

Exploration of electromagnetic outflows of inspiraling binary neutron stars

Dimitrios Skiathas

SURA, NASA GSFC, University of Patras

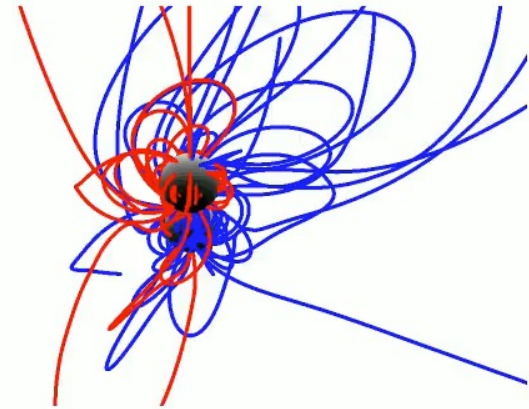
Constantinos Kalapotharakos, *NASA GSFC*

Zorawar Wadiasingh, *NASA GSFC, UMD*

Demos Kazanas, *NASA GSFC*

Alice Harding, *LANL*

Paul Kolbeck, *UW*

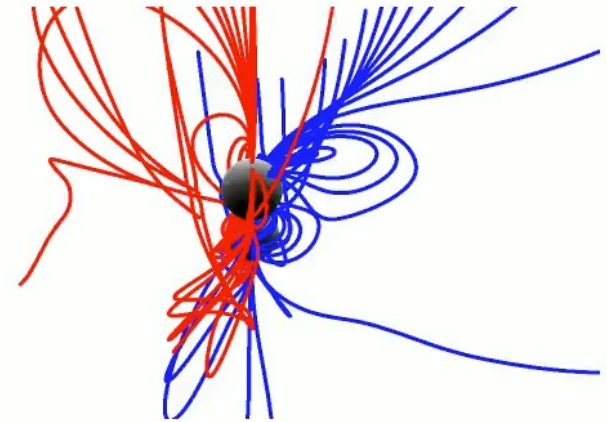


Introduction

- Interacting neutron star magnetospheres can release large amount of energy through electromagnetic emissions.

Piro 2012, Tsang et al. 2012, Pelenzuela et al. 2013, Lyutikov 2019, Sridhar et al. 2021, Most & Philippov 2020,2022, Cooper et al 2023

- Precursor signal to neutron star mergers.
- Echoes contributing to the GRB afterglows.
- Our goal is to explore the parameter space regarding the emission patterns.

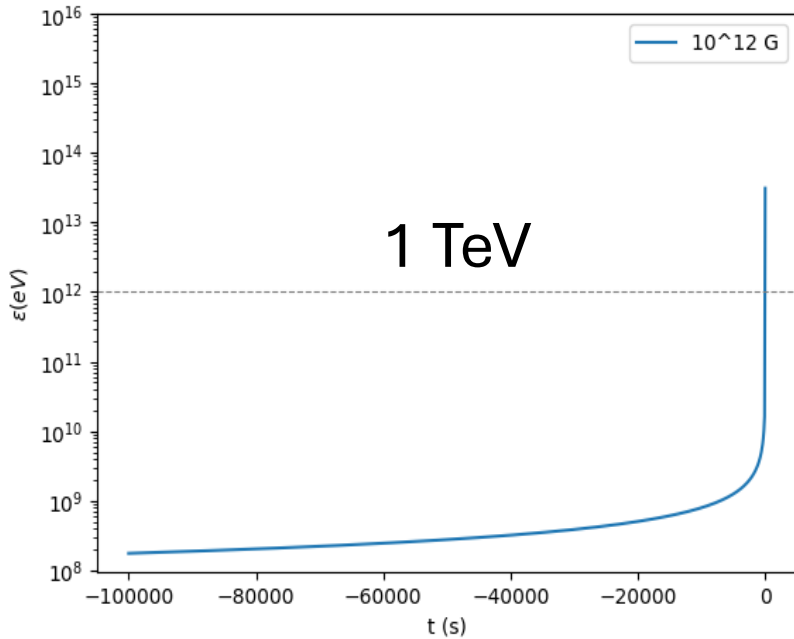


For a review see Suvorov et al. 2024

Assuming emissions due to CR at the radiation reaction limit regime

$$\frac{d\gamma_L}{dt} = \frac{q_e v E_{acc}}{m_e c^2} - \frac{2q_e \gamma_L^4}{3R_c^2 m_e c} = 0, \text{ with } E_{acc} \sim B_{LC} \text{ and } R_c \sim R_{LC}$$

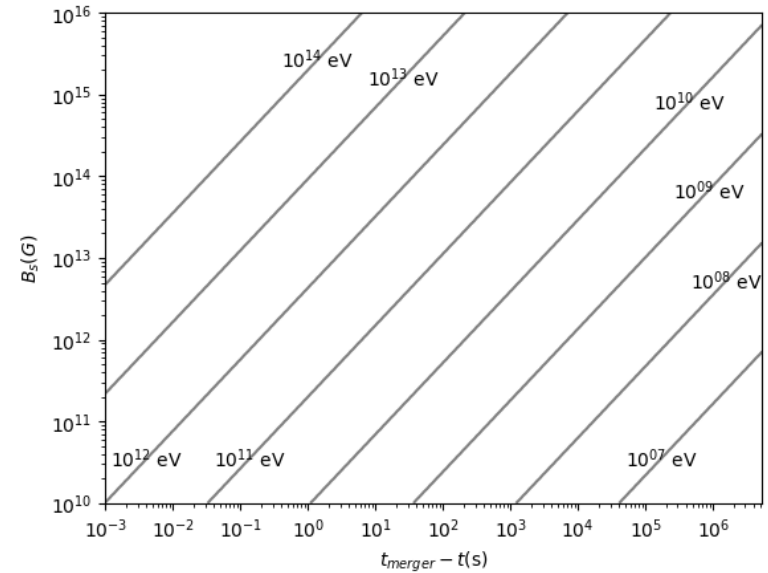
The cut off energy $\epsilon_{cut} = \frac{3}{2} c \hbar \frac{\gamma_L^3}{R_c}$



Likely detectable with CTA with warnings from 3G GW detectors (~2030s)

See Kalapotharakos et al. 2014, 2017, 2022

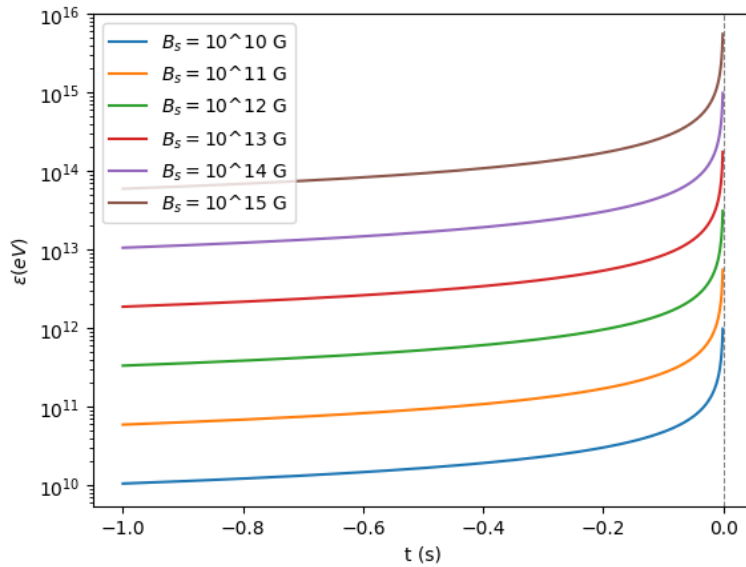
and Konstantinos Poster about the Fundamental Plane



Assuming emissions due to CR at the radiation reaction limit regime

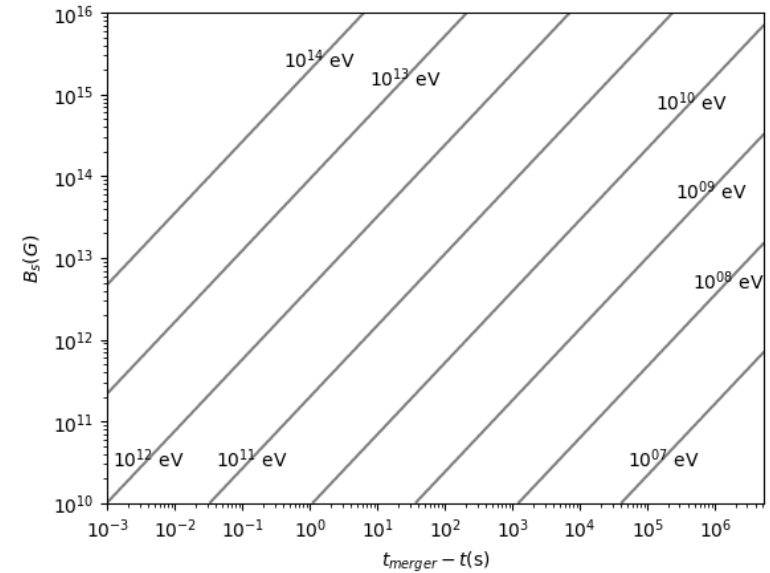
$$\frac{d\gamma_L}{dt} = \frac{q_e v E_{acc}}{m_e c^2} - \frac{2q_e \gamma_L^4}{3R_c^2 m_e c} = 0, \text{ with } E_{acc} \sim B_{LC} \text{ and } R_c \sim R_{LC}$$

The cut off energy $\epsilon_{cut} = \frac{3}{2} c \hbar \frac{\gamma_L^3}{R_c}$



See Kalapotharakos et al. 2014, 2017, 2022

and Konstantinos Poster about the Fundamental Plane



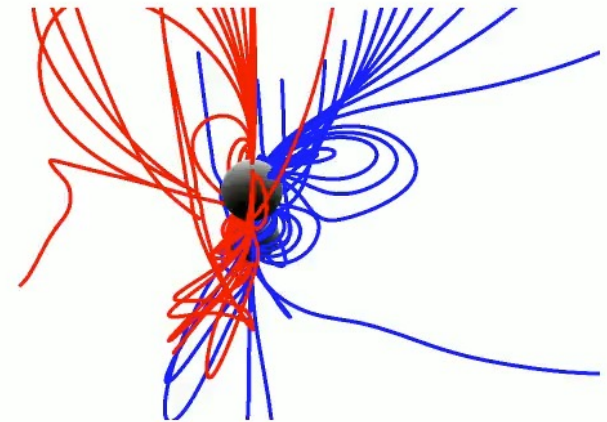
Likely detectable with CTA with warnings from 3G GW detectors (~2030s)

Introduction

- Interacting neutron star magnetospheres can release large amount of energy through electromagnetic emissions.

Piro 2012, Tsang et al. 2012, Pelenzuela et al. 2013, Lyutikov 2019, Sridhar et al. 2021, Most & Philippov 2020,2022, Cooper et al 2023

- Precursor signal to neutron star mergers.
- Echoes contributing to the GRB afterglows.
- Our goal is to explore the parameter space regarding the emission patterns.



For a review see Suvorov et al. 2024

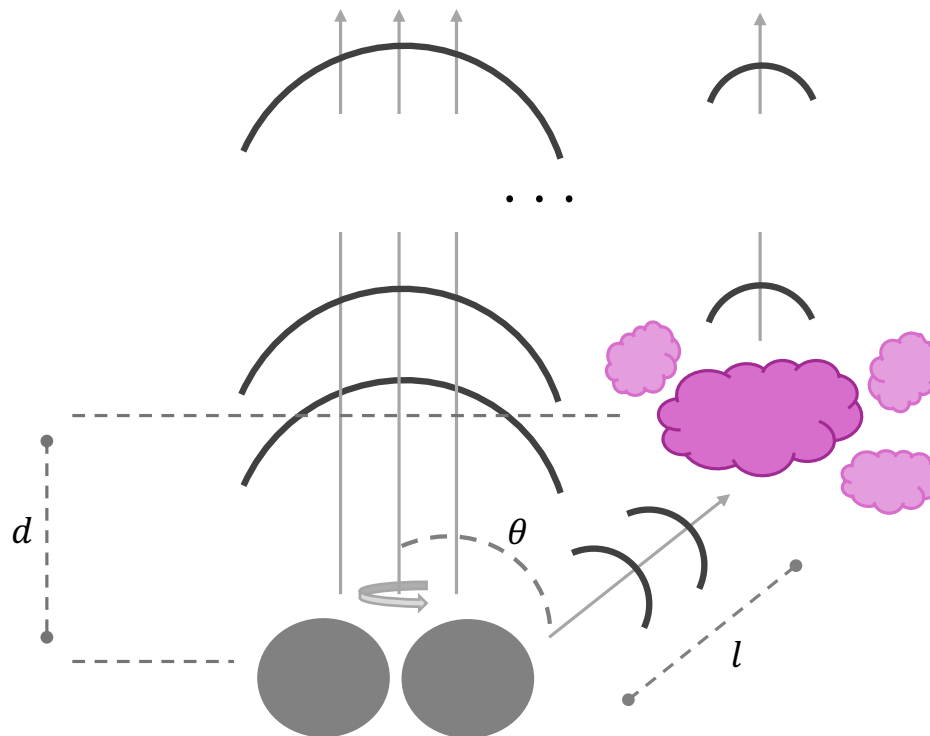
Echoes:



t : GRB Jet, GWs

$t + \Delta t$ later: afterglow rebrightening

$$\Delta t \simeq \frac{l - d}{c}$$
$$= \frac{l}{c} (1 - \cos\theta)$$

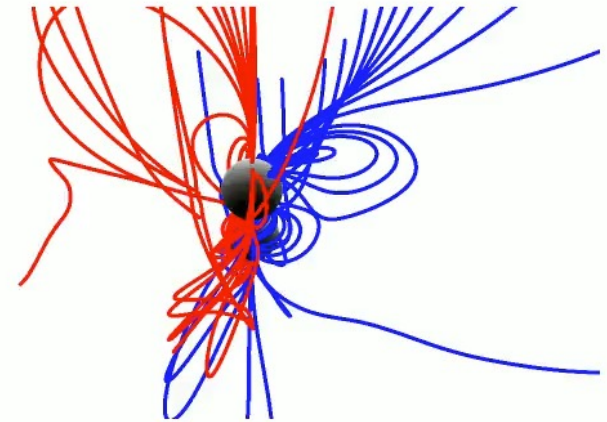


Introduction

- Interacting neutron star magnetospheres can release large amount of energy through electromagnetic emissions.

Piro 2012, Tsang et al. 2012, Pelenzuela et al. 2013, Lyutikov 2019, Sridhar et al. 2021, Most & Philippov 2020,2022, Cooper et al 2023

- Precursor signal to neutron star mergers.
- Creating echoes contributing to the GRB afterglows.
- Our goal is to explore the parameter space regarding the emission patterns.



For a review see Suvorov et al. 2024

Exploration parameters

- Inclination angle of the 1st star

$$a_1: 0^\circ, 45^\circ, 90^\circ$$

- Inclination angle of the 2nd star

$$a_2: 0^\circ, 45^\circ, 90^\circ$$

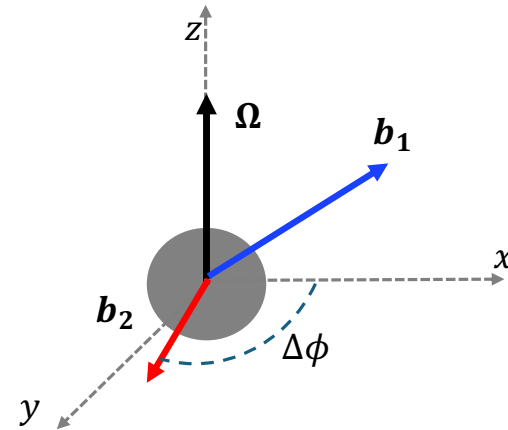
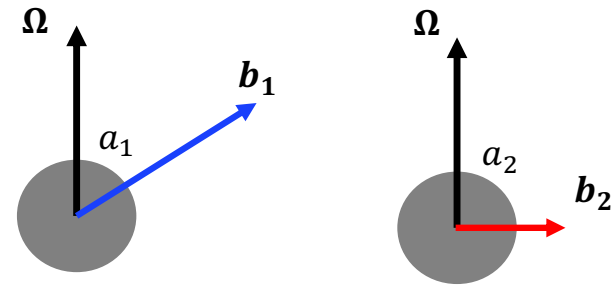
- Ratio of their magnetic moments

$$b_1/b_2: -10^2, -10, -10^{\frac{1}{2}}, -1, 1, 10^{\frac{1}{2}}, 10, 10^2$$

- Spin phase difference*

$$\Delta\phi: 0^\circ, 45^\circ, 90^\circ$$

**Only for some b_1/b_2 ratio*



Our approach

- In the limit of irrotational stars: $\omega_1 = \omega_2 = 0$.
- Dipole magnetic moments.
- Two neutron stars with $R_* = 12\text{km}$ and $M_* = 1.4 M_\odot$

- $\Omega = \left(1 - \frac{t}{t_s}\right)^{-3/8}$ with t_s the time to the merger.

Medvedev & Loeb 2013, Peters 1964

- $\frac{r}{r_i} = \left(\frac{\Omega}{\Omega_i}\right)^{-2/3}$

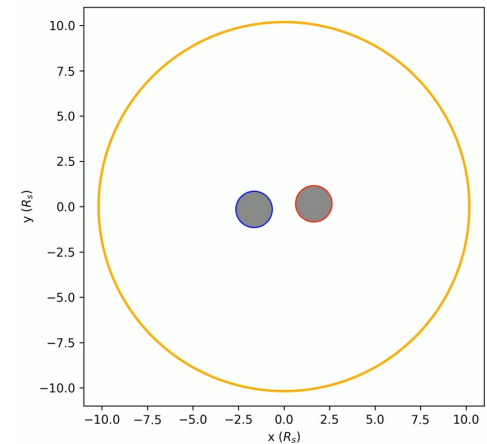
Our approach

- In the limit of irrotational stars: $\omega_1 = \omega_2 = 0$.
- Dipole magnetic moments.
- Two neutron stars with $R_* = 12\text{km}$ and $M_* = 1.4 M_\odot$

- $\Omega = \left(1 - \frac{t}{t_s}\right)^{-3/8}$ with t_s the time to the merger.

- $\frac{r}{r_i} = \left(\frac{\Omega}{\Omega_i}\right)^{-2/3}$

Medvedev & Loeb 2013, Peters 1964



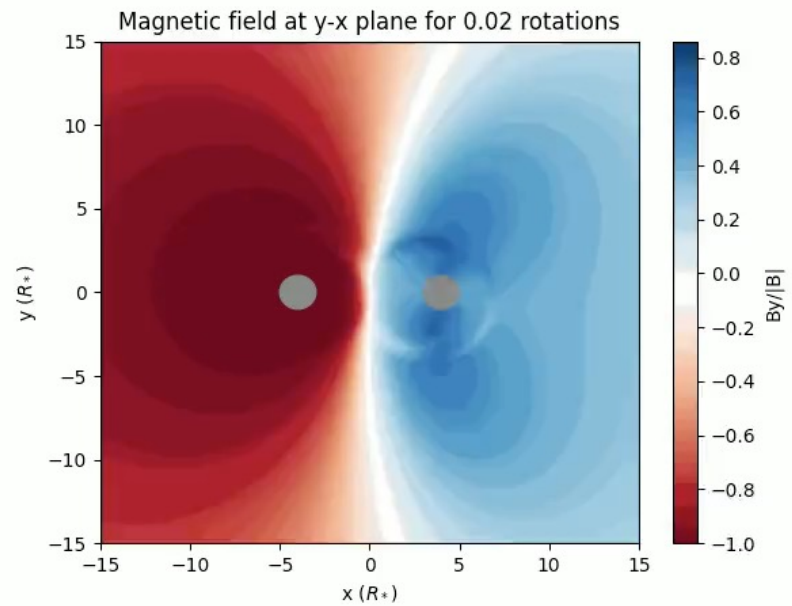
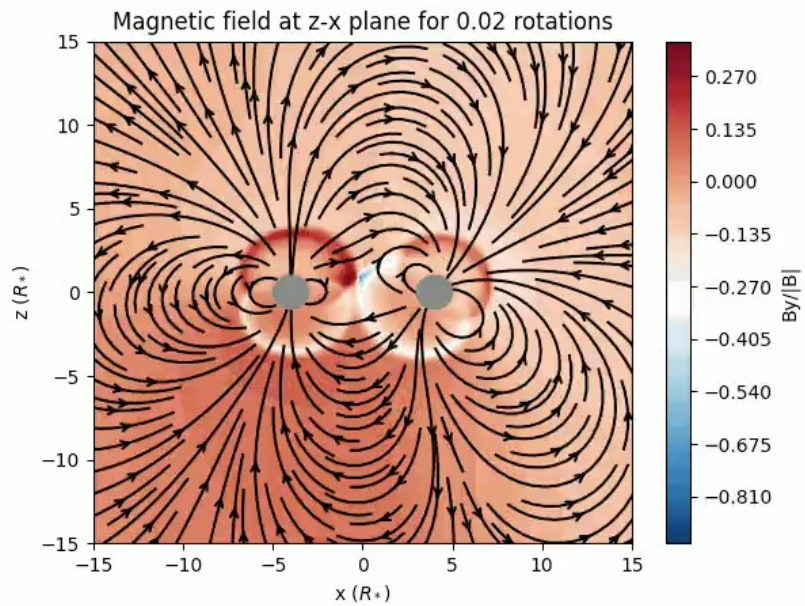
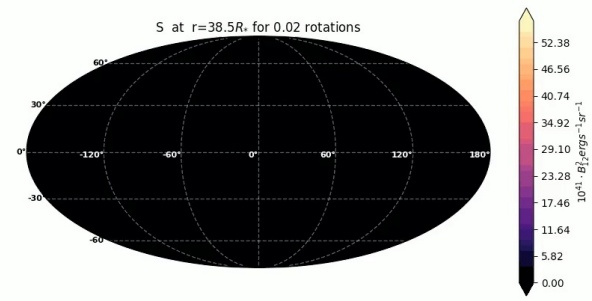
Our approach

- Ignoring tidal distortions and crust cracking.
- Numerical 3D force-free MHD simulation of global magnetosphere.

Kalapothisarakos et al. (2012)

- Grid size: $\sim 500 \times 500 \times 500$.
- Initial separation $d = 3.3R_*$.
- Simulations last for 7.7 ms (**4 rotations**) until neutron stars touch.

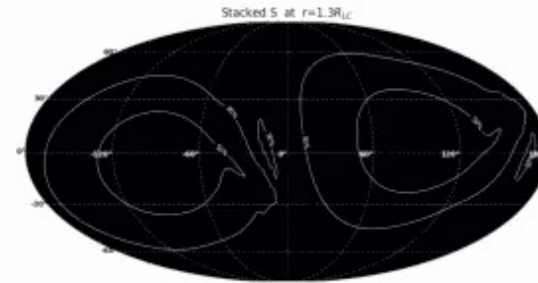
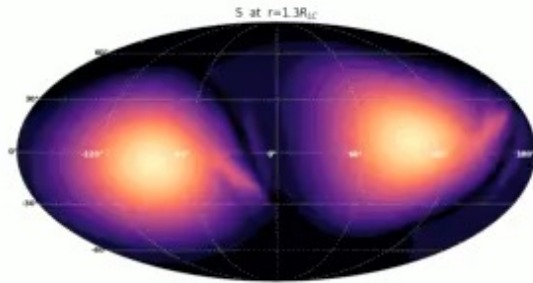
Example Run for ~50 Rotations



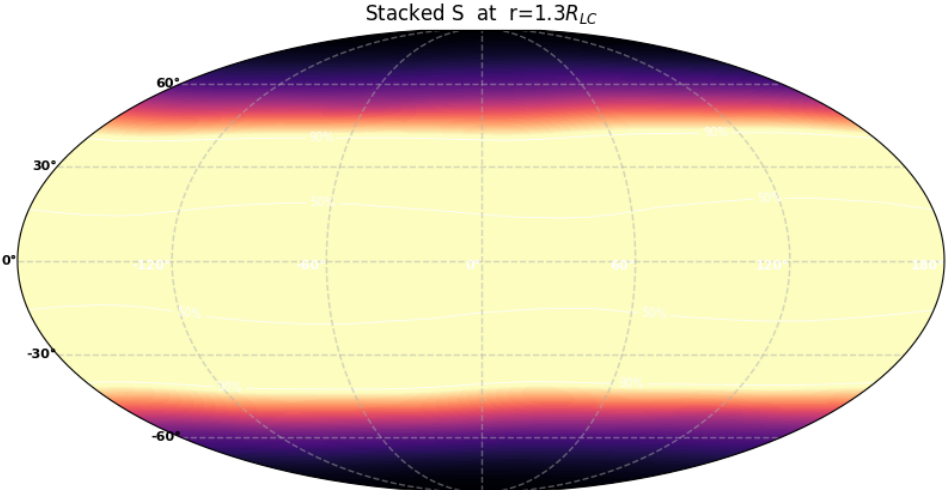
Single Isolated Pulsar

Energy Flux: $L = \int \mathbf{S} \cdot d\mathbf{A}$, where $\mathbf{S} = \frac{c}{4\pi} \mathbf{E} \times \mathbf{B}$, the Poynting vector.
the directional energy flux per time and area unit

For dipole magnetic moment: $L \propto \mu^2 \omega^4 / c^3$

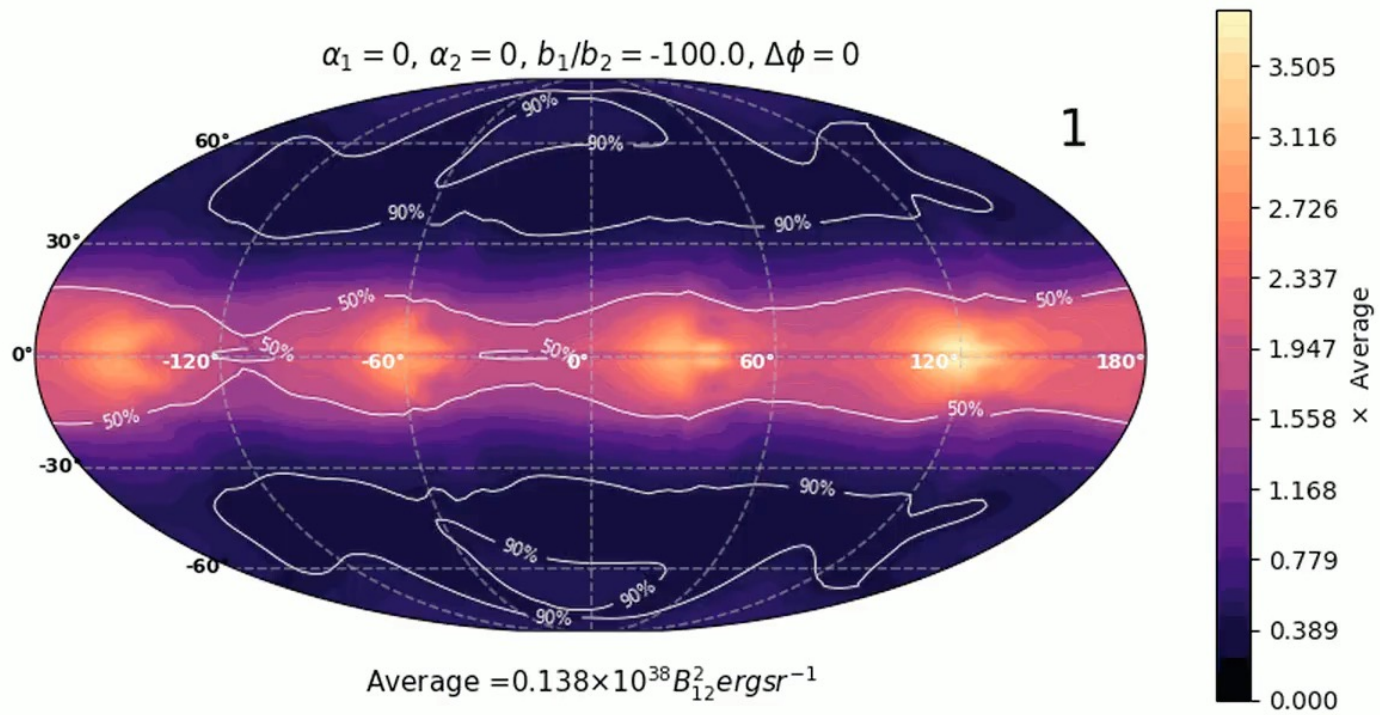


Single Isolated Pulsar



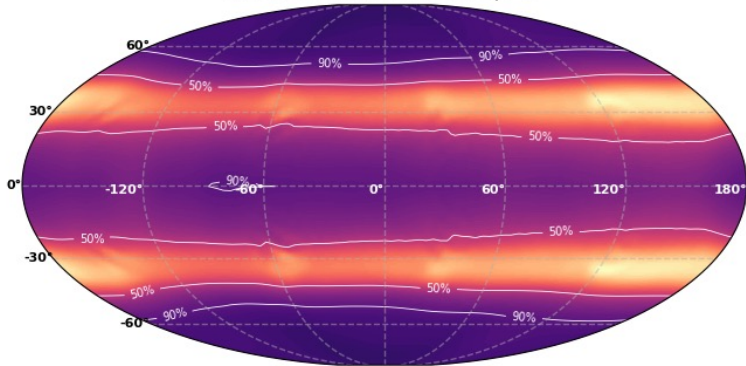
...only θ matters!

144 cases explored

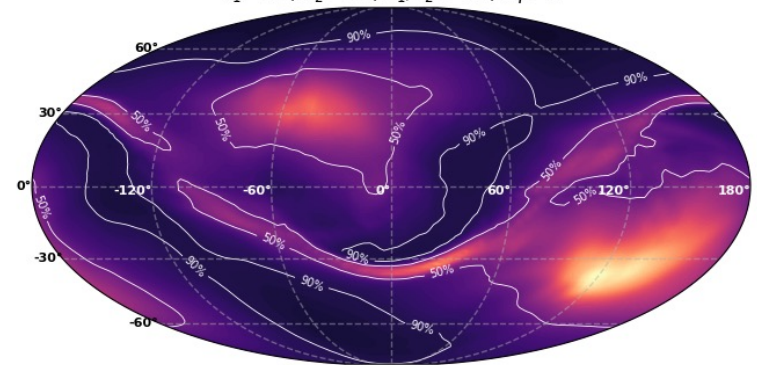


Both θ and ϕ matters

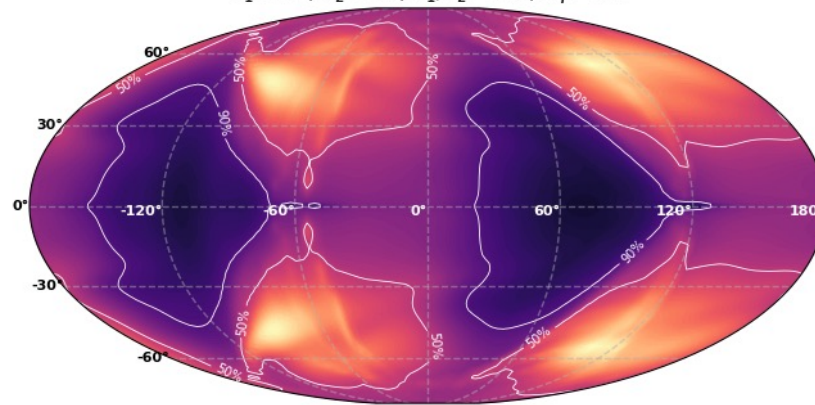
$\alpha_1 = 0, \alpha_2 = 0, b_1/b_2 = 10, \Delta\phi = 0$



$\alpha_1 = 45, \alpha_2 = 90, b_1/b_2 = -10, \Delta\phi = 0$

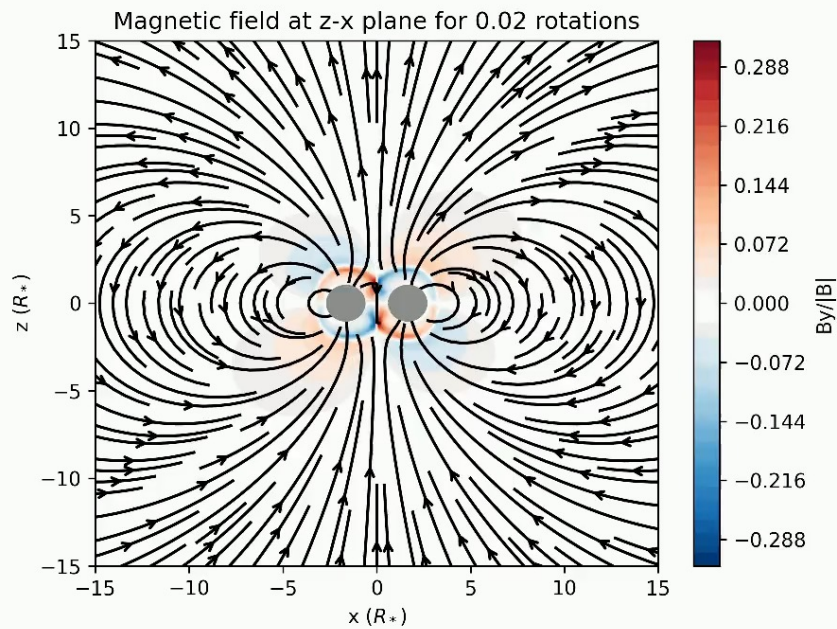


$\alpha_1 = 90, \alpha_2 = 90, b_1/b_2 = -10, \Delta\phi = 90$



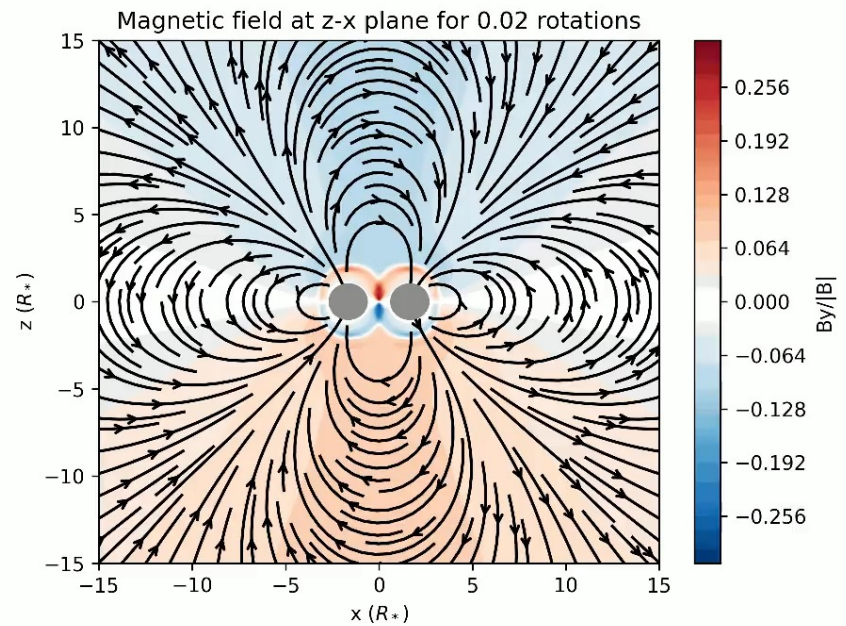
Aligned 

$$\alpha_1 = 0^\circ, \alpha_2 = 0^\circ, b_2 = b_1$$

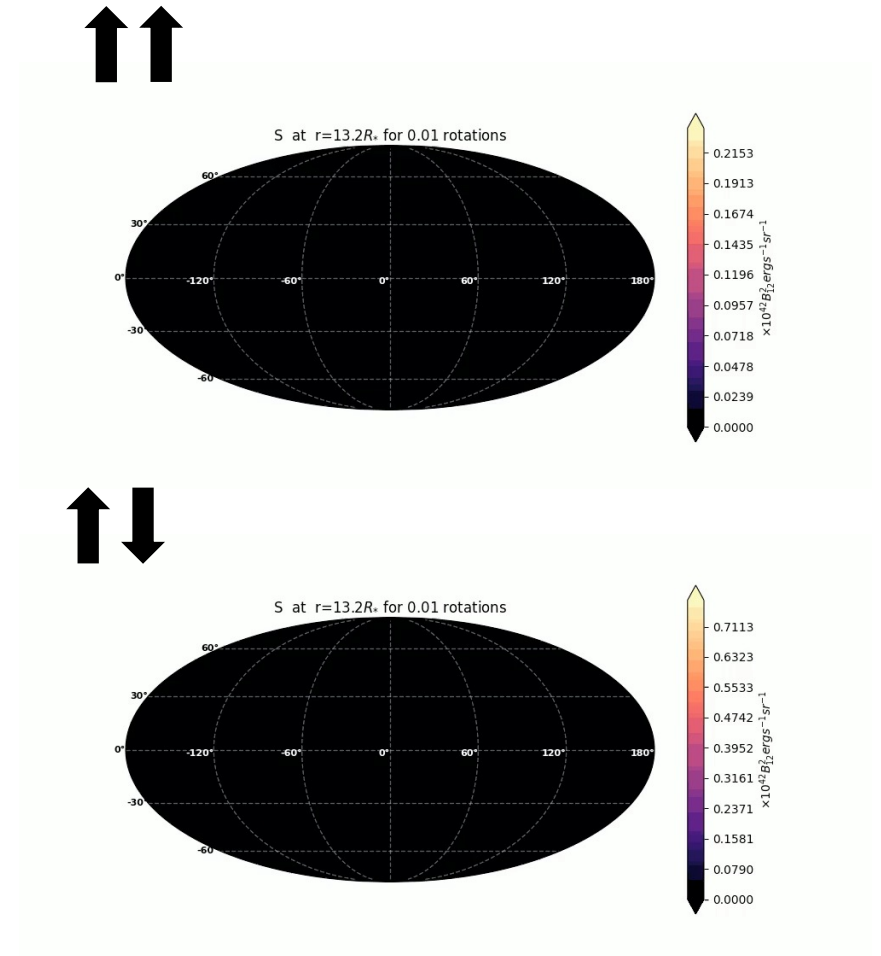
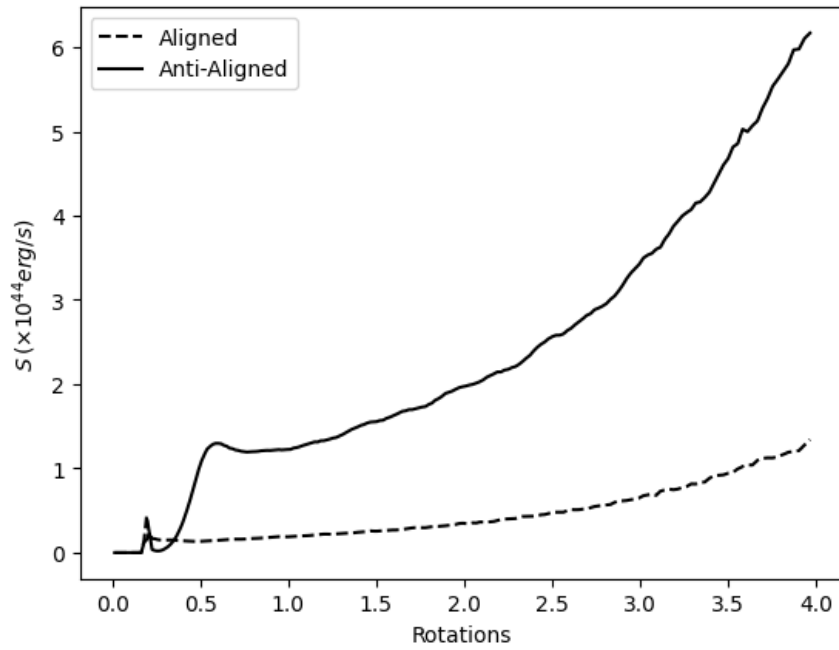


Anti-Aligned 

$$\alpha_1 = 0^\circ, \alpha_2 = 0^\circ, b_2 = -b_1$$

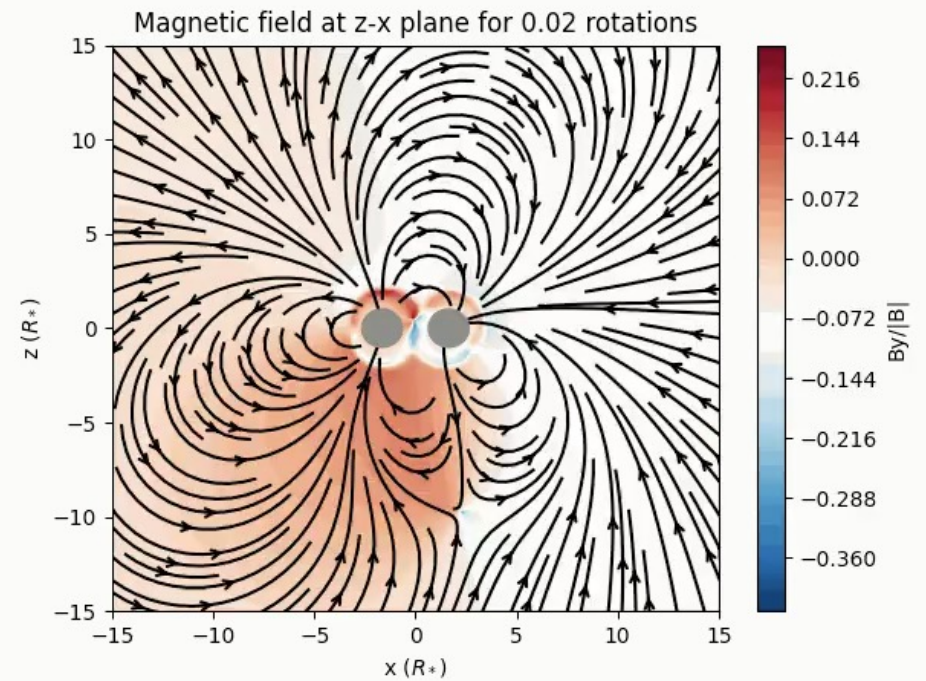
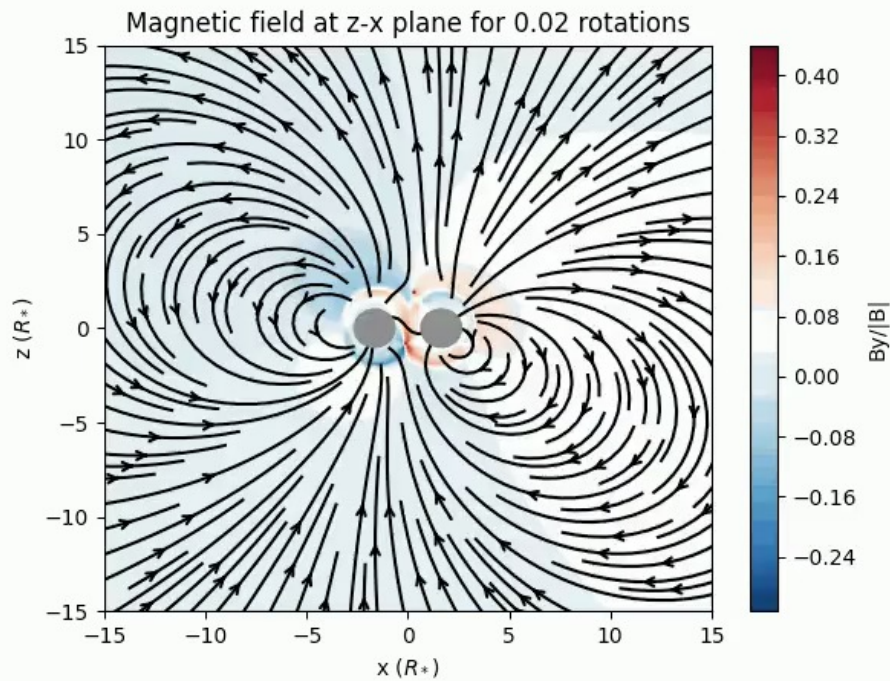




Aligned vs Anti-Aligned





Consistent with Most & Philippov 2020,2022 and Pelenzuella et al. 2013

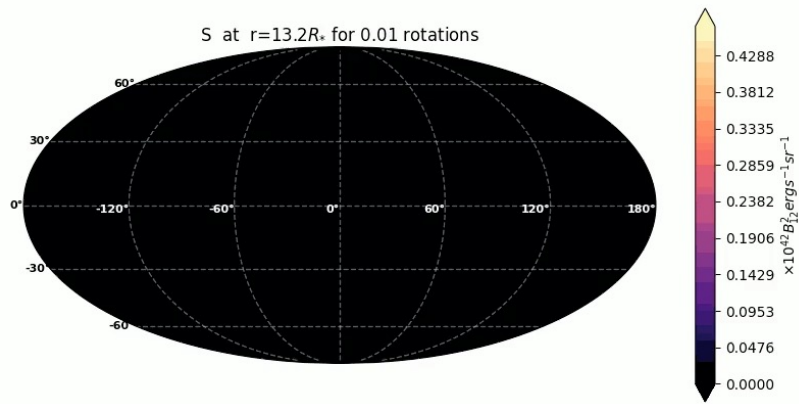
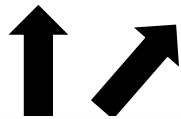
Inclination angles



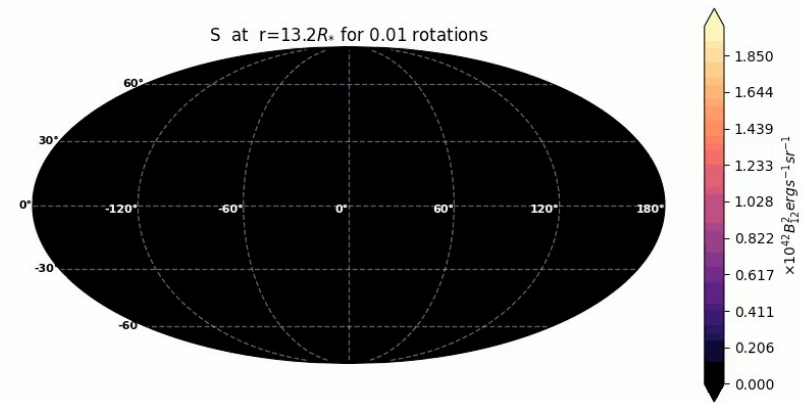
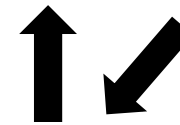


 $\alpha_1 = 0^\circ, \alpha_2 = 45^\circ, b_2 = b_1$

$\alpha_1 = 0^\circ, \alpha_2 = 45^\circ, b_2 = -b_1$



Inclination angles



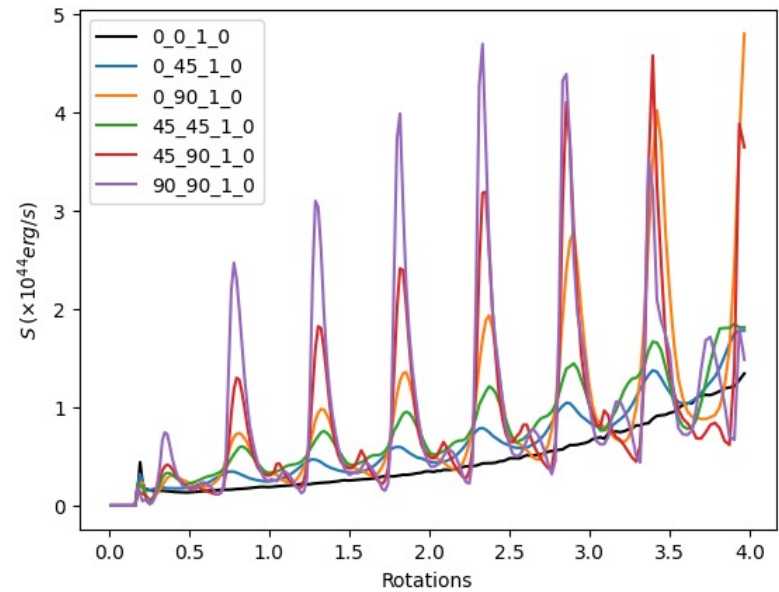
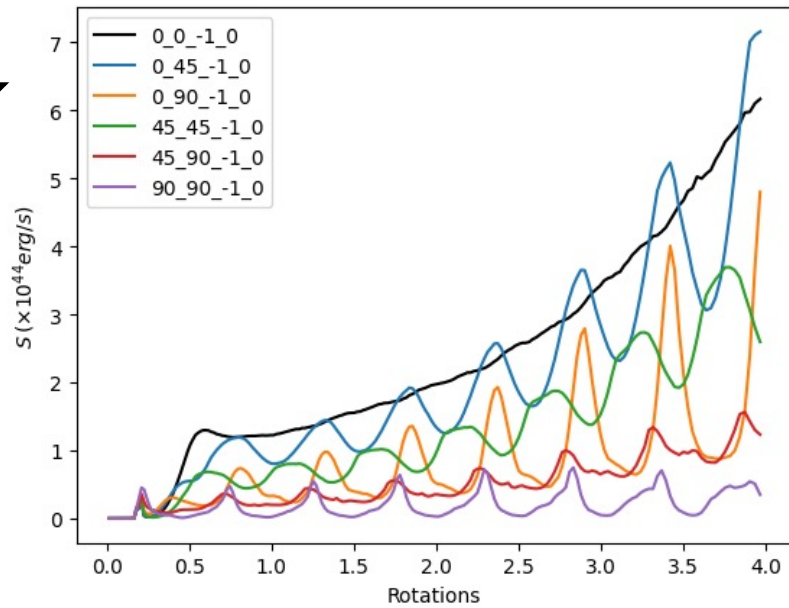
$$\alpha_1 = 0^\circ, \alpha_2 = 45^\circ, b_2 = b_1$$



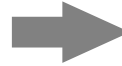
$$\alpha_1 = 0^\circ, \alpha_2 = 45^\circ, b_2 = -b_1$$

Effect of inclination angles:

Aligned stars with inclination give more flux



Anti - aligned stars with inclination give less flux



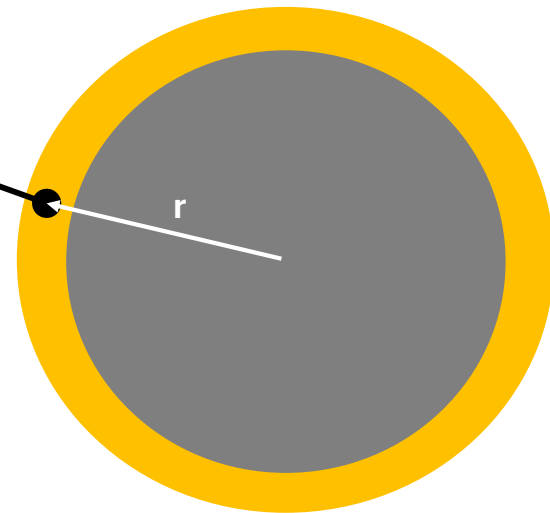
Torques and kicks

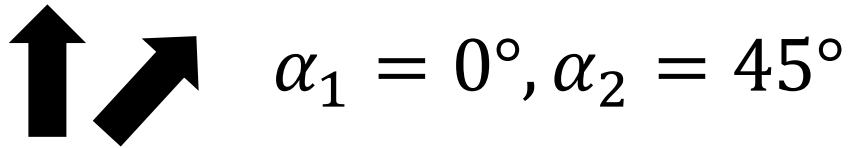
Due to asymmetric pointing flux torques and kicks might develop that will might affect the orbit or magnetic alignment.

$$\mathbf{f} = \nabla \cdot \boldsymbol{\sigma} - \frac{1}{c^2} \frac{\partial \mathbf{S}}{\partial t}$$

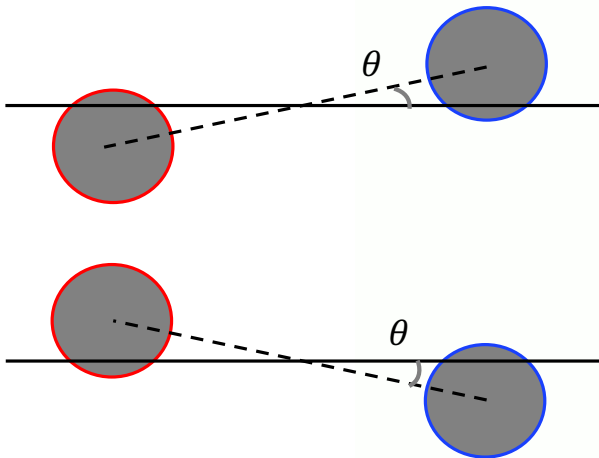
Where σ is the Maxwell tensor:

$$\sigma_{ij} = \frac{1}{4\pi} \left(E_i E_j + B_i B_j - \frac{1}{2} (E^2 + B^2) \delta_{ij} \right)$$





$F=6.20e+28N,$ $F_x=2.48e+28N,$ $F_y=-3.07e+28N,$ $F_z=4.77e+28N,$
 $F1=4.18e+29N,$ $F1_x=-3.63e+29N,$ $F1_y=1.14e+29N,$ $F1_z=-1.73e+29N,$
 $F2=4.69e+29N,$ $F2_x=3.88e+29N,$ $F2_y=-1.45e+29N,$ $F2_z=2.21e+29N,$

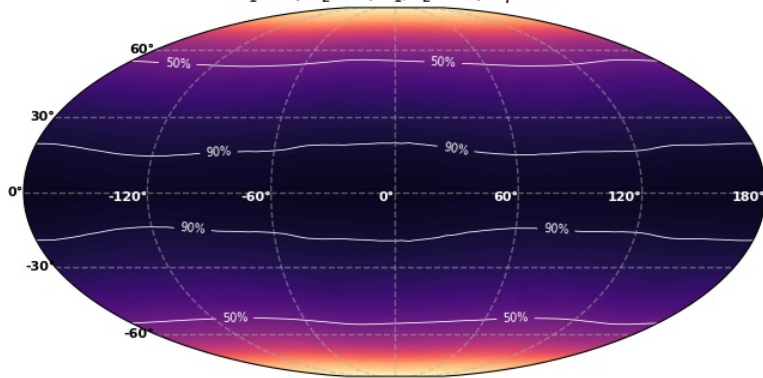


For $3.16 \times 10^{15} G$: $\frac{F_{EM}}{F_G} \leq 10^{-3}$

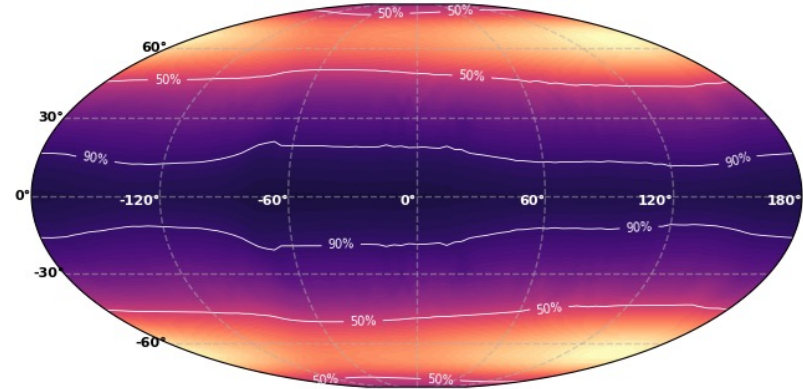
Potentially important for 3G GW detectors.

Effect of b_1/b_2 ratio:

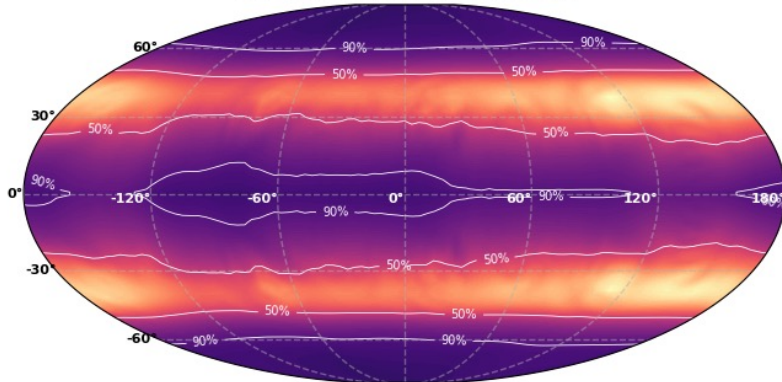
$\alpha_1 = 0, \alpha_2 = 0, b_1/b_2 = -1, \Delta\phi = 0$



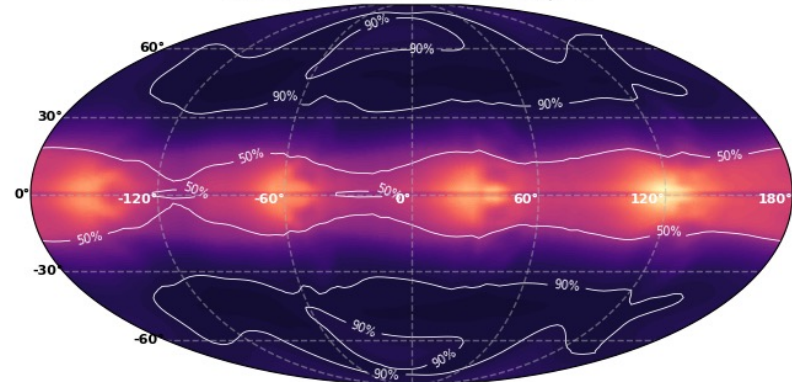
$\alpha_1 = 0, \alpha_2 = 0, b_1/b_2 = -3.16, \Delta\phi = 0$



$\alpha_1 = 0, \alpha_2 = 0, b_1/b_2 = -10, \Delta\phi = 0$

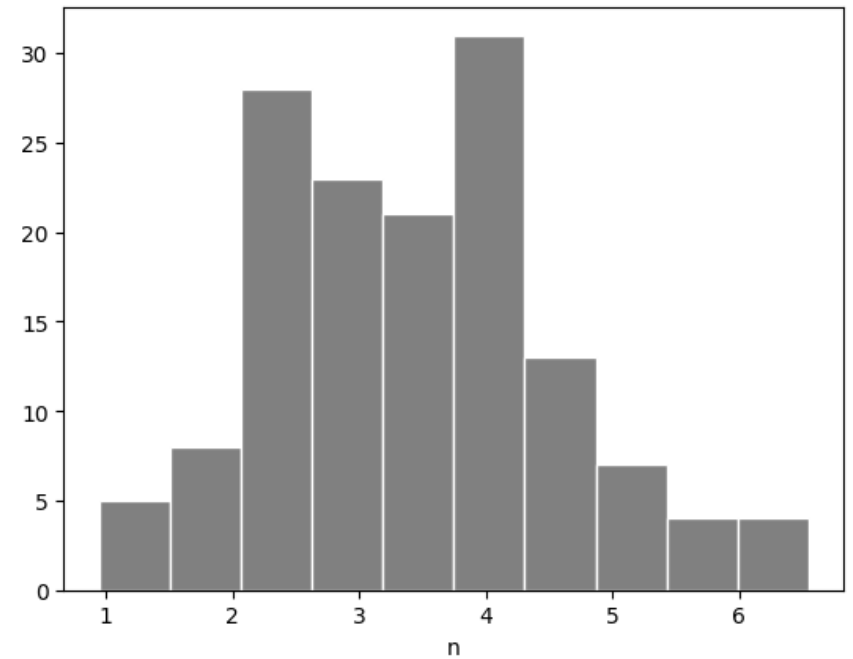
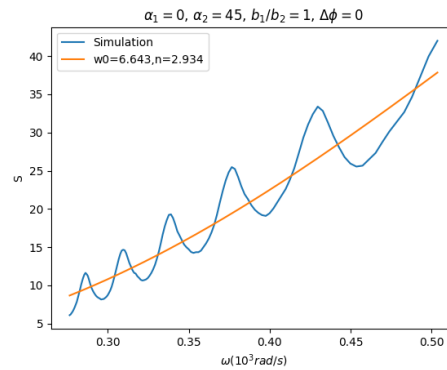
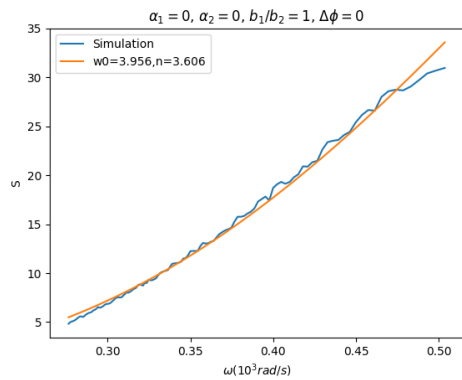


$\alpha_1 = 0, \alpha_2 = 0, b_1/b_2 = -100.0, \Delta\phi = 0$



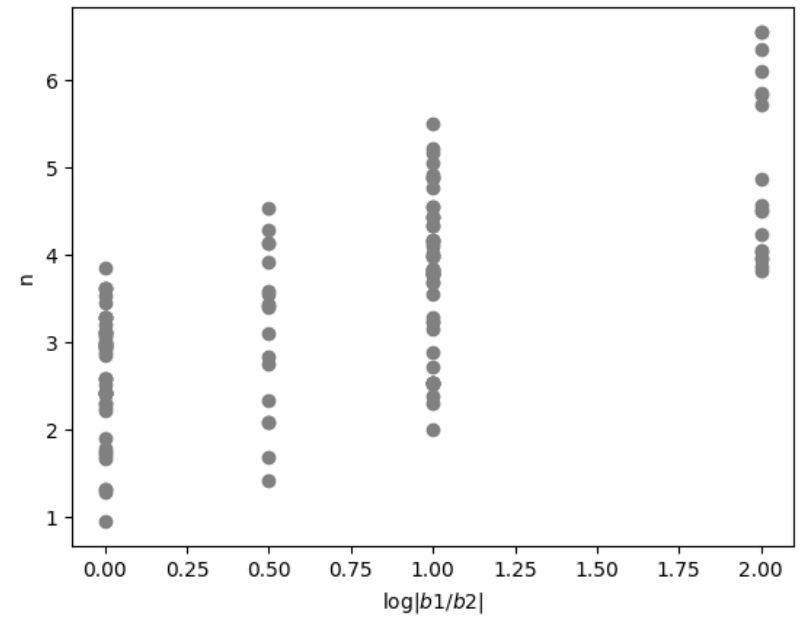
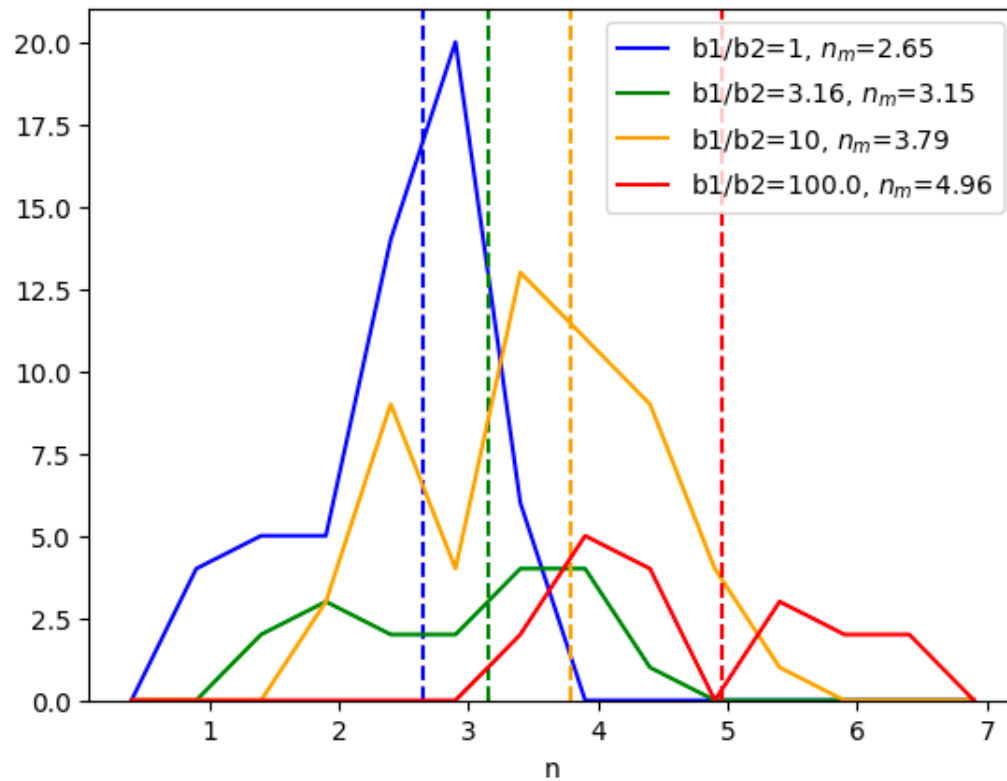
Poynting flux dependency on Ω .

$$L \sim \Omega^n \cdot f(\alpha_1, \alpha_2, \Delta\phi)$$



For single isolated pulsar n=4

Poynting flux dependency on Ω .



γ – rays patterns

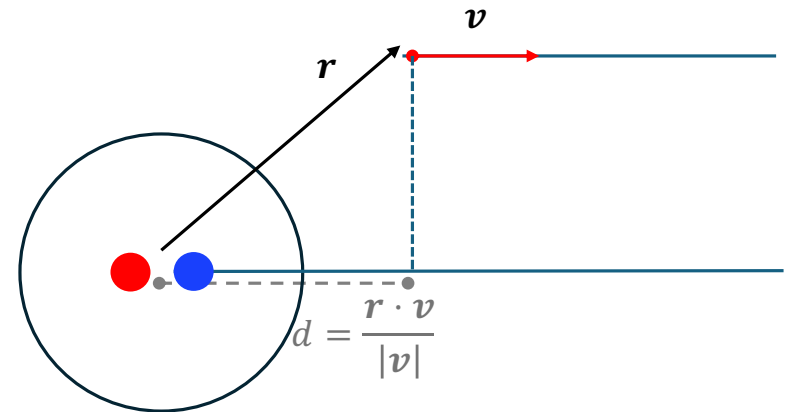
- Acceleration of particles from $E_{||}$ and emission of high energy photons.

- $$\mathbf{E}_{||} = \frac{c|(\nabla \times \mathbf{B})_{||}|}{4\pi\sigma}$$

- Emissions happen in the region outside the orbital light cylinder at points where $\mathbf{E}_{||}$ is high.

- Emission from each point contribute at the direction of particle's velocity on the appropriate time considering the time delay.

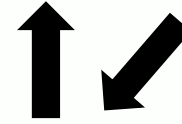
- Luminosity is considered as a function of the Poynting flux with a distance depended efficiency.



$$\mathbf{v} = \frac{\mathbf{E} \times \mathbf{B} \pm (B_0 \mathbf{B} + E_0 \mathbf{E})}{B^2 + E_0^2}$$

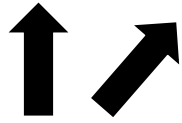
Gruzinov 2012
 Kalapotharakos et al. 2014
 Brambilla et al. 2015

Rotation:0.6



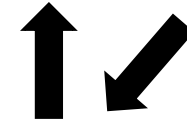
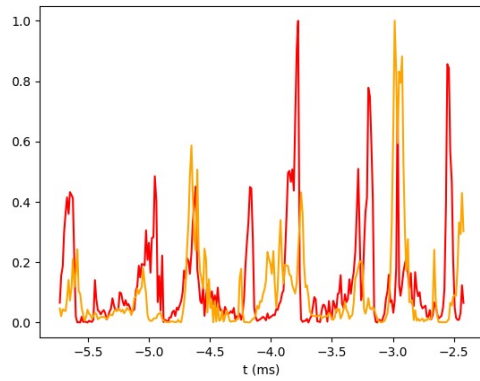
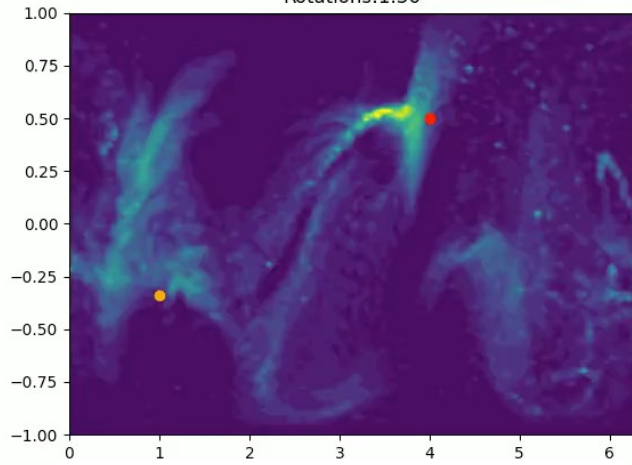
$$\alpha_1 = 0^\circ, \alpha_2 = 45^\circ, b_1 = -b_2$$





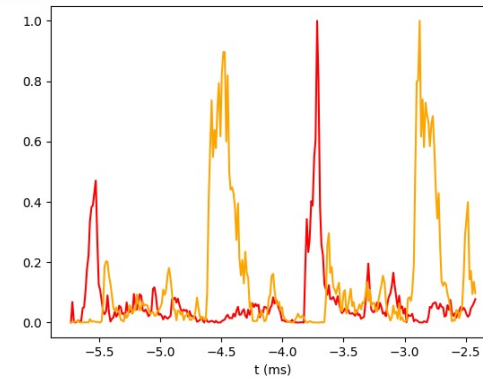
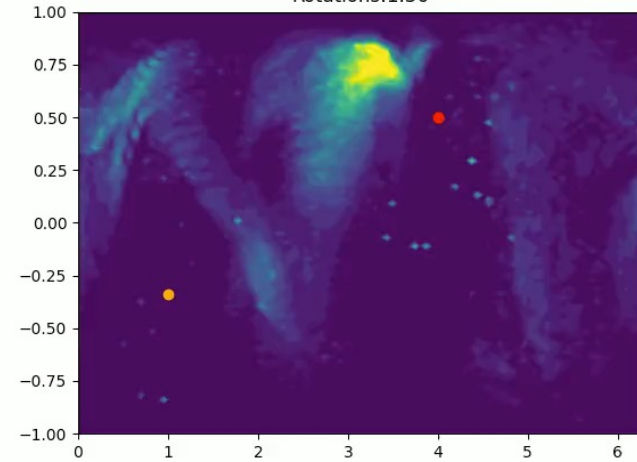
$$\alpha_1 = 0^\circ, a_2 = 45^\circ, b_1 = b_2$$

Rotations:1.36



$$\alpha_1 = 0^\circ, a_2 = 45^\circ, b_1 = -b_2$$

Rotations:1.36



Conclusions

- We explored the parameter space of binary neutron star mergers and their precursor emissions.

144 cases

- The flux varies a lot:

$$\sim 10^{42} - 10^{45} B_{12}^2 \text{ ergs}^{-1}$$

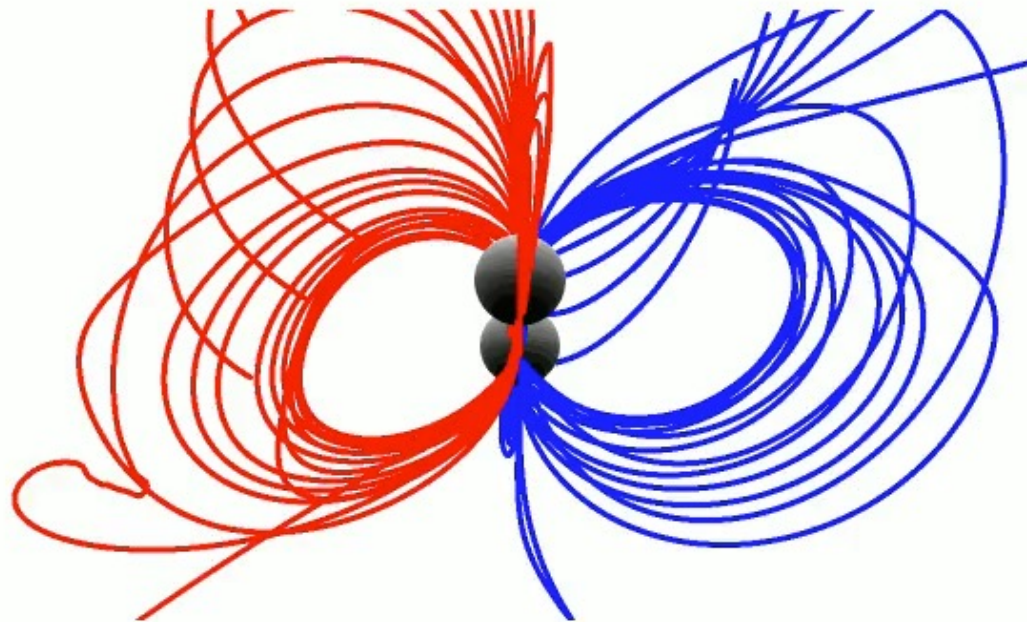
$$\text{or } \sim 10^{49} - 10^{52} B_{15.5}^2 \text{ ergs}^{-1}$$

- If $L_\gamma < 10\% L_{\text{Poynting}}$ γ –rays might be detectable from CTA.
- Non uniform flux distribution over the sky.

50% in less than ~1/3 of the sky

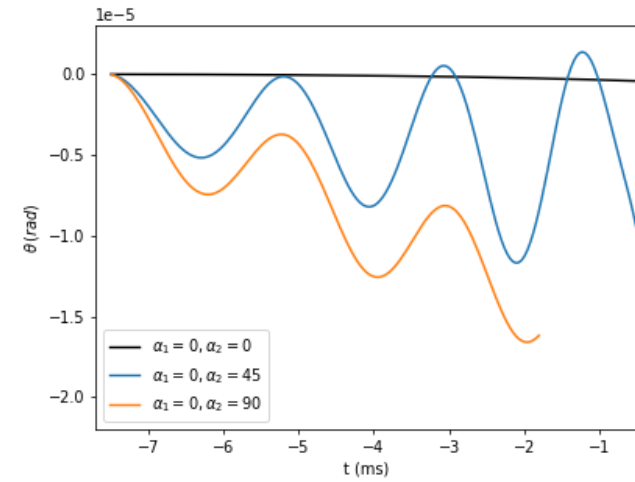
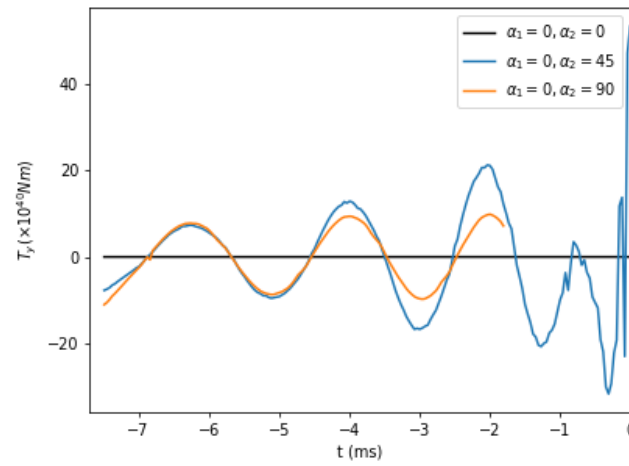
- The Poynting flux dependency on Ω varies.
- Having inclination angles results in different behavior for the aligned and anti-aligned cases and potentially implications on the orbits.

Thank you!



dimitrios.skiathas@nasa.gov

Aligned



Anti-Aligned

