

A Little History, but... in Astronomy its Resolution, Resolution, Resolution

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Gamma Ray Pair Conversion

High Energy



Fig. 2: Photon cross-section σ in lead as a function of photon energy. The intensity of photons can be expressed as $I = I_0 \exp(-\sigma x)$, where x is the path length in radiation lengths, (Review of Particle Properties, April 1980 edition).

Gamma Ray Z = 74Tungsten **e+ QED Process Tungsten Conversion Foils Position Measuring Detectors** • Measured Track co-ordinates \frown **Plastic Scintillation** counters to veto entering charged particles **Total Absorption Calorimeter to** measure gamma ray energy

High Electric

Field

Pair

Energy 1.0 **Splitting Function** 0.9 0.8 0.7 10⁹ 2 × 10 0.6 0.5 4×10^{7} 0.4 2×10^{7} 0.3 $E = 10^{\circ}$ 0.2 E. (= 10 MeV) 0.1 0.4 0.5 0.8 0.9 0.2 0.3 0.6 0.7 0.1 E_{e+}/E_{γ}

Sermi

Opening Angle



At 100 MeV $\theta_{\text{Open}} \sim 1^{\circ}$



From Rossi, High Energy Particles, 1952





 1967-1968, OSO-3 Detected Milky Way as an extended γray source

621 γ-rays

- 1972-1973, SAS-2,
 ~8,000 γ-rays
- 1975-1982, COS-B

orbit resulted in a large and variable background of charged particles

~200,000 **γ**-rays

• 1991-2000, EGRET

Large effective area, good PSF, long mission life, excellent background rejection

 $>1.4 \times 10^{6} \gamma$ -rays





Evolution of GLAST

• April, 1991 CGRO (with EGRET on board) Shuttle Launch

• May, 1992 NASA SR & T Proposal Cycle

1. Select the Technologies



3. Pick the Rocket

Large area SSD systems and CsI Calorimeters resulted from SSC R&D

Original GISMO 1 Event Displays from the first GLAST simulations

2. Make it Modular



4. Fill-it-up!

Another lesson learned in the 1980's: monolithic detectors are inferior to Segmented detectors



Delta II (launch of GP-B)

Cheap, reliable Communication satellite launch vehicle

Diameter sets transverse size

Lift capacity to LEO sets depth of Calorimeter Rocket Payload Fairing



Technology of Choice: Solid State

Silicon Strip Detector Principle

Custom Integrated Circuits (ASICs)



GLAST has 884736 channels. Total Tracker Power = 160 Watts!



Tracker Design and Analysis



Diversion: Review of Covariance



Results:

Gamma-ray Space Telescope

$$x^{2}\left(\frac{\cos^{2}(\theta)}{a^{2}} + \frac{\sin^{2}(\theta)}{b^{2}}\right) + 2xy\cos(\theta)\sin(\theta)\left(\frac{1}{a^{2}} - \frac{1}{b^{2}}\right) + y^{2}\left(\frac{\sin^{2}(\theta)}{a^{2}} + \frac{\cos^{2}(\theta)}{b^{2}}\right) = 1$$

Rotations mix x & y. Major & minor axis plus rotation angle θ complete description. Error Ellipse described by Covariance Matrix:

Distance between a point with an error and another point measured in σ 's: $(n\sigma)^2 = r^T C^{-1} r$ where $r = (\vec{x} - \vec{x})$ and $r\sigma \sim (n\sigma)^2 = r^T C^{-1} r$ where $r = (\vec{x} - \vec{x})$ and $C^{-1} = Inverse(C) = \begin{bmatrix} C_{xx}^{-1} & C_{xy}^{-1} \\ C_{yx}^{-1} & C_{yy}^{-1} \end{bmatrix}$ and $C_{xy}^{-1} = C_{yx}^{-1}$ } Simply weighting the distance by $1/\sigma^2$





Multiplying it out gives:

$$(n\sigma)^{2} = r^{T}C^{-1}r = (x, y)\begin{bmatrix} C_{xx}^{-1} & C_{xy}^{-1} \\ C_{xy}^{-1} & C_{yy}^{-1} \end{bmatrix} \begin{pmatrix} x \\ y \end{pmatrix} = x^{2}C_{xx}^{-1} + 2xyC_{xy}^{-1} + y^{2}C_{yy}^{-1}$$

Where I take $\bar{x} = 0$ without loss of generality. This is the equation of an ellipse! Specifically for a 1σ error ellipse ($n\sigma = 1$) we identify:

 $C_{xx}^{-1} = \frac{\cos^2(\theta)}{a^2} + \frac{\sin^2(\theta)}{b^2} \qquad C_{yy}^{-1} = \frac{\sin^2(\theta)}{a^2} + \frac{\cos^2(\theta)}{b^2} \quad C_{xy}^{-1} = \sin(\theta)\cos(\theta)(\frac{1}{a^2} - \frac{1}{b^2})$ and $C^{-1} = \frac{1}{\det(C)} \begin{bmatrix} C_{yy} & -C_{xy} \\ -C_{xy} & C_{xx} \end{bmatrix}$ where $\det(C) = (C_{xx}C_{yy} - C_{xy}^2)$ The correlation coefficient is defined as: $R^2 = \frac{C_{xy}^2}{C_{yy}C_{yy}}$

Summary: The *Covariance Matrix* describes an ellipse where the major and minor axis and the rotation angle map directly onto its components!





Let the fun begin! To disentangle the two descriptions consider



Also det(C) yields (with a little algebra & trig.): $a \cdot b = \sqrt{\det(C)} = \frac{Area}{\pi}$

Now we're ready to continue the story of GLAST (Fermi-LAT)!



Multi-variate Analysis: Kalman Filter

C_{xx}

C_{sx}

The Kalman filter process is a successive approximation scheme to estimate parameters

Simple Example: 2 parameters - *intercept* and *slope*: $x = x_0 + S_x * z$; $P = (x_0, S_x)$

Errors on parameters $x_0 \& S_x$: Covariance Matrix C =

$$\begin{array}{c} \mathbf{C}_{xs} \\ \mathbf{C}_{ss} \end{array} \right) \qquad \begin{array}{c} \mathbf{C}_{xx} = \langle (\mathbf{x} \cdot \mathbf{x}_m) (\mathbf{x} \cdot \mathbf{x}_m) \rangle \\ & \text{In general} \\ \mathbf{C} = \langle (\mathbf{P} \cdot \mathbf{P}_m) (\mathbf{P} \cdot \mathbf{P}_m)^T \rangle \end{array}$$

Propagation:

$$x(k+1) = x(k) + S_x(k) * (z(k+1)-z(k))$$

$$P(k+1) = F(\delta z) * P(k)$$
 where

$$\mathbf{F}(\delta \mathbf{z}) = \begin{pmatrix} 1 & \mathbf{z}(\mathbf{k}+1) \cdot \mathbf{z}(\mathbf{k}) \\ 0 & 1 \end{pmatrix}$$
$$\mathbf{C} \ (\mathbf{k}+1) = \ \mathbf{F}(\delta \mathbf{z}) \ ^{*}\mathbf{C}(\mathbf{k}) \ ^{*} \ \mathbf{F}(\delta \mathbf{z})^{\mathrm{T}} + \mathbf{Q}(\mathbf{k})$$









Now its repeated for the k+2 planes and so - on. This is called FILTERING - each successive step incorporates the knowledge of previous steps as allowed for by the NOISE and the aggregate sum of the previous hits.





Really a game of "Do you get out what you put in?"

1 GeV Muons used for testing Energy for Kalman Fit = MC Energy Results: Large Tail on χ^2 Solution: Include energy losses in Tracker (Bethe-Bloch)

Next: χ^2 Depends on Angle





Large Angles ⇒ too Narrow!!

...Suspect Meas. Errors





Gaussian Equivalent σ for a Square Distribution



Success! χ² Distributions – Text-Book!



Gamma-ray Space Telescope

 $\cos(\theta) = -1$

$$\langle \sigma_{\rm FIT} \rangle = 3.4 \, \rm{mrad}$$



 $-1 < \cos(\theta) < 0$

 $\langle \sigma_{\rm FIT} \rangle = 4.0 \text{ mrad}$











2. Covariant Averaging of Tracks

Multivariate Averaging:
$$P_{Pair} = \frac{C_1^{-1} \cdot P_1 + C_2^{-1} \cdot P_2}{C_1^{-1} + C_2^{-1}}$$

 $P_{Pair} = (C_1^{-1} + C_2^{-1})^{-1} \cdot (C_1^{-1} \cdot P_1 + C_2^{-1} \cdot P_2)$
 $C_{Pair} = (C_1^{-1} + C_2^{-1})^{-1}$

where P_i are the parameter vectors of the combination(Pair) and tracks (P_1 and P_2) and C_i are the covariance matrices

And

$$\chi^{2} = (P_{1} - P_{Pair})^{T} C_{Res1}^{-1} (P_{1} - P_{Pair}) + (P_{2} - P_{Pair})^{T} C_{Res2}^{-1} (P_{2} - P_{Pair})$$

where

$$C_{\text{Res1,2}} = C_{1,2} - C_{\text{Pair}}$$

The parameter vectors **P** are (x, S_x, y, S_y)





Neutral Energy Concept



Sometimes at the start of the shower the charge pair does not well reflect the direction of the incoming photon.

Bremstrahlung can cause much (most) of the energy to windup in photons.

The Calorimeter centroid is a measure of where these photons impact the calorimeter.

A "Neutral Energy" direction can be inferred by connecting the found vertex with the Cal. Centroid.

One can determine the covariant error matrix for this inferred direction by using the errors on the centroid location.

By having an imaging calorimeter, Fermi-LAT is the first Gamma Ray instrument able to do this! Sermi Where does the Charged Solution go Wrong?

At energies < 1 GeV only the first 2 Tracker Hits determine the direction



Gamma-ray Space Telescope

Internal and *External*

Brems. distort direction *This is in addition to multiple scattering*

Expect effect to be more sever in 18% radiators

(Thick - 18%, Thin - 3%)

Brems. in 2nd and lower decks doesn't effect direction

Due to *Internal* Brems. ratio of effected events (*Thin : Thick Decks*) will be < ratio of Rad. Lens.







<u>Is this Effect Present in the Monte Carlo?</u>

Define θ_{NC} to be the angle between The VTX (Charged Tracks) Direction and the Neutral Energy Direction





Note that the tail of the PSF is strongly correlated with θ_{NC} with unit slope!

Also this correlated tail is more pronounced in the Thick Radiators (as expected!)

And... the tail extends from one side of the FoV to the other!

Neut. Energy Implementation

<u>*First Task:*</u> Determine what the errors are as a function of energy Naively one expects the location error to ~ $R_{Moliere}/\sqrt{E}$

Run γ s [100 MeV, 10 GeV] in a patch (x,y) = ([100,150], [100,150]) at normal incidence ($\hat{e}_z = -1$)

ermi

Gamma-ray Space Telescope



Second Task: Combine Neutral energy direction covariently with Charged Soln.

$$\mathbf{P}_{\text{SUM}} = \mathbf{C}_{\text{SUM}} (\mathbf{C}_{\text{Charged}}^{-1} \mathbf{P}_{\text{Charged}} + \mathbf{C}_{\text{Neutral}}^{-1} \mathbf{P}_{\text{Neutral}})$$

where P_i and C_i are the 4-element parameter vectors and the 4x4 *Covariance Matrices* and $C_{SUM} = (C_{Charged}^{-1} + C_{Neutral}^{-1})^{-1}$





Event Selection: CTBClassLevel > 0 & CTBCORE > .1

Events Using Neutral Energy Solution



Comments

- > 44% events use the Neutral Energy Solution
- The effectiveness of using Neutral Energy increases with increasing energy
- > The far tails on the PSF are reigned in



Making Covariance Usable

The Chicken and Egg problem

In this case the "*Chickens*" are the tracks and the "*Eggs*" are the energy determinations for the tracks.

Tracking needs an estimation of the event energy.

Energy reconstruction benefits greatly from knowing about the tracks.



Trouble: Final Energy can be factors off from the energy used in the 2nd Kalman Fit

Fixing the Energy Problem

Elements of the Cov. Matrix scales as 1/E² below ~ 10 GeV (*Mult. Scat. Dominated*)

Suggests
$$C_{\text{FINAL}} \cong C_{\text{FIT}} \cdot \left(\frac{E_{\text{FINAL}}}{E_{\text{FIT}}}\right)^2$$

Best agreement found with

$$C_{\text{FINAL}} \cong C_{\text{FIT}} \cdot \left(\frac{E_{\text{FINAL}}}{E_{\text{FIT}}}\right)^{1.6}$$

The power 1.6 is consistent with PSF behavior $(1/E^{.8})$ (*Tyrel Johnson, 2007*)

Correct Soln.: Iterate Fit with Final Energy

Find iteration ONLY affects Cov. Matrix, NOT the Parameters!

<u>Covariant Errors on the Sky</u>

Transformation of the Track Covariance Matrix to Sky Coordinates

R.P. Johnson September 19, 2007

The Kalmin fit gives us a covariance matrix in terms of the slope and intercept in each projection. Only the two slopes are relevant to the photon direction, so the first step is to reduce the error matrix from 4×4 to 2×2 , simply by removing the rows and columns related to the intercepts. Let's call the two slopes s_1 and s_2 . The associated 2×2 covariance matrix is σ_s^2 .

Let \hat{v} be the unit vector (direction cosines) denoting the *downward* (per LAT convention) photon direction in the LAT coordinate system and \hat{w} the corresponding unit vector in the sky coordinates (galactic coordinates, for example). Finally, let $\Theta = \{\ell, b\}$ be the photon longitude and latitude in the sky coordinates.

The transformation from \hat{v} to \hat{w} is linear and can be represented by a 3 × 3 orthogonal matrix **R** that includes a space inversion (i.e. negative determinant). The transformations from **s** to \hat{v} and from \hat{w} to Θ are nonlinear, but for purposes of transforming the covariance matrix, all we need are the first-derivative matrices. Let **A** represent the 3 × 2 derivative matrix for the transformation $\mathbf{s} \rightarrow \hat{\mathbf{0}}$ and **B** the 2×3 derivative matrix for $\hat{w} \rightarrow \Theta$. The 2×2 matrix for the overall transformation $\mathbf{s} \rightarrow \Theta$ then is $\mathbf{M} = \mathbf{BRA}$, and the covariance matrix transforms as

 $\sigma_{\Theta}^2 = \mathbf{M} \sigma_s^2 \mathbf{M}^T$.

The transformation equations corresponding to ${\bf A}$ can be written

$$\hat{v}_z = -\left[1 + s_1^2 + s_2^2\right]^{-1/2}$$

 $\hat{v}_x = \hat{v}_z s_1$
 $\hat{v}_y = \hat{v}_z s_2$

and the matrix $\mathbf{A} = \partial \hat{v} / \partial \mathbf{s}$ can be written in terms of the direction cosines \hat{v} as

P

$$\mathbf{A} = \begin{pmatrix} \hat{v}_{z} \left(1 - \hat{v}_{x}^{2} \right) & -\hat{v}_{x} \hat{v}_{y} \hat{v}_{z} \\ -\hat{v}_{x} \hat{v}_{y} \hat{v}_{z} & \hat{v}_{z} \left(1 - \hat{v}_{y}^{2} \right) \\ -\hat{v}_{x} \hat{v}_{z}^{2} & -\hat{v}_{y} \hat{v}_{z}^{2} \end{pmatrix} \,.$$

The transformation equations corresponding to ${\bf B}$ can be written

$$= \tan^{-1}\left(\frac{\hat{w}_y}{\hat{w}_x}\right)$$

 $b = \sin^{-1}\hat{w}_x$

and the matrix $\mathbf{B} = \partial \Theta / \partial \hat{w}$ is

$$\mathbf{B} = \begin{pmatrix} \frac{-\hat{w}_y}{\hat{w}_x^2 + \hat{w}_y^2} & \frac{\hat{w}_x}{\hat{w}_x^2 + \hat{w}_y^2} & 0\\ 0 & 0 & \frac{1}{\sqrt{1 - \hat{w}_x^2}} \end{pmatrix}$$

These equations are singular at the galactic pole ($\hat{w}_z = 1$), which has to be protected, and also at $\hat{v}_z = 0$, which should never occur, given that the Tracker cannot find horizontal tracks.



<u>A First Example: Mrk 421</u>

Neronov & Collaborators selected events around Mrk 421 with E > 100 GeV. This avoids multiple scattering issues. They applied Johnson's Inst.-2-Sky Transform.



Results:

- 1) Error Ellipses (sort of) point back towards Mrk 421
- 2) Smallest Ellipses consistent with limiting track resolution ($\approx \frac{\sigma_{\text{Pitch}} \cdot \sqrt{2}}{60 cm} \approx .01 \text{ deg.}$)
- 3) Covariant Localization ~2 better then PSF Localization





- Fermi-LAT is a highly optimized HEP detector for detecting Gamma Rays in the Space Environment.
- The Reconstruction of the Directions on the Sky were realized through 2nd order (*covariance!*).
- A covariant analysis of source images will bring to bare the full power of the LAT.
- Pass 8 will provide the "correct" covariance matrix using the Final Energy.
- > And... just perhaps we'll finally see Pair Halos