq(root_fn, S_hat, 1e8)
print S_hat, S_lo-S_hat, S_hi-S_hat, sig_S, TS, sqrt(TS)
```

Frequentist upper limits

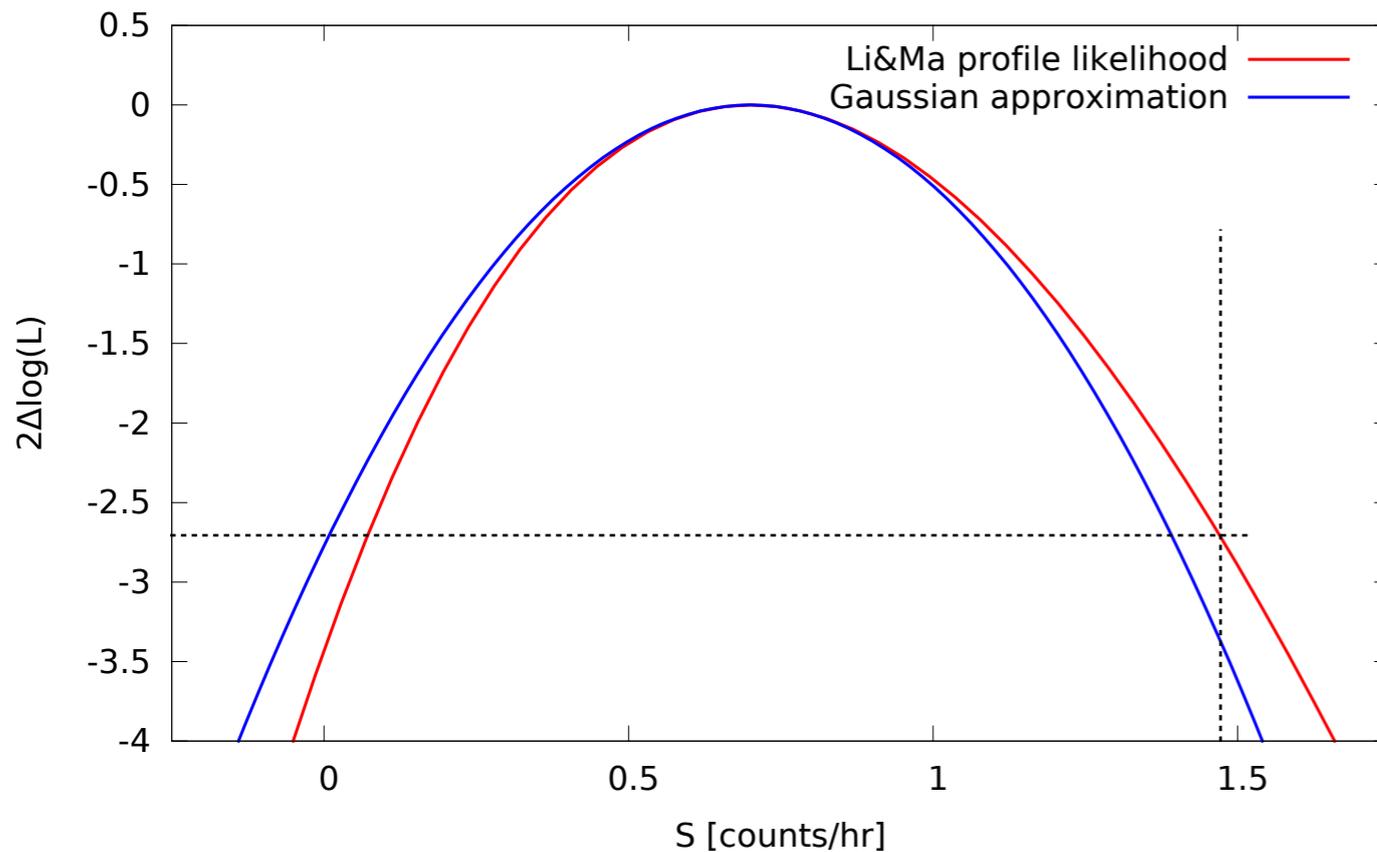
One-sided confidence region using profile likelihood

Rolke, et al., NIM A, 551, 493 (2005)



- In two-sided interval search for two points $S_{1,2}$ where $-2\Delta \ln \mathcal{L}(S_{1,2}) = x$ with $\chi^2(x, 1) = C$
- For one-sided interval (with $C > 0.5$) we need to find single such point with $S_{UL} > \hat{S}$ and for which $0.5 + \chi^2(x, 1)/2 = C$ (or $\chi^2(x, 1) = 2C - 1$)
- E.g. for $C=0.95$ we search $-2\Delta \ln \mathcal{L}(S_{UL}) = 2.71$

Example of profile likelihood



$$\hat{S} = 0.7^{+0.45}_{-0.39} \text{ hr}^{-1}$$

- Frequentist upper limit at 95% confidence level:

$$S_{<95\%} = 1.47 \text{ hr}^{-1}$$

Exercise: adapt 2-sided interval code to calculate this

- Our 1ES1218 example isn't very enlightening here, so take:

$$n_{off} = 24$$

$$n_{on} = 15$$

$$\alpha = 1/3$$

$$T = 10.0 \text{ hr}$$

- Giving:

$$\hat{S} = 0.7 \text{ hr}^{-1}$$

$$\sigma_S = 0.42 \text{ hr}^{-1}$$

$$TS = 3.43$$

$$\sigma = 1.85$$

Bayesian statistics

- Likelihood function has no meaning itself, e.g., it is not a probability. Its usefulness comes from theorems such as the LRT.
- MLE belongs to the class of “frequentist” statistical methods: talk about the results of repeated hypothetical experiments.
- Saw how to produce confidence intervals: true parameter value would lie inside the interval in a certain % of hypothetical expts.
- Somewhat awkward language ???

Bayesian statistics

- In Bayesian statistics we talk about the “probability” that the parameters have certain values.

- Bayes’ theorem:

Posterior probability density →
$$P(\Theta|X) = \frac{P(\Theta)P(X|\Theta)}{P(X)} \propto P(\Theta)\mathcal{L}(\Theta|X)$$

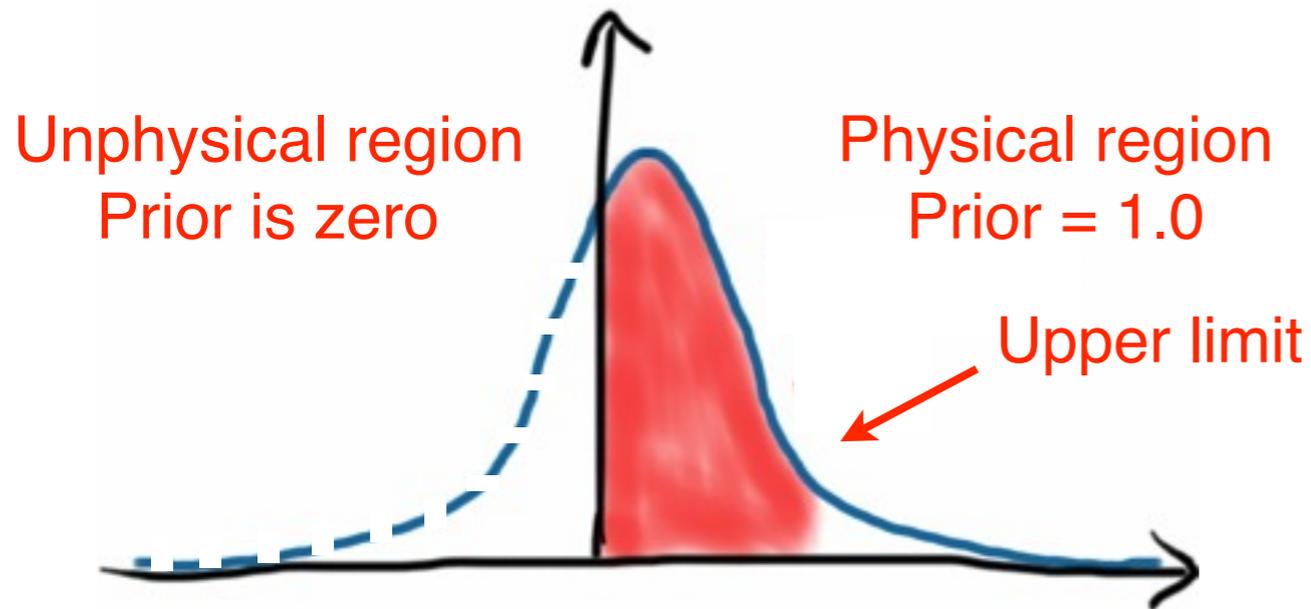
Prior probability density ↓
Likelihood ↓

relates probability after experiment has been done to probability before.

- Can think of this as refining our belief about the model through experimental results.

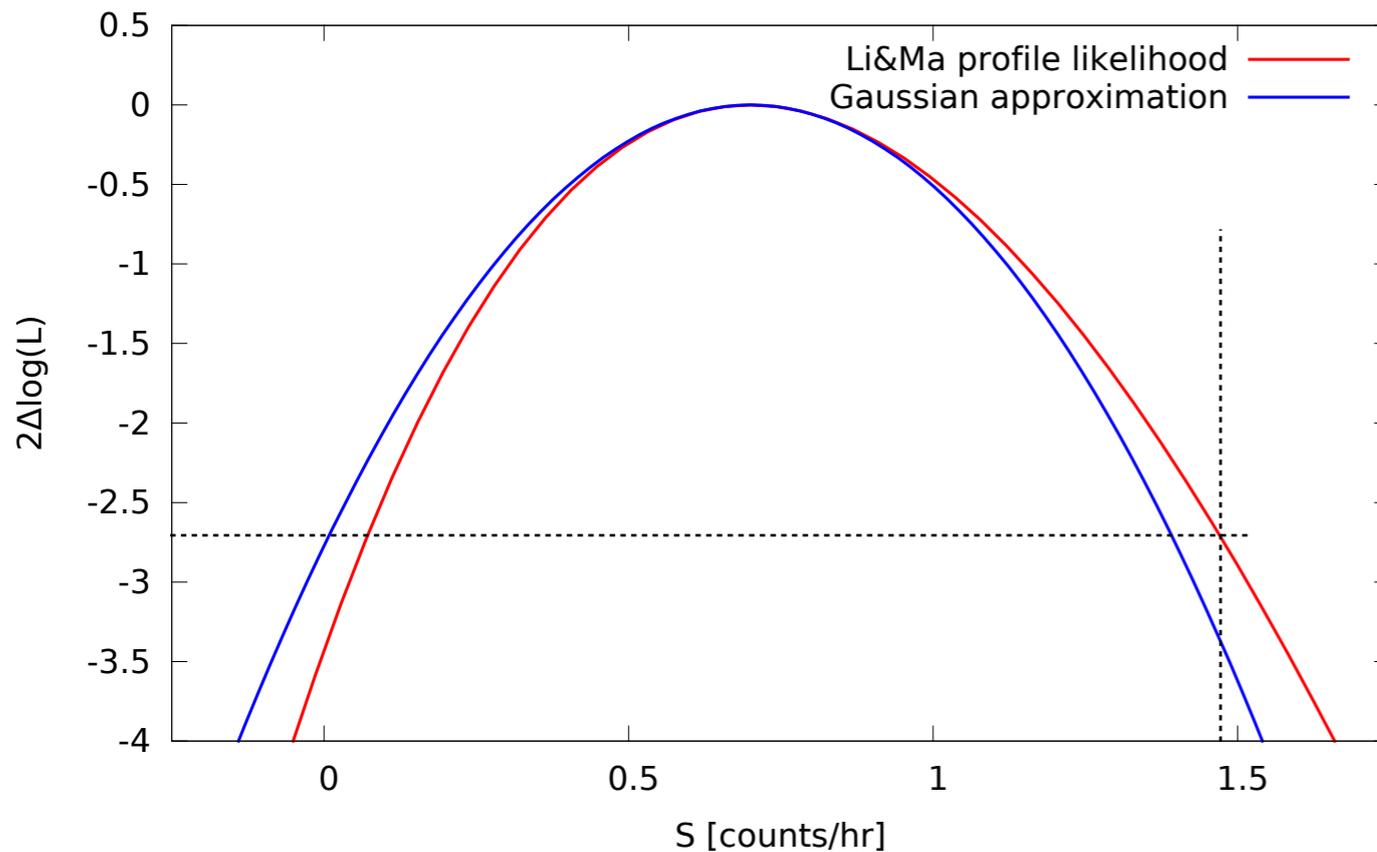
Bayesian upper limits

Or more correctly “Quasi-Bayesian” or “Bayesian-like”



- Bayesian confidence regions correspond to what you would expect...
- ... they are regions that contain a certain fraction of the posterior probability.
- Integrate over parameter from lower bound to find point where integral reaches C% of total.
- In case of multiple parameters, use the profile likelihood. Not strictly a Bayesian approach.

Example of profile likelihood



$$\hat{S} = 0.7^{+0.45}_{-0.39} \text{ hr}^{-1}$$

- Frequentist upper limit at 95% confidence level:

$$S_{<95\%} = 1.47 \text{ hr}^{-1}$$

- Bayesian 95% upper limit:

$$S_{<95\%} = 1.54 \text{ hr}^{-1}$$

- Our 1ES1218 example isn't very enlightening here, so take:

$$n_{off} = 24$$

$$n_{on} = 15$$

$$\alpha = 1/3$$

$$T = 10.0 \text{ hr}$$

- Giving:

$$\hat{S} = 0.7 \text{ hr}^{-1}$$

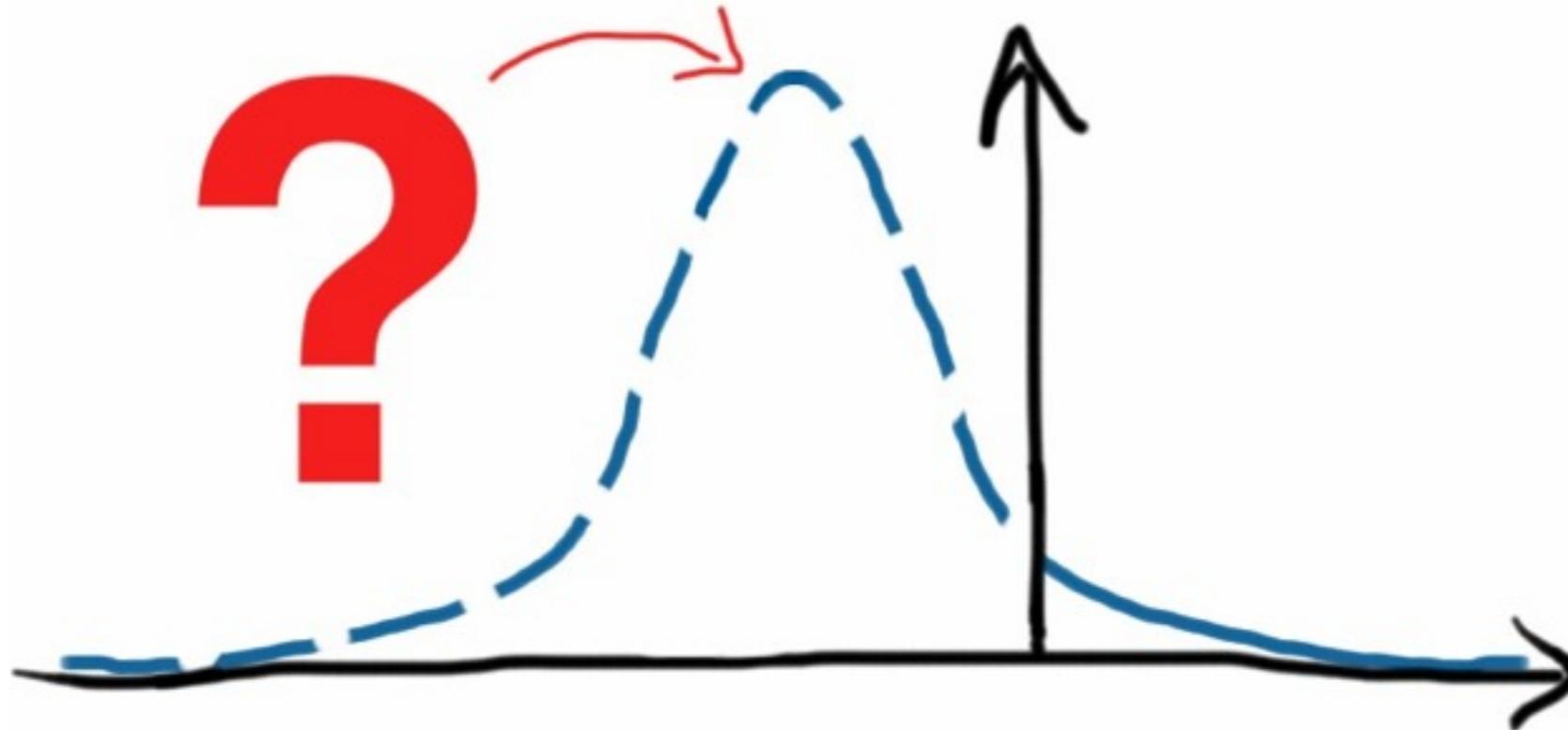
$$\sigma_S = 0.42 \text{ hr}^{-1}$$

$$TS = 3.43$$

$$\sigma = 1.85$$

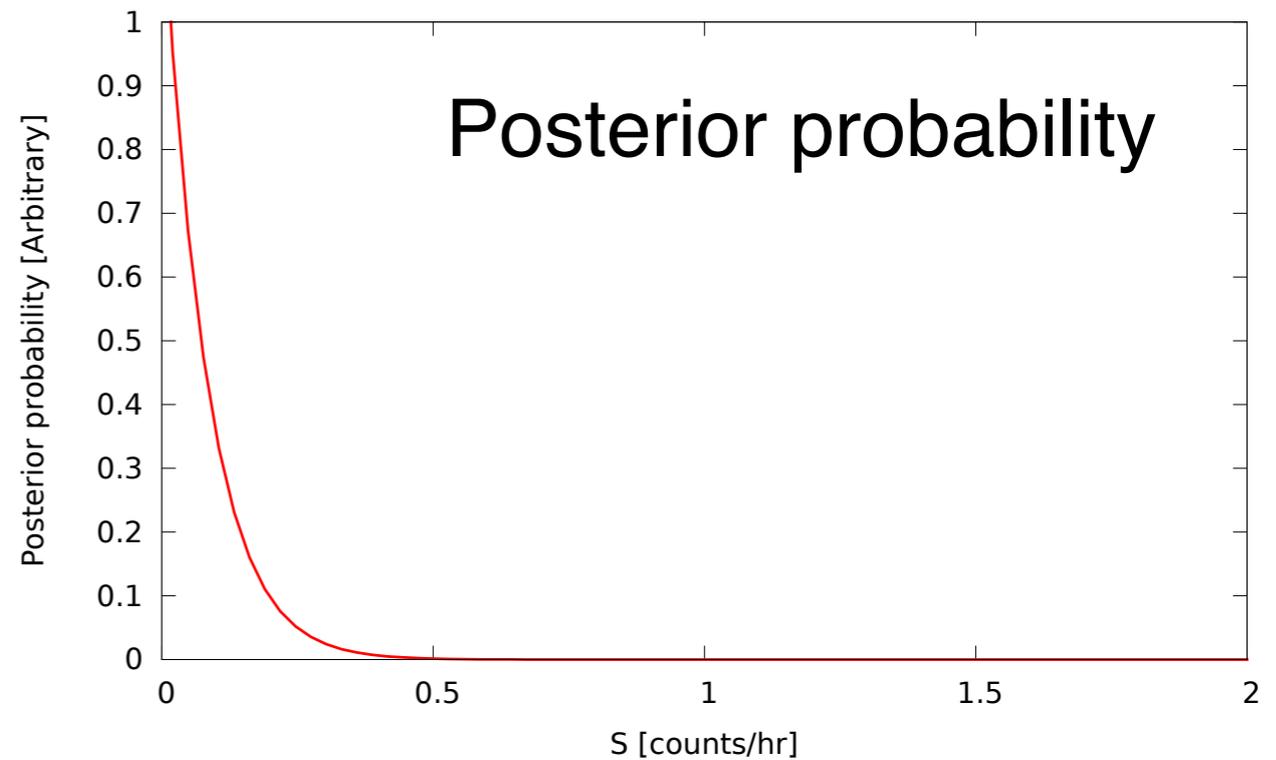
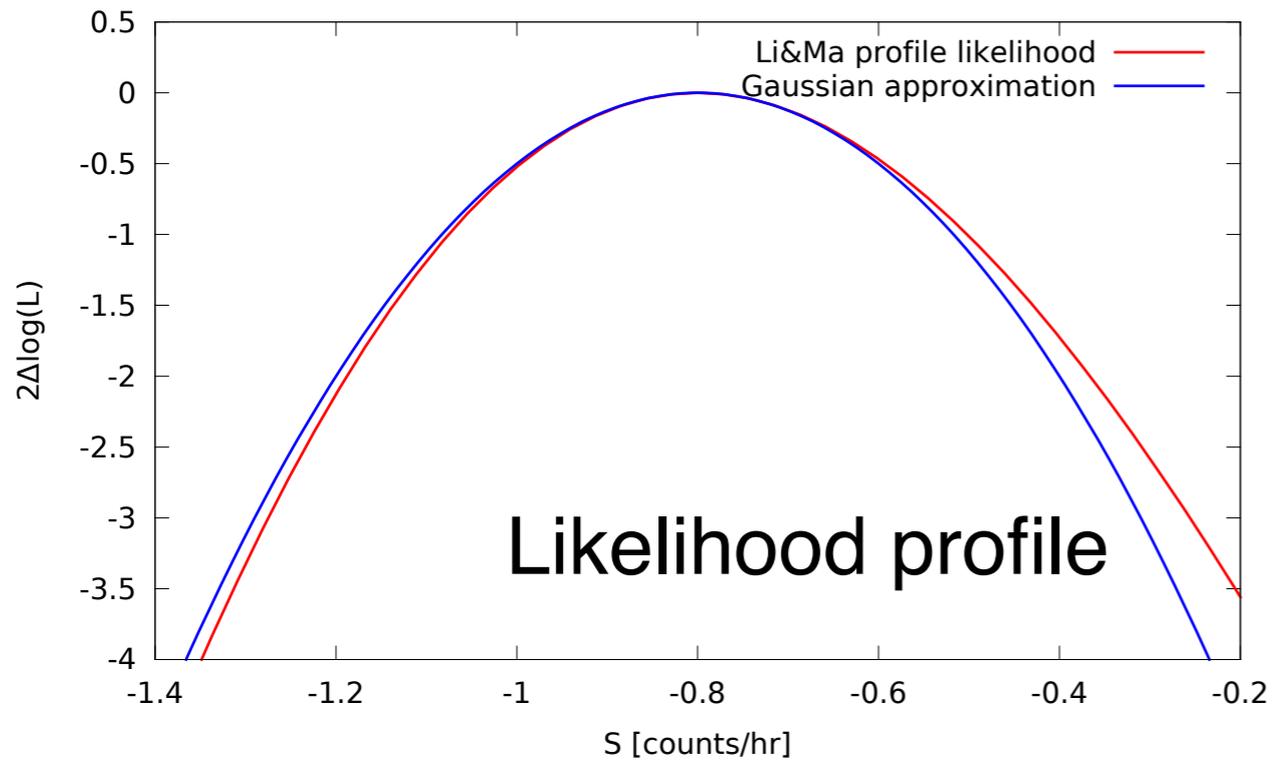
Why have two methods?

The problem of unphysical upper limits



- Unphysical frequentist upper limits occur can occur if the peak of the likelihood is in an unphysical region of the parameter space.
- More complex (or ad hoc) approaches fix this.
- But Bayesian upper limits are not affected.

Example of unphysical MLE



$n_{off} = 36$
 $n_{on} = 4$
 $\alpha = 1/3$
 $T = 10.0 \text{ hr}$

MLE is negative - not physical for source flux.

$\hat{S} = -0.8 \text{ hr}^{-1}$
 $\sigma_S = 0.28 \text{ hr}^{-1}$
 $TS = 5.80$
 $\sigma = -2.41$

“Background fluctuation”,
 fewer “On” counts than
 expected given “Off”

Frequentist UL: $S_{<95\%} = -0.29 \text{ hr}^{-1}$ - unphysical
 Bayesian UL: $S_{<95\%} = 0.43 \text{ hr}^{-1}$ - OK!

```

# ul_lima_bayes_1d.py - 2013-05-25 SJF
# Bayesian upper limit in Li & Ma problem
from math import *
import scipy.stats, scipy.optimize, scipy.integrate, sys
# non, noff, alpha, T = (2808, 4959, 1.0/3, 27.2)
# non, noff, alpha, T = (15, 24, 1.0/3, 10.0)
non, noff, alpha, T = (4, 36, 1.0/3, 10.0)
C = 0.95; # Use 95% confidence region
def logL(S,B):
    return non*log(max((S+alpha*B)*T,sys.float_info.min)) + \
    noff*log(max(B*T,sys.float_info.min)) - (S+(1+alpha)*B)*T
def profileLogL(S):
    opt_fn = lambda B: -logL(S,B)
    opt_res = scipy.optimize.minimize(opt_fn, 1)
    return -opt_res.fun
S_hat      = (non-noff*alpha)/T
sig_S      = sqrt(non+noff*alpha**2)/T
logL_max   = profileLogL(S_hat)
def logPrior(S):
    return log(1);
def logPosterior(S):
    return logPrior(S)+profileLogL(S)-logL_max
def integralPosterior(Smax):
    integrand = lambda S: exp(logPosterior(S))
    y, err = scipy.integrate.quad(integrand,0,Smax)
    return y
total_integral = integralPosterior(S_hat+100*sig_S);
root_fn = lambda S: integralPosterior(S) - total_integral*C
S_ul = scipy.optimize.brentq(root_fn, 0, S_hat+100*sig_S)
print S_ul, integralPosterior(S_ul)/total_integral, total_integral

```

Good practices

- It is always best to define all the parameters of an analysis before looking at the data.
 - Data selection “cuts”
 - Thresholds for claiming detection.
- It is tempting to adjust the analysis procedure to enhance some small signal, **BUT THIS IS FRAUGHT WITH DANGER!**
- Best practice is to do a blind analysis. Use MC or side-band data to refine analysis in advance.
- But this is not always possible...

Trials factors

Or the “look-elsewhere effect”

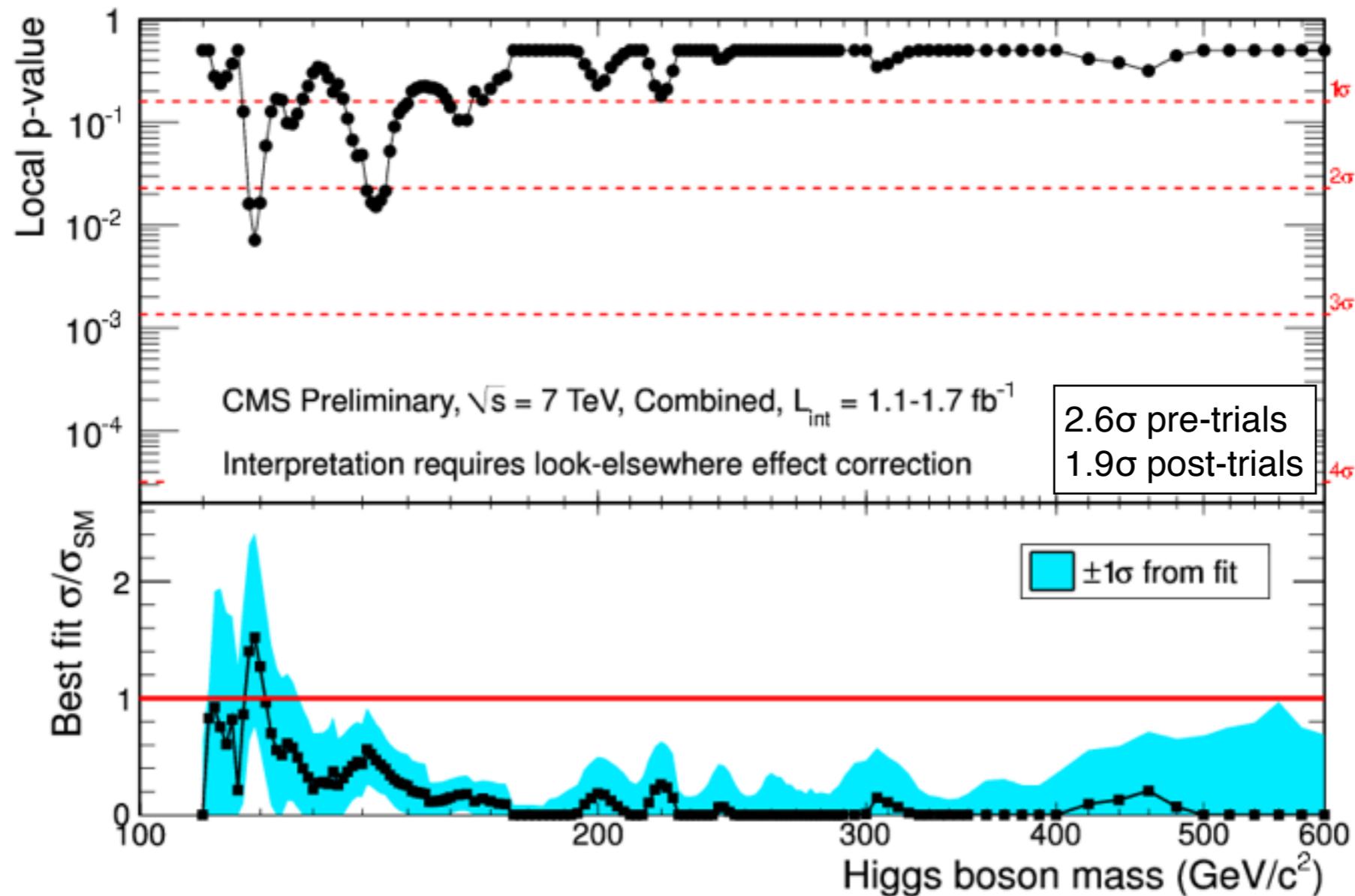
- Often you simply don't know enough in advance to fully determine the analysis, e.g.
 - the mass of the DM particle (or Higgs)
 - the locations of sources in the sky etc...
- So, you must look through the data and search for a significant excess signal ...
- ... and unfortunately you must pay a statistical penalty for doing so.

Trials factors

Or the “look-elsewhere effect”

- If after making N independent tests of for a significant event (e.g N energy channels)
- the most significant test had a P-value of: P_{pre}
- then to account for the number of “trials” you must scale the P-value as: $P_{post} = 1 - (1 - P_{pre})^N$
- For example, a 4σ event has a P-value of $P_{pre} = 6.3 \times 10^{-5}$. With 1000 trials, the post-trial P-value of $P_{post} = 1 - (1 - 6.3 \times 10^{-5})^{1000} = 0.06$ which is equivalent to a 1.9σ event.

For example... (Higgs@CMS)



- But.. how many truly independent samples?
- Depends on resolution & can be difficult to estimate

Detectability / Sensitivity

- Interested in detectability of sources, i.e. sensitivity of instrument for given threshold.
- Consider “no fluctuations” case where:

$$n_{on}^{NF} = (S_t + \alpha B_t)T, \quad n_{off}^{NF} = B_t T$$

- Then test statistic is:

$$TS^{NF} = 2 \left[(S_t + \alpha B_t)T \ln \frac{(1 + \alpha)(S_t + \alpha B_t)T}{\alpha(S_t + (1 + \alpha)B_t)T} + B_t T \ln \frac{(1 + \alpha)B_t T}{(S_t + (1 + \alpha)B_t)T} \right]$$

Detectability / Sensitivity

- Weak source case: $S_t \ll \alpha B_t$

$$\sigma^{NF} = \sqrt{TS^{NF}} \approx \frac{\sqrt{T}}{\sqrt{1+\alpha}} \frac{S_t}{\sqrt{\alpha B_t}}$$

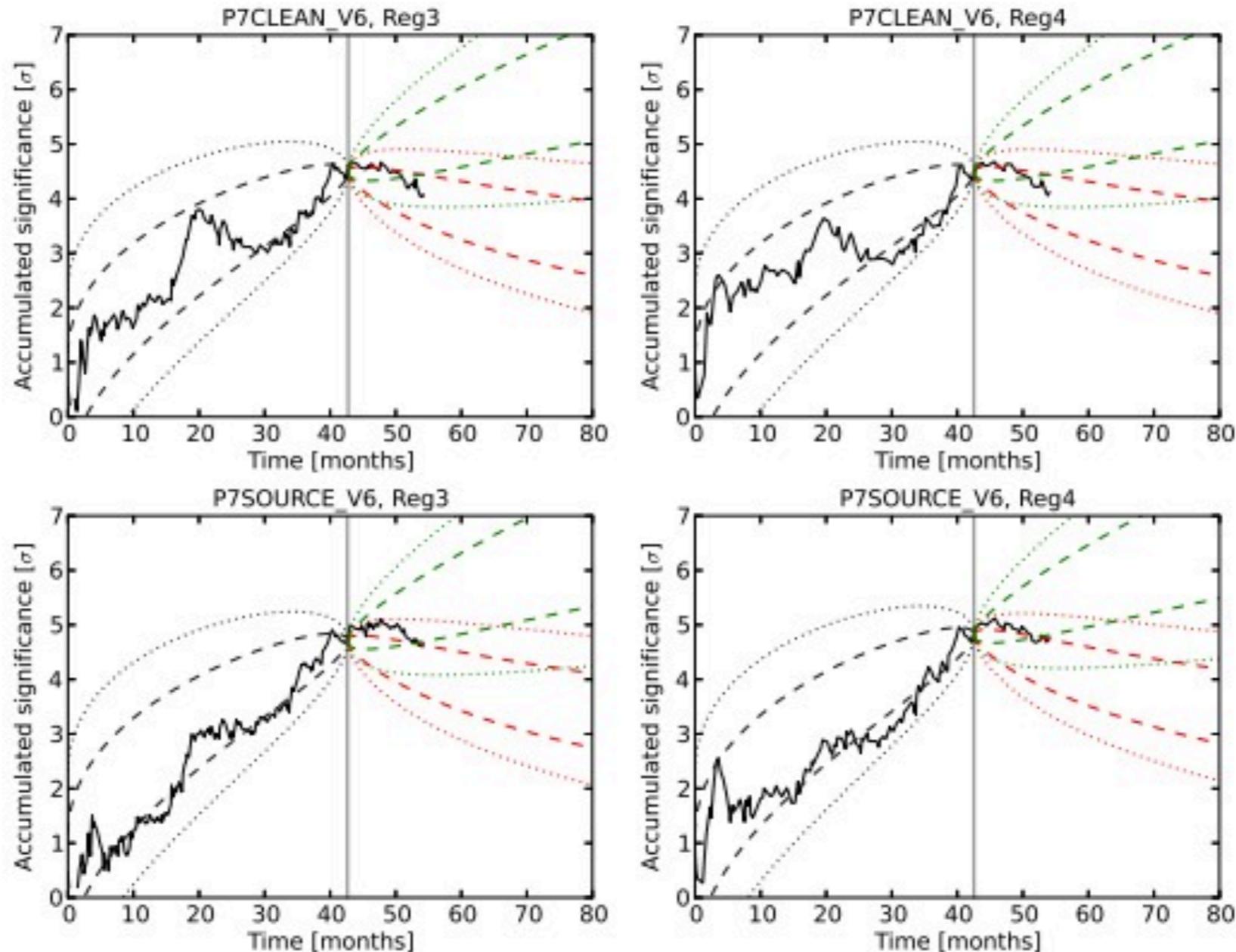
Grows as sqrt(T)

- Strong source case: $S_t \gg \alpha B_t$

$$\sigma^{NF} = \sqrt{TS^{NF}} \approx \sqrt{2S_t T \ln(1 + 1/\alpha)}$$

Note what happens here when $\alpha \rightarrow 0$ (which corresponds to perfectly well determined “zero” on-source background) the significance becomes infinite. If you have no background then even one event is a significant.

Meanwhile, over in real life...



DM line with LAT
Weniger, Proc 2nd
Fermi Symposium
(2012)

Figure 2: Time evolution of the accumulated significance of the line feature in Gaussian sigma in comparison with the expectations. The dashed (dotted) lines show the 68%CL (95%CL) bands corresponding to a real signal (green), a statistical fluke (red) and a steady source in the past (black). The solid line shows the actual behaviour of the feature in the LAT data. The vertical line indicates the 8th of March 2012, which we use as reference point for a new trial-free measurement. In the fit, the line energy is fixed to $E_\gamma = 129.8$ GeV (see text for details). The four panels show results for the ROIs Reg3 (left) and Reg4 (right) from Weniger [2012], for P7CLEAN_V6 (top) and P7SOURCE_V6 (bottom) class events. Data until the 22nd of February 2013 is taken into account.

Detectability / Sensitivity

Minimum source strength to achieve detection at some threshold σ_{det}

- Weak source case: $S_t \ll \alpha B_t$

$$S_t > \frac{\sigma_{det} \sqrt{\alpha B_t}}{\sqrt{T}} \sqrt{1 + \alpha}$$

Minimum detectable flux decreases as $1/\sqrt{T}$ and depends on B_t : “Background-dominated regime”

- Strong source case: $S_t \gg \alpha B_t$

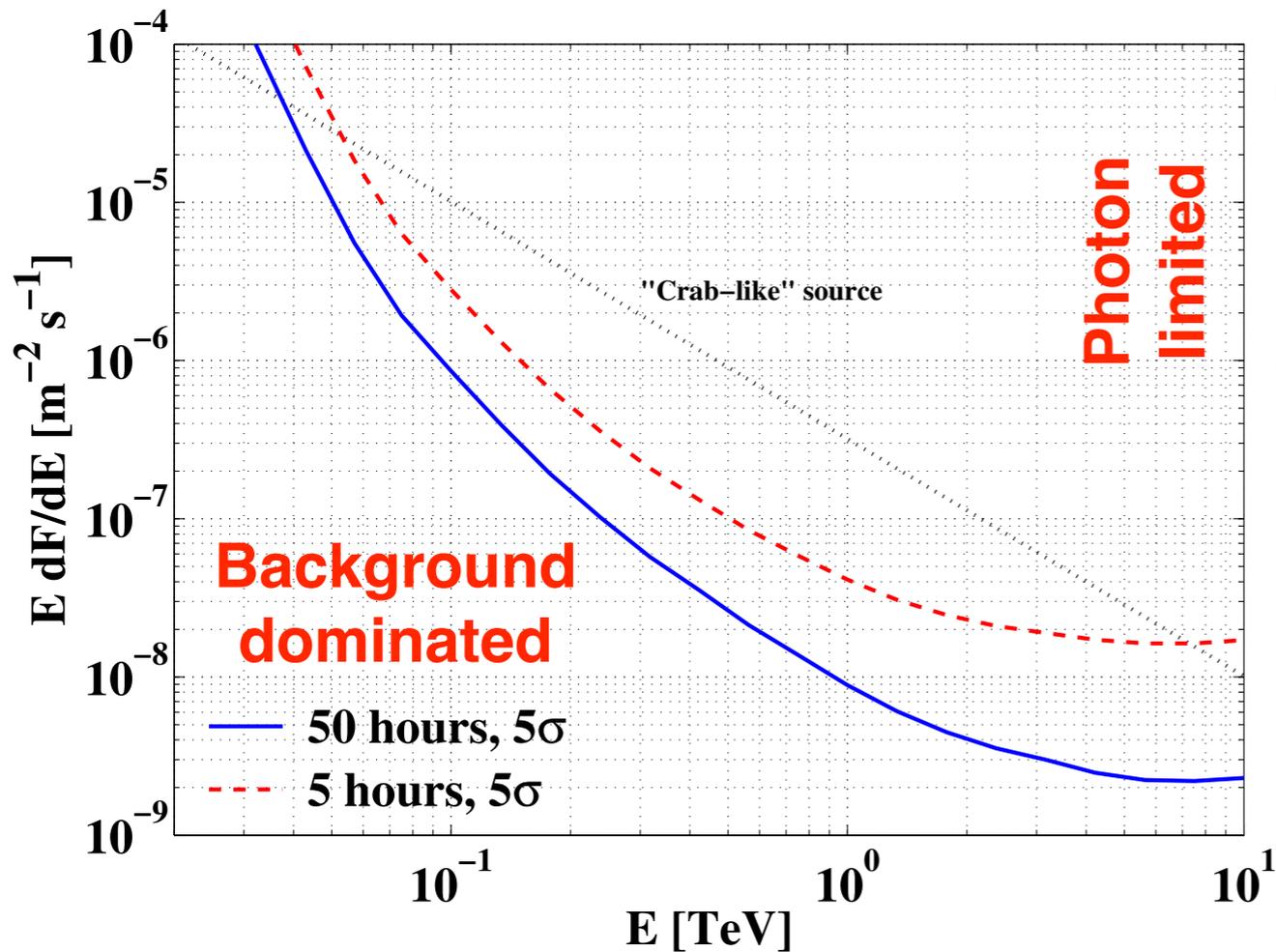
$$S_t > \frac{\sigma_{det}^2}{T} \frac{1}{2\sqrt{1 + 1/\alpha}}$$

Roughly this says that the number of detected photons must be larger than σ^2 (times some constant): $S_t T = n_{det} > C \sigma_{det}^2$
eg. must detect 25 photons for 5σ .

Minimum detectable flux decreases as $1/T$ and is independent of B_t : “Photon-limited regime”

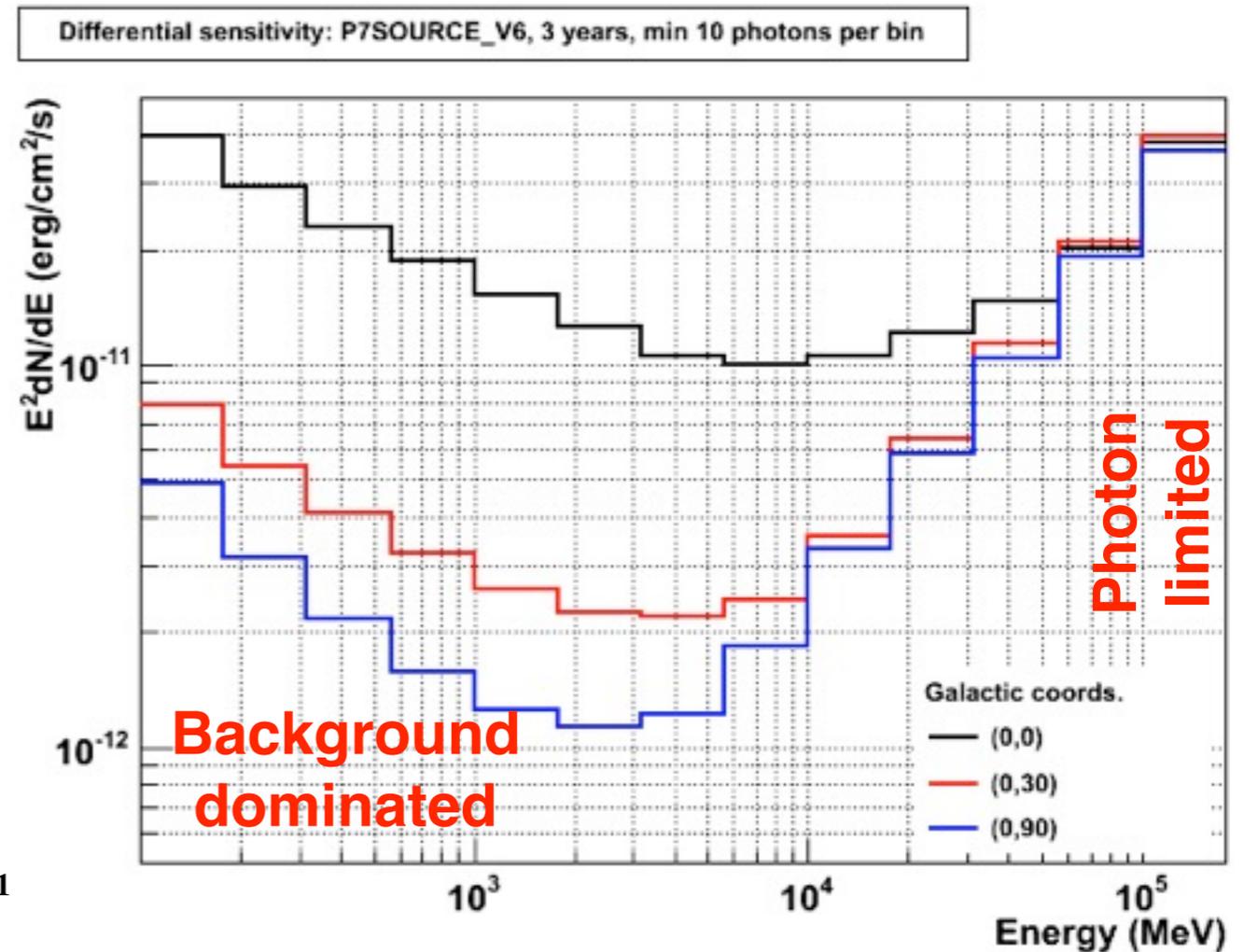
Detectability / Sensitivity

“Differential sensitivity” plots, i.e. sensitivity in logarithmic energy bands



Sensitivity for ACT array of 4 telescopes for 5 and 50 hours of observation.

Low energies: $\sqrt{10}$ × improvement.
High energies: 10 × improvement.



LAT sensitivity from FSSC site for different background levels (Galactic or extra-galactic).

Low energies: big dependency
High energies: almost no dependency.

Systematic errors

What if assumed value of alpha is incorrect?

- Assume there is no real source:

$$n_{on}^{NF} = \alpha_t B_t T = \alpha(1 + \delta) B_t T, \quad n_{off}^{NF} = B_t T$$

where the error in alpha is small: $\delta \ll 1$

- Then: $\hat{S}^{NF} = B_t \alpha \delta$

$$\sigma^{NF} = \sqrt{T S^{NF}} \approx \frac{\sqrt{T}}{\sqrt{1 + \alpha}} \delta \sqrt{\alpha B_t}$$

- This looks like a real signal. Accurate knowledge of experimental response is critical. **MLE is only as good as the model!**

Review

- ML provides “cookbook” for estimation and hypothesis testing:
 - estimates: maximum of likelihood
 - errors: curvature of log-likelihood surface
 - TS and significance: is improvement in $\log-\mathcal{L}$ over null hypothesis consistent with χ^2 ?
- Significance expected to grow as \sqrt{T} , but sensitivity can improve as $1/T$ if photon limited.
- Systematic errors important to consider

Onwards to LAT analysis...

- LAT ML analysis is fundamentally the same as what we have seen here (but more complex).
- Channels organized by sky position and energy (i.e. 3-dimensions). Million channels typical.
- Model is Poisson for each channel with mean determined by:
 - spatial-spectral model provided by user
 - observational response (calculated by software from IRFs provided by LAT team)
- MLE by software: errors, covariances, TS, etc